

MENHYDRO 2010

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NON LINEAR SYSTEMS
OF CONSERVATION LAWS

Linear Systems - Hyperbolicity

Let A be a constant matrix in $\mathcal{M}(p)$
and $W_0 \in L^\infty(\mathbb{R})^p$, consider the system:

$$\begin{cases} \frac{\partial W}{\partial t} + A \frac{\partial W}{\partial x} = 0 & x \in \mathbb{R} \\ & t > 0 \\ W(x, 0) = W_0(x) \end{cases}$$

We assume that the system is strictly hyperbolic

$$\rightarrow A = R \Lambda R^{-1}, \quad \Lambda = \text{Diag}(\lambda_1, \lambda_2, \dots, \lambda_p)$$

and $\lambda_1 < \lambda_2 < \dots < \lambda_p$

$R = [r_1, r_2, \dots, r_p]$ is the right eigenvectors matrix

$$\text{i.e.: } A r_j = \lambda_j r_j$$

The rows l_i of R^{-1} are the left eigenvectors of the system:

$$\text{i.e.: } l_i A = \lambda_i l_i \quad \text{and} \quad l_i r_j = \delta_{ij}$$

$$R^{-1} = \begin{bmatrix} - & l_1 & - \\ - & l_2 & - \\ - & \vdots & - \\ - & l_p & - \end{bmatrix}$$

①

SOLUTION OF THE LINEAR SYSTEM

Proposition: The solution of the System (SL) is given by:

$$W(x, t) = \sum_{j=1}^P [P_j \cdot W_0(x - \lambda_j t)] \cdot \pi_j$$

Proof: let's make the change of variable $P_e: V = R^{-1} W$ ($U_j = P_j \cdot W$)

$$\Leftrightarrow W = R V = \sum_{j=1}^P U_j \pi_j$$

System (SL) \Rightarrow

$$\begin{cases} \frac{\partial V}{\partial t} + \Lambda \frac{\partial V}{\partial x} = 0 & \left(\frac{\partial U_j}{\partial t} + \lambda_j \frac{\partial U_j}{\partial x} = 0 \right) \\ V(x, 0) = R^{-1} W_0(x) = V_0(x) \end{cases}$$

$$U_j(x, t) = U_j(x - \lambda_j t, 0) = P_j \cdot W_0(x - \lambda_j t)$$

$$\text{and } W(x, t) = R V(x, t) = \sum_{j=1}^P U_j(x, t) \pi_j$$

②

SELF SIMILARITY

Proposition:

The solution of the Riemann problem

$$\frac{\partial W(x,t)}{\partial t} + \frac{\partial}{\partial x} F(W(x,t)) = 0$$

$$W(x,0) = W_0(x) = \begin{cases} w_L & \text{if } x < 0 \\ w_R & \text{if } x > 0 \end{cases}$$

is self-similar

$$\text{i.e.: } W(x,t) = H\left(\frac{x}{t}\right)$$

Proof: for $\alpha > 0$ put $y = \alpha x$ and $\tau = \alpha t$

$$\text{Let } U(y, \tau) = W(x, t) = W\left(\frac{y}{\alpha}, \frac{\tau}{\alpha}\right)$$

$$\text{Remark } U(y, \tau=0) = W_0\left(\frac{y}{\alpha}\right) = \begin{cases} w_L & \text{if } y < 0 \\ w_R & \text{if } y > 0 \end{cases}$$

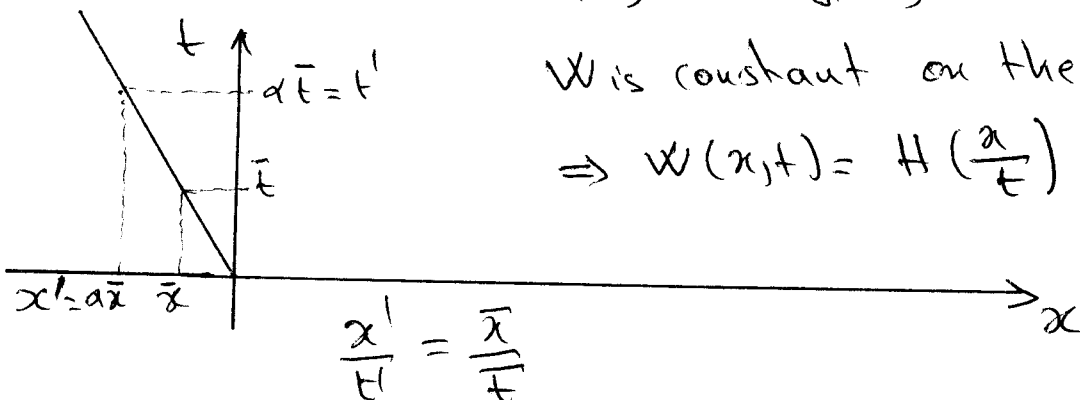
$$\frac{\partial U}{\partial \tau} = \frac{1}{\alpha} \frac{\partial W}{\partial t} \quad \text{and} \quad \frac{\partial F(U)}{\partial y} = \frac{1}{\alpha} \frac{\partial F(W)}{\partial x}$$

$$\text{So } \frac{\partial U}{\partial \tau} + \frac{\partial}{\partial y} F(U) = 0$$

$$\text{then } U(y, \tau) = W(y, \tau) = W(\alpha x, \alpha t) = W(x, t)$$

W is constant on the rays $\frac{x}{t} = \text{const}$

$$\Rightarrow W(x,t) = H\left(\frac{x}{t}\right)$$



THE RIEMANN PROBLEM

consider the initial value problem:

$$(RP) \begin{cases} \frac{\partial W}{\partial t} + A \frac{\partial W}{\partial x} = 0 \\ W_0(x) = \begin{cases} W_L & \text{if } x < 0 \\ W_R & \text{if } x > 0 \end{cases} \end{cases}$$

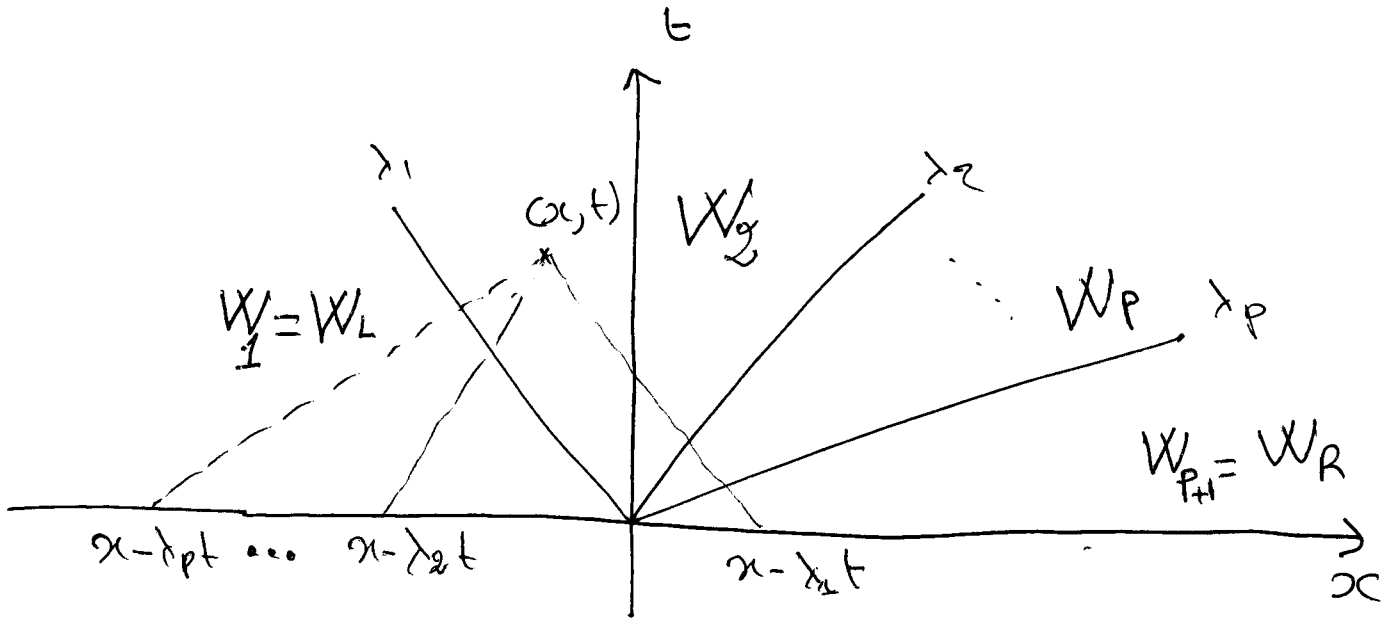
Remark: $W_R - W_L = \sum_{k=1}^p (\beta_k - \alpha_k) \pi_k$
if $W_L = \sum_{k=1}^p \alpha_k \pi_k$ and $W_R = \sum_{k=1}^p \beta_k \pi_k$

Proposition:

The solution of problem (RP) is made of constant states, separated by characteristic curves $C_k: \frac{x}{t} = \lambda_k$ in the frame (x, t) .

The solution shows a Jump $[W]_k = (\beta_k - \alpha_k) \pi_k$ across the k -characteristic C_k .
 λ_k is the speed of propagation of the discontinuity $[W]_k$.

Definition : The jump of the family k propagating at constant velocity λ_k , is called : k -wave .



Proof of The proposition :

Remark : $U_{0,k}(x) = P_k \sum_{d=1}^P \delta_j(x) \pi_j$

with $\delta_j(x) = \begin{cases} \alpha_j & \text{if } x < 0 \\ \beta_j & \text{if } x > 0 \end{cases}$

But $P_k \cdot \alpha_j = \delta_k \pi_j \Rightarrow U_k(x, t) = \delta_k(x - \lambda_k t)$

and $W(x, t) = \sum_{k=1}^P U_k(x, t) \pi_k$

gives : $\forall(x, t) / \frac{x}{t} \neq \lambda_k$

$$W(x, t) = \sum_{\frac{x}{t} < \lambda_k} \alpha_k \pi_k + \sum_{\frac{x}{t} > \lambda_k} \beta_k \pi_k \quad (5)$$

THE PHASE FRAME

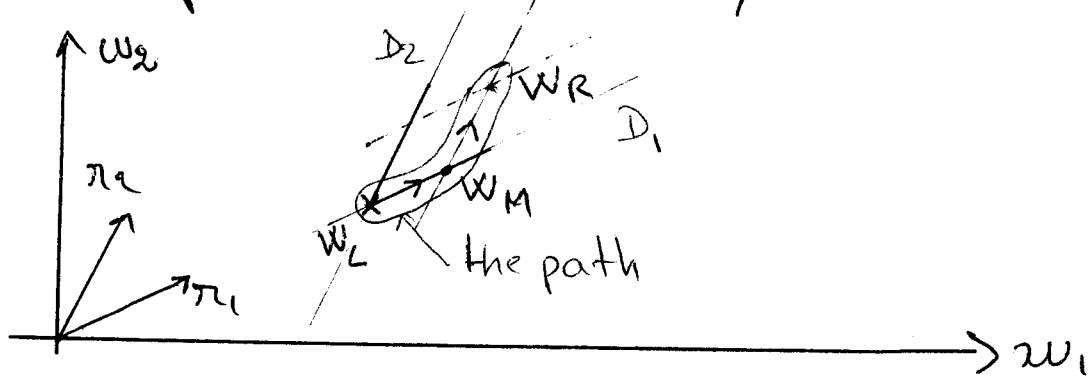
Remark:
$$W(x, t) = W_L + \sum_{\frac{x}{t} > \lambda_k} (\beta_k - \alpha_k) r_k$$

$$= W_R + \sum_{\frac{x}{t} < \lambda_k} (\beta_k - \alpha_k) r_k$$

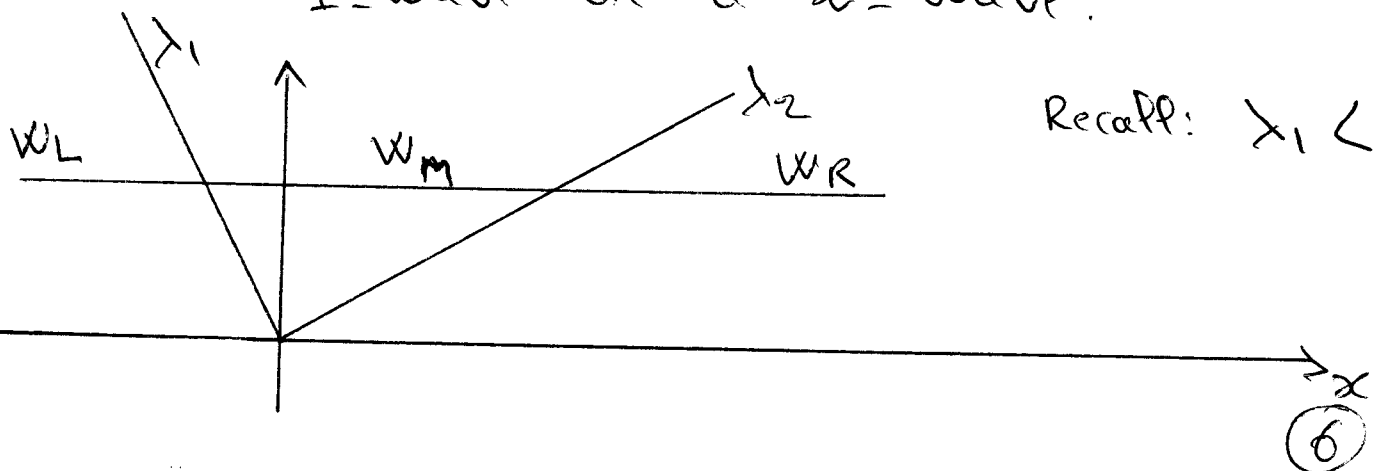
Remark: Solving the R.P. consists in a decomposition of the initial discontinuity in several jumps

$$W_R = W_L + \sum_{k=1}^{p'} (\beta_k - \alpha_k) r_k$$

Example with a 2×2 system:



Lines D_1 and D_2 give all the states that can be connected to W_L by a 1-wave or a 2-wave.



Recall: $\lambda_1 < \lambda_2$

NON LINEAR RIEMANN PROBLEM

consider the problem:

$$(NLRP) \left\{ \begin{array}{l} \frac{\partial W}{\partial t} + \frac{\partial}{\partial x} F(W) = 0 \\ W(x, 0) = W_0(x) = \begin{cases} W_L & \text{if } x < 0 \\ W_R & \text{if } x > 0 \end{cases} \end{array} \right.$$

$$F'(W) = A(W) \equiv \text{strictly } \mathbb{R}\text{-diagonalisable} \\ \lambda_1(W) < \lambda_2(W) < \dots < \lambda_p(W) \quad \forall W$$

Hugoniot Waves

goal: construct a weak solution
made of m discontinuities

propagating at speeds: $\lambda_1 < \lambda_2 < \dots < \lambda_m$

consider the discontinuity (\hat{W}, \tilde{W}) , speed λ

The jump (Rankine-Hugoniot) condition:

$$F(\tilde{W}) - F(\hat{W}) = \lambda (\tilde{W} - \hat{W})$$

m equations of $m+1$ unknowns $(\tilde{W} + \lambda)$

→ a one-parameter family solution

Like in linear case, one writes:

$$\tilde{W}_p = \hat{W} + \mu \vec{\pi}_p$$

$$\begin{cases} \tilde{W}_p(\mu, \hat{W}) = \hat{W} + \mu \vec{\pi}_p \\ \Delta_p(\mu, \hat{W}) \end{cases}$$

\tilde{W} connected to \hat{W} by a p -wave.

Remark: $\tilde{W}(0, \hat{W}) = \hat{W}$

Proposition: The curve (Hugoniot Locus)

$\tilde{W}_p(\mu, \hat{W})$ is tangent to the eigenvector $\vec{\pi}_p(\hat{W})$ at $\tilde{W} = \hat{W}$

Proof: R.H.C $\rightarrow F(\tilde{W}_p) - F(\hat{W}) = \Delta_p(\tilde{W}_p - \hat{W})$

$$\frac{d}{d\mu} \rightarrow A(\tilde{W}_p) \frac{d\tilde{W}_p}{d\mu} = \Delta_p \frac{d\tilde{W}_p}{d\mu} + (\tilde{W}_p - \hat{W}) \frac{d\Delta_p}{d\mu}$$

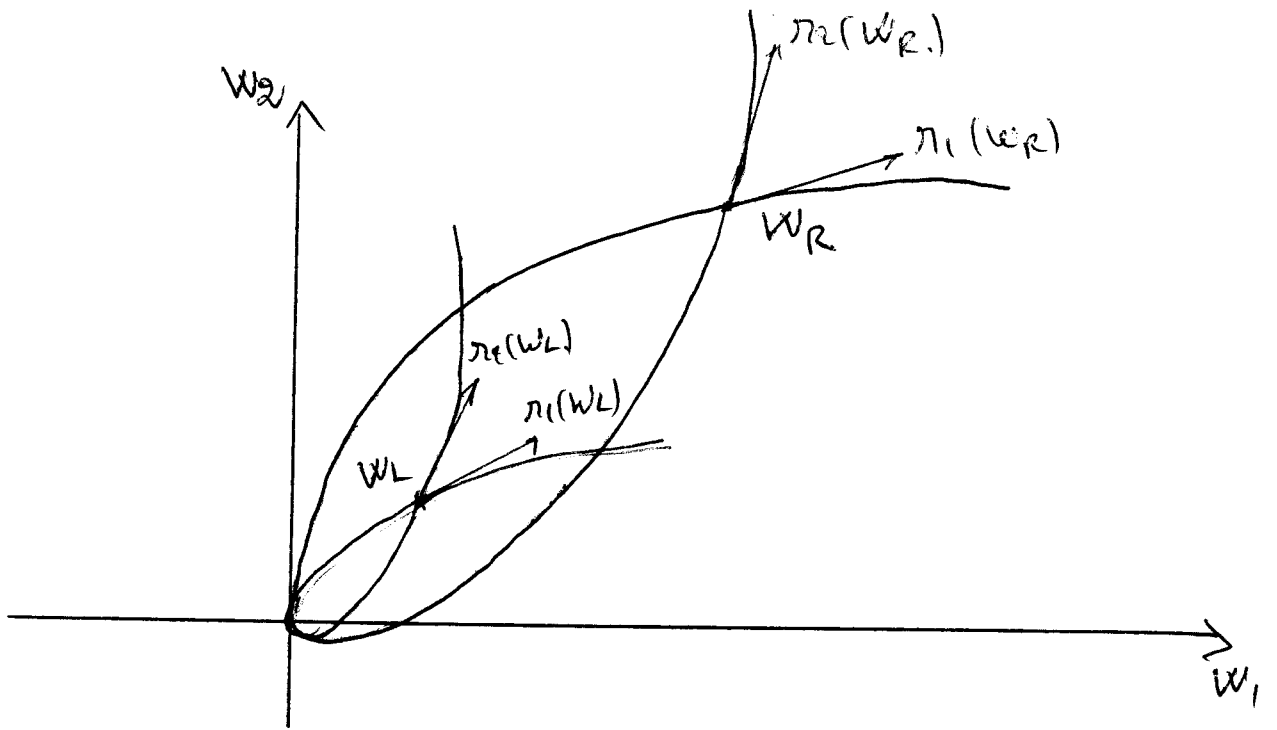
$$\mu=0 \rightarrow A(\hat{W}) \frac{d\tilde{W}_p(0, \hat{W})}{d\mu} = \Delta_p(0, \hat{W}) \frac{d\tilde{W}_p(0, \hat{W})}{d\mu}$$

Then $\frac{d\tilde{W}_p(0, \hat{W})}{d\mu}$ is an eigenvector of $A(\hat{W})$

associated to an eigenvalue $\Delta_p(0, \hat{W})$

but $\lambda_1 < \lambda_2 < \dots < \lambda_p$ and $\Delta_1 < \Delta_2 < \dots < \Delta_p$

$$\rightarrow \Delta_p(0, \hat{W}) = \lambda_p \quad \text{and} \quad \mu = p.$$



SHALLOW WATER PROBLEM

EXACT SOLUTION FOR DAM BREAK

consider the system:

$$\begin{cases} \frac{\partial h}{\partial t} + \frac{\partial}{\partial x} (hu) = 0 \\ \frac{\partial}{\partial t} (hu) + \frac{\partial}{\partial x} \left(hu^2 + \frac{g}{2} h^3 \right) = 0 \end{cases}$$

Note $q = hu$, $w = \begin{pmatrix} h \\ q \end{pmatrix}$, $F(w) = \begin{pmatrix} q \\ \frac{q^2}{h} + \frac{g}{2} h^3 \end{pmatrix}$

One gets: $\frac{\partial w}{\partial t} + \frac{\partial}{\partial x} F(w) = 0$

Then $F'(w) = A(w) = \begin{pmatrix} 0 & 1 \\ c^2 - u^2 & 2u \end{pmatrix}$

with $c = \sqrt{gh}$

So $A(w) = R(w) \Lambda(w) R^{-1}(w)$

$\lambda_1(w) = u - c$ $\lambda_2(w) = u + c$

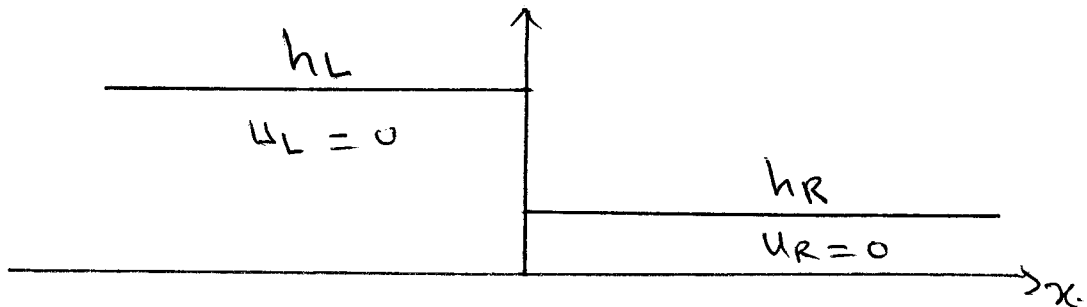
$\pi_1(w) = \begin{pmatrix} 1 \\ \lambda_1(w) \end{pmatrix}$ $\pi_2(w) = \begin{pmatrix} 1 \\ \lambda_2(w) \end{pmatrix}$

Rankine - Hugoniot \Rightarrow

$$[F(w)] = s [w]$$

$$\left\{ \begin{array}{l} \tilde{q} - \hat{q} = s (\tilde{h} - \hat{h}) \\ \left(\frac{\tilde{q}^2}{\tilde{h}} + \frac{g}{2} \tilde{h}^2 \right) - \left(\frac{\hat{q}^2}{\hat{h}} + \frac{g}{2} \hat{h}^2 \right) = s (\tilde{q} - \hat{q}) \end{array} \right.$$

particular case : $\hat{u} = 0 \rightarrow \hat{q} = 0$

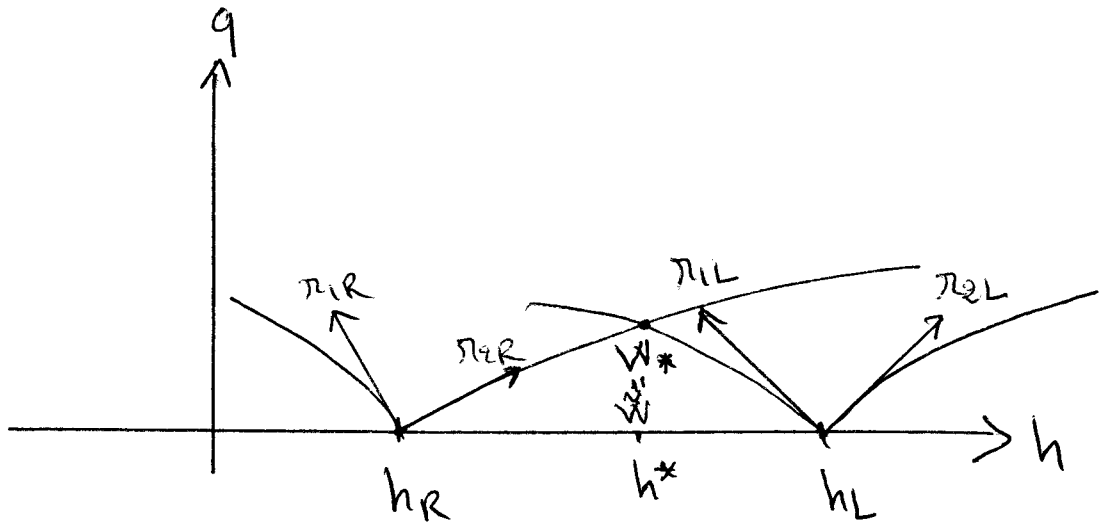


$$\tilde{q} = s (\tilde{h} - \hat{h}) \rightarrow s = \frac{\tilde{q}}{\tilde{h} - \hat{h}}$$

$$\frac{\tilde{q}^2}{\tilde{h}} + \frac{g}{2} (\tilde{h} - \hat{h}) = \frac{\tilde{q}^2}{\tilde{h} - \hat{h}}$$

$$\rightarrow \tilde{q} = \pm (\hat{h} - \tilde{h}) \tilde{h} \sqrt{\frac{g}{2} \left(\frac{1}{\hat{h}} + \frac{1}{\tilde{h}} \right)}$$

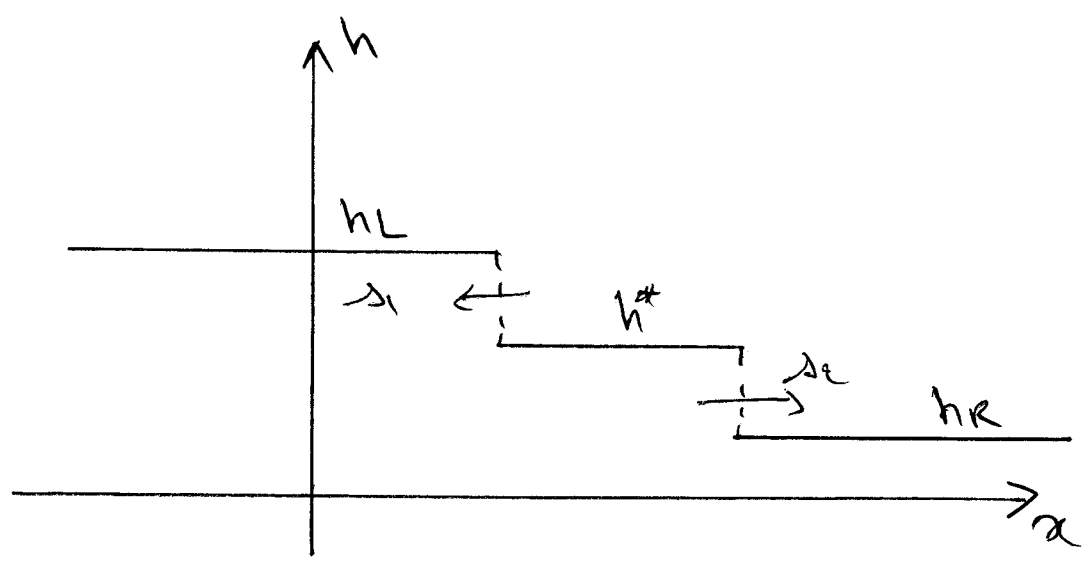
$$\tilde{u} = \pm (\hat{h} - \tilde{h}) \sqrt{\frac{g}{2} \left(\frac{1}{\hat{h}} + \frac{1}{\tilde{h}} \right)}$$



Then:

$$\left\{ \begin{array}{l} u^* = (h_L - h^*) \sqrt{\frac{g}{2} \left(\frac{1}{h_L} + \frac{1}{h^*} \right)} \\ s_1 = \frac{-u^* h^*}{h_L - h^*} \end{array} \right.$$

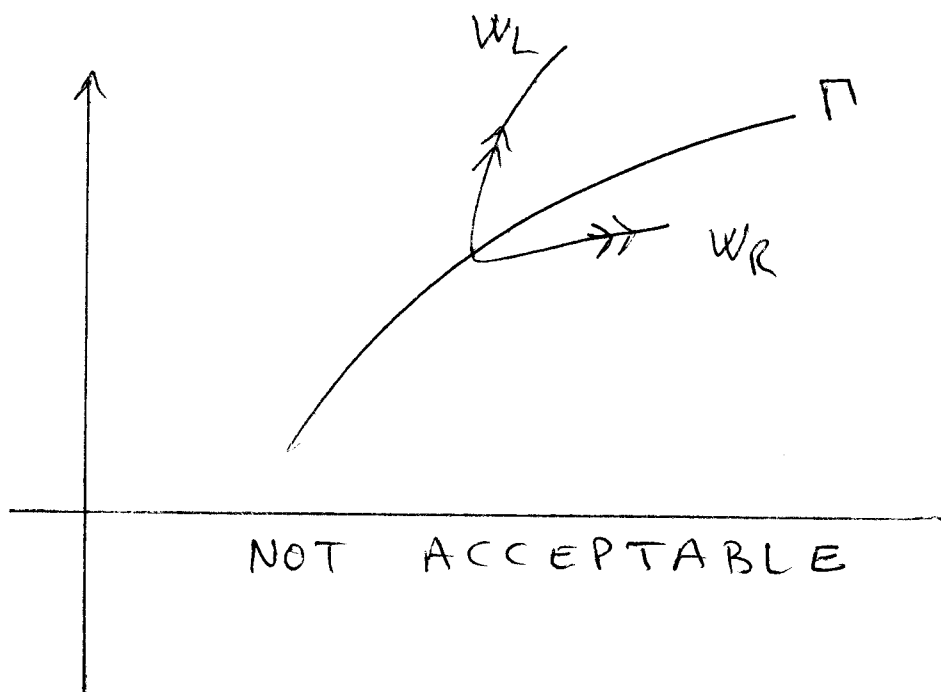
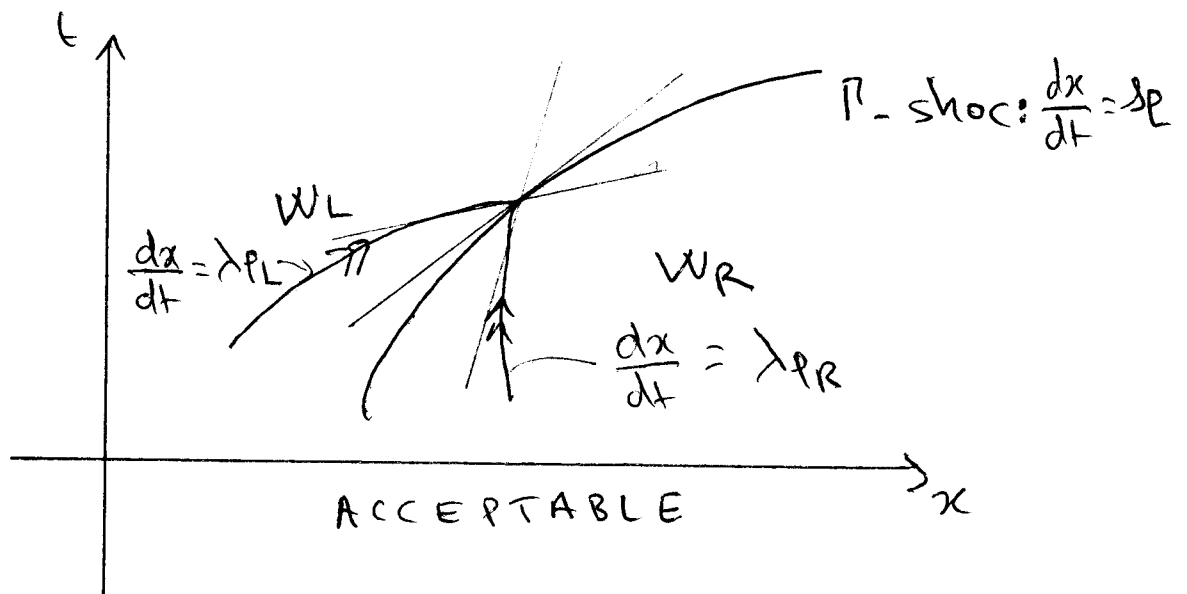
$$\left\{ \begin{array}{l} u^* = (h^* - h_R) \sqrt{\frac{g}{2} \left(\frac{1}{h_R} + \frac{1}{h^*} \right)} \\ s_2 = \frac{u^* h^*}{h^* - h_R} \end{array} \right.$$



Proposition: (Σ_{ax} Entropy condition)

A shock in a Π -wave family is admissible (entropy satisfying)

if: $\lambda_P(W_L) > \Delta P > \lambda_P(W_R)$
 $\lambda_{PL}^{-1} < \Delta P^{-1} < \lambda_{PR}^{-1}$



Hence: the first shoe (w_L, w^*)
is not acceptable

Indeed:

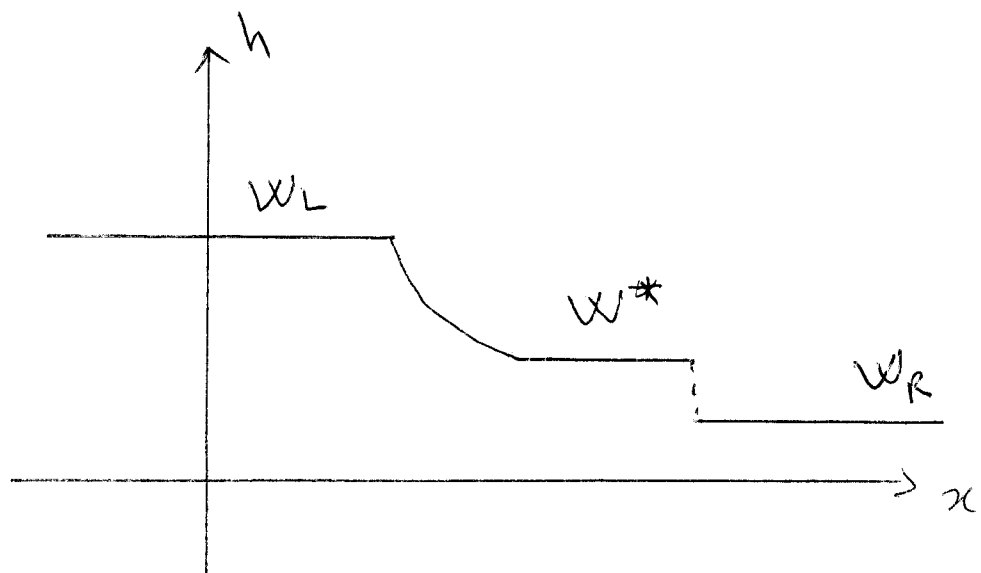
$$\lambda_P(w_L) > \Delta P > \lambda_P(w^*) \Rightarrow$$

$$0 - \sqrt{gh_L} \stackrel{\textcircled{1}}{>} \frac{-u^* h^*}{h_L - h^*} \stackrel{\textcircled{2}}{>} u^* - \sqrt{gh^*}$$

$$\textcircled{1} \Rightarrow \sqrt{gh_L} < h^* \sqrt{\frac{g}{2} \left(\frac{1}{h_L} + \frac{1}{h^*} \right)}$$

$$\Rightarrow \sqrt{gh_L} < \sqrt{gh^*} \frac{1}{2} (1 + h^*/h_L)$$

Impossible since $h^* < h_L$



RAREFACTION WAVE

Proposition: Let $w(x, t) = H\left(\frac{x}{t}\right)$

the solution of the R.P in

a region where it is regular,

then H is solution of the system:

$$H'(\xi) = \left[\nabla \lambda_p[H(\xi)] \cdot \vec{\eta}_p[H(\xi)] \right]^{-1} \vec{\eta}_p[H(\xi)]$$

$$\xi_1 < \xi < \xi_2$$

$$H(\xi_1) = W_L$$

Proof: $\frac{\partial w}{\partial t} + A(w) \frac{\partial w}{\partial x} = 0 \Rightarrow$

$$-\frac{x}{t^2} H'\left(\frac{x}{t}\right) + A\left(H\left(\frac{x}{t}\right)\right) \left(\frac{1}{t}\right) H'\left(\frac{x}{t}\right) = 0$$

$$\Rightarrow A[H(\xi)] H'(\xi) = \xi H'(\xi) / \xi = \frac{x}{t}$$

$$\Rightarrow H'(\xi) \text{ colinear to an eigenvector}$$

$$\Rightarrow \begin{cases} H'(\xi) = \alpha(\xi) \vec{\eta}_p[H(\xi)] & \textcircled{1} \\ \xi = \lambda_p[H(\xi)] & \textcircled{2} \end{cases}$$

$$\frac{d}{d\xi} \textcircled{2} \Rightarrow 1 = \nabla \lambda_P[H(\xi)] H'(\xi)$$

$$\textcircled{1} \Rightarrow 1 = \nabla \lambda_P[H(\xi)] \alpha(\xi) \vec{\pi}_P[H(\xi)]$$

$$\text{So } \alpha(\xi) = \underline{\left[\nabla \lambda_P[H(\xi)] \cdot \vec{\pi}_P[H(\xi)] \right]^{-1}}$$

Application to the Rarefaction
in the 1-wave family

$$\lambda_1(w) = u - c = \frac{q}{h} - \sqrt{gh} ; \pi_1 = \frac{1}{\lambda_1}$$

$$\nabla \lambda_1 \cdot \pi_1 = -\frac{2}{3} \sqrt{\frac{g}{h}} \Rightarrow$$

$$\int h'(\xi) = -\frac{2}{3} \sqrt{\frac{h}{g}}$$

$$\int q'(\xi) = -\frac{2}{3} \sqrt{\frac{h}{g}} \left(\frac{q}{h} - \sqrt{gh} \right) = \frac{2}{3} \left(h - \frac{q}{c} \right)$$

$$\text{with } \xi_1 = \xi(w_L) = \lambda_1(w_L) = 0 - \sqrt{gh_L}$$

$$h(\xi_1) = h_L \quad q(\xi_1) = 0$$

Solution of the O.D.E :

$$\left. \begin{aligned} h &= \frac{1}{9g} \left(2\sqrt{gh_L} - \frac{x}{t} \right)^2 \\ u &= \frac{2}{3} \left(\sqrt{gh_L} + \frac{x}{t} \right) \end{aligned} \right\} \frac{x}{t} = \xi$$

How to compute h^* and u^* ?

Remark: $\xi_2 = \lambda_1(w^*) = u^* - \sqrt{gh^*}$

$$\text{gives } \begin{cases} h^* = \frac{1}{3g} (2\sqrt{gh_L} + \sqrt{gh^*} - u^*)^2 \\ u^* = \frac{2}{3} (\sqrt{gh_L} - \sqrt{gh^*} + u^*) \end{cases}$$

$$\Rightarrow u^* = 2(\sqrt{gh_L} - \sqrt{gh^*}) \quad (1)$$

And from the g -shoc (w^*, w_R)

$$u^* = (h^* - h_R) \sqrt{\frac{g}{2} \left(\frac{1}{h_R} + \frac{1}{h^*} \right)} \quad (2)$$

(1) + (2) \rightarrow a non linear equation to solve, e.g. by Newton iterations

NON DIMENSIONALIZED FORM OF THE S-W. Equations

note \bar{W} the variables with dimension

$$(SWD) \left\{ \begin{array}{l} \frac{\partial \bar{h}}{\partial \bar{t}} + \frac{\partial}{\partial \bar{x}} (\bar{h} \bar{u}) = 0 \\ \frac{\partial}{\partial \bar{t}} (\bar{h} \bar{u}) + \frac{\partial}{\partial \bar{x}} \left(\bar{h} \bar{u}^2 + \frac{g}{2} \bar{h}^3 \right) = 0 \end{array} \right.$$

put $h = \bar{h} / h_L \Rightarrow \bar{h} = h_L \cdot h$

$u = \bar{u} / \bar{U}$ with $\bar{U} = \sqrt{g h_L}$

$x = \bar{x} / h_L \quad t = \bar{t} / t_0 \quad / \quad t_0 = \sqrt{\frac{h_L}{g}}$

Then (SWD) becomes

$$(SWND) \left\{ \begin{array}{l} \frac{\partial h}{\partial t} + \frac{\partial}{\partial x} (h u) = 0 \\ \frac{\partial}{\partial t} (h u) + \frac{\partial}{\partial x} \left(h u^2 + \frac{h^3}{2} \right) = 0 \end{array} \right.$$