

# Construction of a multi-soliton blow-up solution to the semilinear wave equation in one space dimension

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# Semilinear wave equations

$u : \mathbb{R}_t \times \mathbb{R}_x \rightarrow \mathbb{R}$  solution to

$$\begin{cases} \partial_{tt}u - \partial_{xx}u - |u|^{p-1}u = 0, \\ (u, \partial_t u)|_{t=0} = (u_0, u_1). \end{cases} \quad (\text{NLW})$$

where  $p > 1$ .

Local well-posedness:  $(u_0, u_1) \in H^1(\mathbb{R}) \times L^2(\mathbb{R})$  (Ginibre, Soffer & Velo, Lindblad & Sogge).

# Blow-up curve

Blow-up criterion (Levine 74):

$$\text{If } \int \left( \frac{1}{2}|u_1|^2 + \frac{1}{2}|\partial_x u_0|^2 - \frac{1}{p+1}|u_0|^{p+1} \right) dx < 0,$$

then the solution can not be global in time.

If  $u$  is a blow-up solution:

- let  $D \subset \mathbb{R}^2$  be the maximal domain of influence of  $u$  (in space-time), write  $D = \{(x, t) \mid 0 \leq t < T(x)\}$ .
- Blow-up curve  $\Gamma = \{(x, T(x))\}$ .
- $T$  is 1-Lipschitz (finite speed of propagation).
- $\bar{T} = \inf_{x \in \mathbb{R}} T(x)$  is the blow-up time.

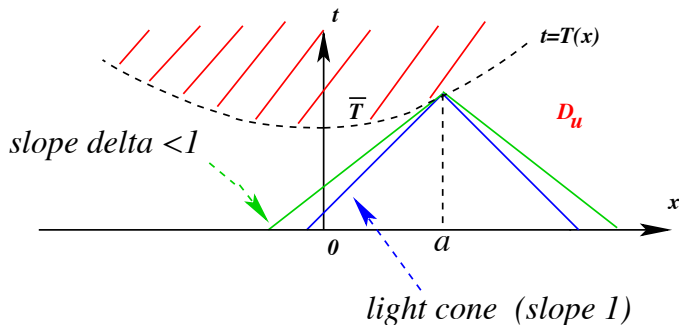
Goal:

- Description of *any* arbitrary blow-up solution;
- Construction of *examples* for each of the blow-up modalities.

# Characteristic points

A point  $a \in \mathbb{R}$  is *non-characteristic* if  $D$  contains a splaying cone

$$\mathcal{C}_\delta(a, T(a)) := \{(x, t) \in \mathbb{R} \times \mathbb{R}^+ \mid |x - a| \leq \frac{T(a) - t}{\delta}\} \subset D \text{ for some } \delta < 1.$$



A point is *characteristic* if it is not the case.

Notation:

- $\mathcal{R}$  is the set of non-characteristic points.
- $\mathcal{S}$  is the set of characteristic points.

Known results:

- $\mathcal{R}$  is never empty ( $\bar{x}$  such that  $T(\bar{x}) = \bar{T}$ ).
- $\mathcal{S}$  can be empty (Caffarelli and Friedman 85, 86).

Introduce the similarity variables: for any point  $(x_0, T_0) \in \bar{D}$

$$w_{x_0, T_0}(y, s) = (T_0 - t)^{\frac{2}{p-1}} u(x, t), \quad y = \frac{x - x_0}{T_0 - t}, \quad s = -\ln(T_0 - t).$$

Functional space:

$$\mathcal{H} = \left\{ (q, p) \left| \int_{-1}^1 (|p(y)|^2 + |\partial_y q|^2 (1 - y^2) + |q(y)|^2) (1 - y^2)^{\frac{2}{p-1}} dy < +\infty \right. \right\}.$$

Stationary solutions (for the “ $w$ ” equation):

$$\kappa(d, y) = \pm \kappa_0 \frac{(1 - d^2)^{\frac{1}{p-1}}}{(1 + dy)^{\frac{2}{d-1}}}, \quad \kappa_0 = \left( \frac{2(p+1)}{(p-1)^2} \right)^{\frac{1}{p-1}}.$$

# Regularity of $\mathcal{R}$

## Theorem (Merle & Z. 2007, 2008)

- $\mathcal{R}$  is open and  $T(x)$  is  $\mathcal{C}^1$  on  $\mathcal{R}$ .
- There exists  $\mu_0 > 0$ ,  $C_0$ , such that for all  $x_0 \in \mathcal{R}$ , there exist  $\varepsilon(x_0) = \pm 1$ ,  $s(x_0) \geq -\ln T(x_0)$  such that  $\forall s \geq s(x_0)$ ,

$$\left\| \begin{pmatrix} w_{x_0, T(x_0)}(s) \\ \partial_s w_{x_0, T(x_0)}(s) \end{pmatrix} - \varepsilon(x_0) \begin{pmatrix} \kappa(T'(x_0)) \\ 0 \end{pmatrix} \right\|_{\mathcal{H}} \leq C_0 e^{-\mu_0(s-s_0)}.$$

(Nouaili improved the regularity of  $T$  to  $\mathcal{C}^{1, \mu_0}$ ).

- Only one  $\kappa(d)$ .
- Its parameter is the slope of the blow-up curve.
- Exponential convergence to the profile.

# Description of $\mathcal{S}$ and refined asymptotics of characteristic blow-up points

## Theorem (Merle & Z. 2012, improved in Côte & Z. 2012)

- $\mathcal{S}$  is a discrete set.
- If  $x_0 \in \mathcal{S}$ , there exist  $k = k(x_0) \geq 2$ ,  $\varepsilon(x_0) = \pm 1$  and  $\zeta_0 = \zeta_0(x_0)$  s.t.  $\forall s \geq s_0$

$$\left\| \begin{pmatrix} w_{x_0, T(x_0)}(s) \\ \partial_s w_{x_0, T(x_0)}(s) \end{pmatrix} - \varepsilon(x_0) \sum_{i=1}^k (-1)^i \begin{pmatrix} \kappa(d_i(s)) \\ 0 \end{pmatrix} \right\|_{\mathcal{H}} \leq C_0 \left( \frac{s_0}{s} \right)^\eta \quad \text{for some } \eta > 0.$$

where  $d_i(s) = -\tanh \zeta_i(s)$  and  $\zeta_i(s)$  is defined on the next slide.

- Furthermore, the blow-up curve is corner shaped at  $x_0$ : for some  $\gamma = \gamma(p) > 0$ ,

$$T(x) = T(x_0) - |x - x_0| + \frac{\gamma e^{2\zeta_0 \operatorname{sgn}(x_0 - x)} |x - x_0| (1 + o(1))}{|\ln |x - x_0||^{\frac{(k-1)(p-1)}{2}}}.$$



# Illustration with hyperbolic coordinates

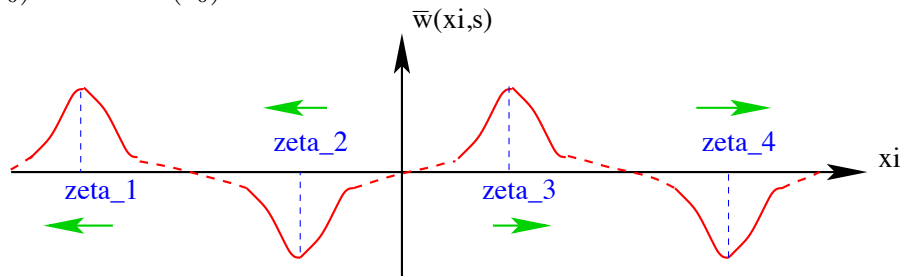
Introducing

$$\bar{w}_{x_0}(\xi, s) = (1 - y^2)^{\frac{1}{p-1}} w_{x_0}(y, s) \text{ with } y = \tanh \xi \text{ and } \zeta_i(s) = -\tanh^{-1} d_i(s),$$

we get

$$\|\bar{w}_{x_0}(\xi, s) - \epsilon(x_0) \kappa_0 \sum_{i=1}^{k(x_0)} (-1)^i \cosh^{-\frac{2}{p-1}}(\xi - \zeta_i(s))\|_{H^1 \cap L^\infty(\mathbb{R})} \rightarrow 0 \text{ as } s \rightarrow \infty,$$

and with  $k(x_0) = 4$  and  $\epsilon(x_0) = -1$ :



$(\zeta_i)_{i=1,\dots,k}$  is a solution to the system

$$\dot{\zeta}_i = e^{-\frac{2}{p-1}(\zeta_i - \zeta_{i-1})} - e^{-\frac{2}{p-1}(\zeta_{i+1} - \zeta_i)}, \quad i = 1, \dots, k,$$

with the convention  $\zeta_0(s) \equiv -\infty$ ,  $\zeta_{k+1}(s) \equiv +\infty$ , and barycenter

$\frac{1}{k}(\zeta_1(s) + \dots + \zeta_k(s)) = \zeta_0$ . One can compute explicitly for

$$\zeta_i(s) = \left( i - \frac{k+1}{2} \right) \frac{(p-1)}{2} \ln s + \alpha_i + \zeta_0(x_0),$$

where  $\alpha_i = \alpha_i(p, k)$  are chosen adequately.

Notice that

- The solitons are alternating.
- The number of solitons can be seen on the blow-up curve.
- The blow-up curve is never symmetric with respect to  $x_0$ , unless maybe if the barycenter of the solitons  $\zeta_0(x_0) = 0$ .

# Characteristic blow-up points with prescribed asymptotics

## Theorem (Côte & Z. 2012)

For any integer  $k \geq 2$  and  $\zeta_0 \in \mathbb{R}$ , there exists a blow-up solution  $u(x, t)$  in  $H^1 \times L^2(\mathbb{R})$  with  $0 \in \mathcal{S}$  such that  $T(0) = 1$  and

$$\left\| \begin{pmatrix} w_{0,1}(s) \\ \partial_s w_{0,1}(s) \end{pmatrix} - \sum_{i=1}^k (-1)^{i+1} \begin{pmatrix} \kappa(d_i(s)) \\ 0 \end{pmatrix} \right\|_{\mathcal{H}} \rightarrow 0 \quad \text{as } s \rightarrow \infty,$$

with

$$d_i(s) = -\tanh \zeta_i(s), \quad \zeta_i(s) = \left( i - \frac{k+1}{2} \right) \frac{(p-1)}{2} \ln s + \alpha_i + \zeta_0.$$

# Some ideas of proofs : DESCRIPTION, then CONSTRUCTION

Equation on  $w$ : let  $\rho = (1 - y^2)^{\frac{2}{p-1}}$  and  $\mathcal{L}w = \frac{1}{\rho} \partial_y (\rho (1 - y^2) \partial_y w)$ .

$$\partial_{ss} w = \mathcal{L}w - \frac{2(p+1)}{(p-1)^2} w + |w|^{p-1} w - \frac{p+3}{p-1} \partial_s w - 2y \partial_{ys}^2 w, \quad (\text{eqw})$$

Starting point: Monotonicity property in the  $w$  variable.

$$E(w) = \int_{-1}^1 \left( \frac{1}{2} |\partial_s w|^2 + \frac{1}{2} |\partial_y w|^2 (1 - y^2) + \frac{p+1}{(p-1)^2} w^2 - \frac{1}{p+1} |w|^{p+1} \right) \rho dy.$$

## Theorem (Lyapunov functional, Antonini & Merle 02)

$$E(w(s_2)) - E(w(s_1)) = -\frac{4}{p-1} \int_{s_1}^{s_2} \int_{-1}^1 |\partial_s w|^2 (1 - y^2)^{\frac{2}{p-1}-1} dy ds \leq 0.$$

If  $E(w(s_0)) < 0$  for some  $s_0 \in \mathbb{R}$ ,  $w$  blows-up in finite time.

# DESCRIPTION: Decomposition into a sum of solitons

- We have two bounds

$$\begin{aligned} & \|w_{x_0, T(x_0)}, \partial_s w_{x_0, T(x_0)}\|_{\mathcal{H}} \leq C, \\ & \int_{-\ln s_0}^{+\infty} \int_{-1}^1 |\partial_s w|^2 (1 - y^2)^{\frac{2}{p-1}-1} dy ds \leq C. \end{aligned}$$

- We find local limits: for some sequences  $s_n \rightarrow +\infty$ , in the  $\xi = \arg \tanh(y)$  variable,

$$w_{x_0, T(x_0)}(\xi + \xi_n, s + s_n) \rightarrow w^* \text{ stationary solution, in } H_{\text{loc}}^1.$$

- Nonzero stationary solutions are exactly the  $\pm \kappa(d, \cdot)$ . In the  $(\zeta, \xi) = (\arg \tanh d, \arg \tanh y)$  variables,

$$\kappa(\zeta, \xi) = \frac{\kappa_0(p)}{\cosh(\xi - \zeta)^{\frac{2}{p-1}}} \text{ is a soliton.}$$

## Proposition

There exist an integer  $k(x_0)$  and  $\varepsilon_i \in \{\pm 1\}$  and continuous functions  $d_i(s)$  such that

$$\left\| \begin{pmatrix} w_{x_0, T(x_0)}(s) \\ \partial_s w_{x_0, T(x_0)}(s) \end{pmatrix} - \sum_{i=1}^{k(x_0)} \varepsilon_i \begin{pmatrix} \kappa(d_i(s), \cdot) \\ 0 \end{pmatrix} \right\|_{\mathcal{H}} \rightarrow 0 \quad \text{as } s \rightarrow +\infty.$$

and, with  $\zeta_i = \arg \tanh(d_i)$ ,  $\zeta_{i+1}(s) - \zeta_i(s) \rightarrow +\infty$ .

At this point:

- We may have  $k(x_0) = 0$  or  $1$  or  $k(x_0) \geq 2$ .
- No control on the signs  $\varepsilon_i$ .
- If  $k(x_0) \geq 2$ , we have no control on the size  $\zeta_{i+1}(s) - \zeta_i(s)$ .

# DESCRIPTION at non-characteristics points

If  $x_0$  is a non-characteristic point

- Control in a splaying cone.
- A covering argument gets rid of the weight.

$$\|w_{x_0, T(x_0)}, \partial_s w_{x_0, T(x_0)}\|_{H^1 \times L^2} \leq C.$$

- We get one single limit:  $k(x_0) = 1$  (otherwise we quit  $H^1(-1, 1)$ ).
- Modulation + linear version of the Lyapunov yields

$$\forall s \geq s(x_0), \quad \left\| \begin{pmatrix} w_{x_0, T(x_0)}(s) \\ \partial_s w_{x_0, T(x_0)}(s) \end{pmatrix} - \varepsilon(x_0) \begin{pmatrix} \kappa(d(x_0)) \\ 0 \end{pmatrix} \right\|_{\mathcal{H}} \leq C_0 e^{-\mu_0(s-s_0)}.$$

- Stability property  $\rightarrow d(x_0) = T'(x_0)$  and  $\mathcal{R}$  is open.

# DESCRIPTION at characteristics points

- The covering argument does not hold:
- We may have  $k(x_0) \geq 2$ .

Equation on the solitons centers  $\zeta_i = \arg \tanh d_i$ :

$$\frac{1}{c_1(p)} \dot{\zeta}_i = \left( \varepsilon_i \varepsilon_{i-1} e^{-\frac{2}{p-1}(\zeta_i - \zeta_{i-1})} - \varepsilon_i \varepsilon_{i+1} e^{-\frac{2}{p-1}(\zeta_{i+1} - \zeta_i)} \right) + o(1).$$

- By construction  $\zeta_{i+1} - \zeta_i \rightarrow +\infty$ .
- This implies  $\varepsilon_i = (-1)^{i-1} \varepsilon_1$ .
- To study further the dynamics, we need an adequate modulated decomposition.
- Introduction of the  $\kappa^*(d, \nu, y)$ .



# Modulation

Define for  $\nu > -1 + |d|$ ,  $\kappa^*(d, \nu, y) = (\kappa_1^*(d, \nu, y), \kappa_2^*(d, \nu, y))$ , where

$$\kappa_1^*(d, \nu, y) = \kappa_0 \frac{(1 - d^2)^{\frac{1}{p-1}}}{(1 + dy + \nu)^{\frac{2}{p-1}}}, \quad \kappa_2^*(d, \nu, y) = \nu \partial_\nu \kappa_1^*(d, \nu, y).$$

- For  $\mu \in \mathbb{R}$ ,  $\kappa_1^*(d, \mu e^s, y)$  is solution of (eqw).
- If  $\mu = 0$ , it is  $\kappa(d, y)$ .
- If  $\mu > 0$ , it converges to 0 as  $s \rightarrow +\infty$ .
- If  $\mu < 0$ , it blows up at time  $s = \ln\left(\frac{|d|-1}{\mu}\right)$ .
- We can write a decomposition

$$\begin{pmatrix} w_{x_0, T(x_0)}(s) \\ \partial_s w_{x_0, T(x_0)}(s) \end{pmatrix} = q(s) + \sum_{i=1}^{k_0} (-1)^j \kappa^*(\hat{d}_i, \hat{\nu}_i), \quad \|q(s)\|_{\mathcal{H}} \rightarrow 0$$

with the projection  $\Pi_{\lambda, i}(q(s)) = 0$ , for all  $i = 1, \dots, k(x_0)$  and eigenvalues  $\lambda = 0, 1$  (we have  $2k$  nonnegative eigenvalues).

Define

$$J = \sum_{i=2}^k e^{-\frac{2}{p-1}(\zeta_i - \zeta_{i-1})}, \quad \bar{J} = \sum_{i=1}^k \frac{|\nu_i|}{1 - d_i^2}, \quad \hat{J} = \sum_{i=2}^k e^{-\frac{\bar{p}}{p-1}(\zeta_i - \zeta_{i-1})},$$

where  $\bar{p} = \min(p, 2 - 1/100)$ .

**Proposition (Dynamics of the parameters)**

$$\begin{aligned} \left| \frac{\dot{\nu}_i - \nu_i}{1 - d_i^2} \right| &\leq C(\|q\|_{\mathcal{H}}^2 + J + \|q\|_{\mathcal{H}} \bar{J}) \\ \left| \frac{1}{c_1(p)} \dot{\zeta}_i - \left( e^{-\frac{2}{p-1}(\zeta_i - \zeta_{i-1})} - e^{-\frac{2}{p-1}(\zeta_{i+1} - \zeta_i)} \right) \right| &\leq C(\|q\|_{\mathcal{H}}^2 + (J + \|q\|_{\mathcal{H}}) \bar{J} + J^{1+\eta}) \\ \|q(s)\|_{\mathcal{H}}^2 &\leq C e^{-\eta(s-s_0)} \|q(s_0)\|_{\mathcal{H}}^2 + C \hat{J}(s)^2 \end{aligned}$$

with  $\zeta_i(s) = -\arg \tanh(d_i(s))$ .

- The decomposition in generalized solitons  $\kappa^*$  is *stable* in some sense:

$$\zeta_i \sim \left( i - \frac{k+1}{2} \right) \frac{(p-1)}{2} \ln s, \quad J \sim s^{-2}, \quad \hat{J} \sim s^{-\bar{p}}, \quad \|q\|_{\mathcal{H}} \leq s^{-\bar{p}}, \quad \frac{|\nu_i|}{1 - d_i^2} \leq s^{-\bar{p}}.$$

# CONSTRUCTION: Prescribed non-characteristic blow-up

Parameters are given: number of solitons  $k \geq 2$  integer and their barycenter  $\zeta_0 \in \mathbb{R}$ . Define  $\bar{\zeta}_i, \bar{d}_i$  be the “perfect” centers of mass:

$$\bar{d}_i(s_0) = \tanh \bar{\zeta}_i(s), \quad \bar{\zeta}_i(s) = \left( i - \frac{k+1}{2} \right) \frac{(p-1)}{2} \ln s + \alpha_i.$$

Step 1: Construction of  $w$  decomposing into  $k$  solitons as  $s \rightarrow +\infty$ , without the condition on the barycenter.

Recall there is a 1 to 1 correspondence between  $w$  and  $(q, (d_i)_i, \nu_i)$  around a sum of  $k$  decoupled soliton.

Goal: Find initial conditions  $(q(s_0), (d_i(s_0))_i, (\nu_i(s_0))_i)$  such that  $w$  is defined on  $[s_0, +\infty)$  and

$$q(s) \rightarrow 0, \quad d_i(s) \sim \bar{d}_i(s) \quad \text{and} \quad \nu_i(s) \rightarrow 0 \quad \text{as} \quad s \rightarrow +\infty.$$

Equations describing the dynamics up to leading order:

$$\begin{aligned}\dot{\nu}_i &\sim \nu_i, \\ \frac{1}{c_1(p)} \dot{\zeta}_i &\sim \left( e^{-\frac{2}{p-1}(\zeta_i - \zeta_{i-1})} - e^{-\frac{2}{p-1}(\zeta_{i+1} - \zeta_i)} \right), \\ \|q(s)\|_{\mathcal{H}}^2 &\leq C e^{-\eta(s-s_0)} \|q(s_0)\|_{\mathcal{H}}^2 + C \hat{J}(s)^2.\end{aligned}$$

- $q$  has some (spectral) stability property: **Negative part of the spectrum  $\lambda \leq -\eta$** ,
- The  $\nu_i$  correspond to  $\lambda = 1$ . They are *transversally* unstable.
- The  $\zeta_i$  correspond to  $\lambda = 1$ . Fortunately, we have almost a Lyapunov Theorem for the ODE system of  $\zeta_i$ : stability property except for the barycenter.  
More precisely, consider the linearized system of  $\zeta_i$  around  $\bar{\zeta}_i$ .

Let  $\xi = (\zeta_i - \bar{\zeta}_i)_i$ , up to a linear change of variable  $\phi = P\xi$ , it writes

$$\dot{\phi} \sim \frac{1}{s} M \phi, \quad \text{where} \quad (M\phi, \phi) \leq - \sum_{i=2}^k \phi_i^2, \quad \text{and} \quad M\phi_1 = 0.$$

Hence, the  $\phi_i$  for  $i \geq 2$  are controlled, and  $\phi_1$  “doesn’t change much”.

**CONCLUSION:** we only need to control  $\nu_i$  for  $i = 1, \dots, k$ .

Define the rescaling  $\Gamma_s : (\nu_1, \dots, \nu_k) \mapsto (s^{-1/2-|\gamma_1|}\nu_1, \dots, s^{-1/2-|\gamma_k|}\nu_k)$  where

$$\gamma_i = \left(i - \frac{k+1}{2}\right) \frac{(p-1)}{2}.$$

Consider initial data of the type

$$(0, (\bar{d}_i(s_0))_i, (\nu_i(s_0))_i) \quad \text{that is} \quad w(s_0) = 0 + \sum_{i=1}^k \kappa^*(\bar{d}_i(s_0), \nu_i(s_0))$$

where  $\nu_i(s_0)$  belongs to  $\mathbb{B} := [-1, 1]^k$  after rescaling

$$(\nu_i(s_0))_i = \Gamma_{s_0}((\nu_{i,0})_i), \quad \nu_{i,0} \in [-1, 1].$$

For any  $\nu = (\nu_{i,0})_i$ , let  $(q(s), ((d_i(s))_i, (\nu_i(s))_i))$  the evolution with such initial conditions at  $s = s_0$ .

Define

- The rescaled flow  $\Phi : (s, \nu) \mapsto \Gamma_s^{-1}((\nu_1(s), \dots, \nu_k(s)))$ , and
- The exit time  $s^*(\nu) = \sup\{s \geq s_0 \mid \forall \tau \in [s_0, s], \Phi(\nu) \in \mathbb{B}\}$ .

Goal: find  $\nu \in \mathbb{B}$  such that  $s^*(\nu) = +\infty$ .

We argue by contradiction: assume that for all  $\nu \in \mathbb{B}$ ,  $s^*(\nu) < +\infty$ . Define

$$\Psi : \mathbb{B} \rightarrow \mathbb{B}, \quad \nu \mapsto \Phi(s^*(\nu), \nu).$$

Then, denoting  $\mathbb{S} = \partial\mathbb{B}$  the boundary

$$\textcircled{1} \quad \forall \nu \in \mathbb{B}, \quad \Psi(\nu) \in \mathbb{S}.$$

$$\textcircled{2} \quad \forall \nu \in \mathbb{B}, s \in [s_0, s^*(\nu)], \quad s^{1/2+\eta} \|q(s)\|_{\mathcal{H}} + \sum_{i=2}^k s^\eta |\phi_i|(s) + s_0^\eta |\phi_1(s)| \leq 1.$$

$$\textcircled{3} \quad \forall \nu \in \mathbb{S}, \quad s^*(\nu) = s_0 \text{ and } \Psi(\nu) = \nu.$$

$$\textcircled{4} \quad \Psi \text{ is continuous.}$$

We used:

- Stability properties for 2).
- Transversality of the flow on the boundary  $\mathbb{S}$  for 3) and 4).

$\Psi : \mathbb{B} \rightarrow \mathbb{S}$  continuous such that  $\Psi|_{\mathbb{S}} = \text{Id}$ . This contradicts Brouwer's Theorem.

Hence there exists  $\nu^\# \in \mathbb{B} = [-1, 1]^k$  such that  $s^*(\nu^\#) = +\infty$ .

- Conclusion: existence of  $w^\sharp \in \mathcal{C}([s_0, +\infty), \mathcal{H})$  satisfying the alternating  $k$ -soliton decomposition, with barycenter  $|\zeta_0| \leq s_0^\eta$ .
- Define  $u^\sharp \in H_{\text{uloc}}^1 \times L_{\text{uloc}}^2$  solution to (NLW), such that the trace on  $(-1, 1)$  is  $w^\sharp$ :

$$u^\sharp(0)|_{(-1,1)} = w(x, s), \quad \partial_t u^\sharp(x, 0)|_{(-1,1)} = \partial_s w^\sharp(x, s_0) + \frac{2}{p-1} w^\sharp(x, s_0) + x \partial_y w^\sharp(x, s_0).$$

- Check that for  $t \in [0, 1)$  and  $|x| < 1 - t$ ,

$$u^\sharp(x, t) = (1 - t)^{-\frac{2}{p-1}} w^\sharp \left( \frac{x}{1 - t}, s_0 - \ln(1 - t) \right).$$

$u^\sharp$  has characteristic point 0 as desired (without the barycenter condition).

- Fix barycenter using a Lorentz transform on  $u^\sharp$ .
- Corollary (prescribing multiple characteristic points) follows from finite speed of propagation.

# Conclusion

## Summary

- Complete description of the blow-up.
- Results are specific to semilinear equation:  
e.g.  $\partial_{tt}u - \partial_{xx}u = \partial_x u \partial_t u$ , explicit solution with blow-up curve  $\{(x, |x| + 1)\}$ ,  $\mathcal{S} = \mathbb{R}^*$ .
- Results are very sensitive to the nonlinearity, especially to sign change, but the method is robust:  
e.g.  $\partial_{tt}u - \partial_{xx}u = |u|^p$ : one always has  $\mathcal{S} = \emptyset$ .

## Some open questions

- Extension to  $\mathbb{R}^d$ ? Ok in the radial case.  
Problem in the general case: classification of stationary solutions.
- Can one construct a blow-up solution with prescribed characteristic set  $\mathcal{S}$ ?
- Given a 1-Lipshitz (smooth) curve  $\Gamma$ , can one give a solution  $u$  with blow-up curve  $\Gamma$ ?  
(See Killip & Visan).