

# Non characteristic points for the semilinear wave equation

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## The equation

$$\begin{cases} \partial_t^2 u = \partial_x^2 u + |u|^{p-1}u, \\ u(0) = u_0 \text{ and } u_t(0) = u_1, \end{cases}$$

where  $p > 1$ ,

$u(t) : x \in \mathbb{R} \rightarrow u(x, t) \in \mathbb{R}$ ,

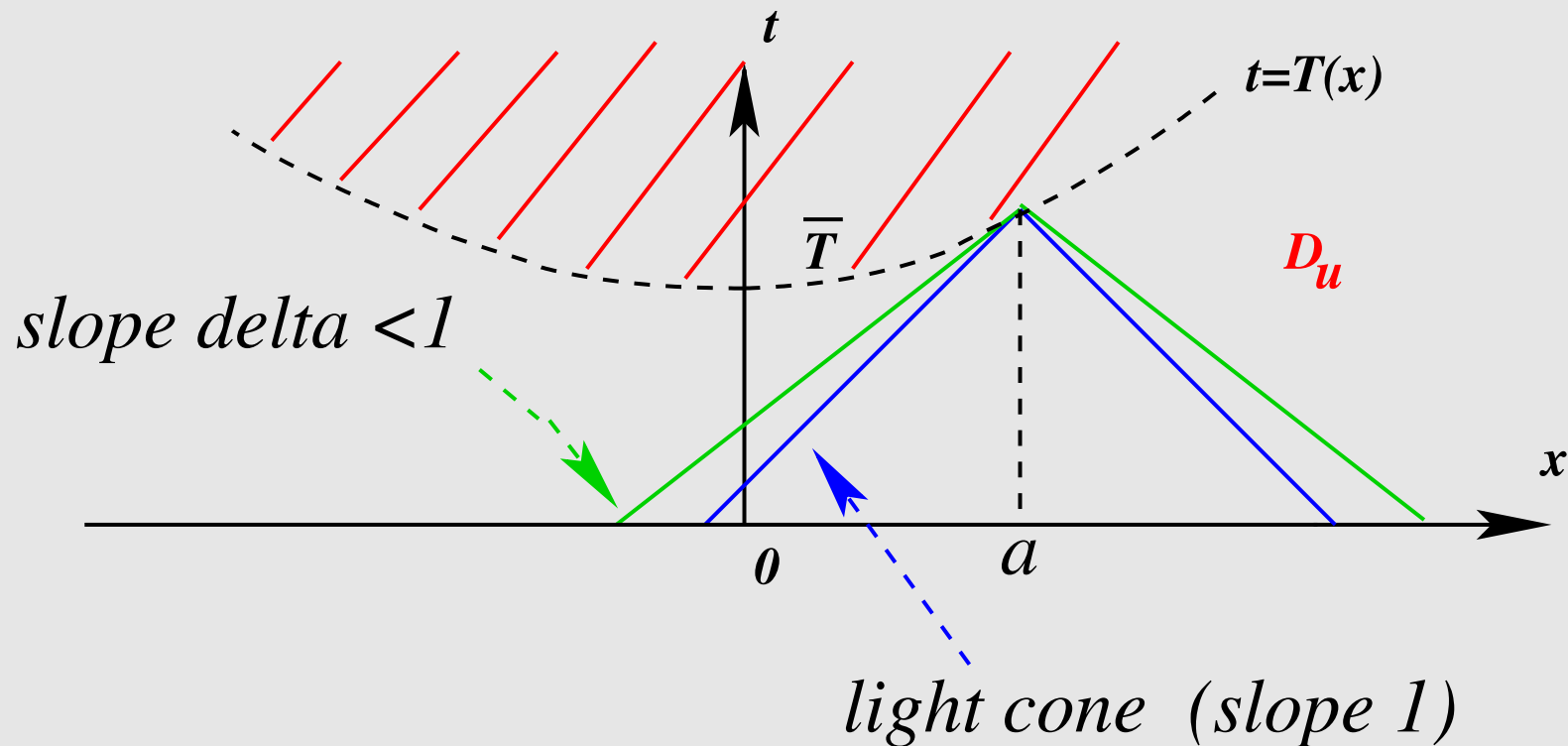
$u_0 \in H_{loc,u}^1(\mathbb{R})$  and  $u_1 \in L_{loc,u}^2(\mathbb{R})$

and

$$\|v\|_{L_{loc,u}^2(\mathbb{R})} = \sup_{a \in \mathbb{R}} \left( \int_{a-1}^{a+1} |v(x)|^2 dx \right)^{1/2}.$$

## Definition: Non characteristic points and characteristic points

A point  $a$  is said *non characteristic* if the domain contains a cone with vertex  $(a, T(a))$  and slope  $\delta < 1$ .



The point is said *characteristic* if not.

- Notation:  $\mathcal{R} \subset \mathbb{R}$  is the set of all *non* characteristic points.
- Notation:  $\mathcal{S} \subset \mathbb{R}$  is the set of all characteristic points ( $\mathcal{S} \cup \mathcal{R} = \mathbb{R}$ ).

## A Lyapunov functional and the blow-up rate

**Selfsimilar transformation for all  $x_0 \in \mathbb{R}$**

$$w_{x_0}(y, s) = (T(x_0) - t)^{\frac{2}{p-1}} u(x, t), \quad y = \frac{x - x_0}{T(x_0) - t}, \quad s = -\log(T(x_0) - t).$$

$(x, t)$  in the light cone of vertex  $(x_0, T(x_0)) \iff (y, s) \in (-1, 1) \times [-\log T(x_0), \infty)$ .

**Equation on  $w = w_{x_0}$ :** For all  $(y, s) \in (-1, 1) \times [-\log T(x_0), \infty)$ :

$$\begin{aligned} & \partial_{ss}^2 w - \frac{1}{\rho} \partial_y (\rho(1 - y^2) \partial_y w) + \frac{2(p+1)}{(p-1)^2} w - |w|^{p-1} w \\ &= -\frac{p+3}{p-1} \partial_s w - 2y \partial_{sy}^2 w \end{aligned}$$

$$\text{where } \rho(y) = (1 - |y|^2)^{\frac{2}{p-1}}$$

## A Lyapunov functional (Antonini-Merle)

$$E(w) = \int_{-1}^1 \left( \frac{1}{2} (\partial_s w)^2 + \frac{1}{2} (\partial_y w)^2 (1 - y^2) + \frac{(p+1)}{(p-1)^2} w^2 - \frac{1}{p+1} |w|^{p+1} \right) \rho dy,$$

Thanks to a Hardy-Sobolev inequality,  $E = E(w, \partial_s w)$  is well defined in the energy space

$$\mathcal{H} = \left\{ q \in H_{loc}^1 \times L_{loc}^2(B) \mid \|q\|_{\mathcal{H}}^2 \equiv \int_{-1}^1 \left( q_1^2 + (\partial_y q_1)^2 (1 - y^2) + q_2^2 \right) \rho dy < +\infty \right\}.$$

## Properties of the Lyapunov functional $E$

**Lemma 1 (Monotonicity (Antonini-Merle))** *For all  $s_1$  and  $s_2$ :*

$$E(w(s_2)) - E(w(s_1)) = -\frac{4}{p-1} \int_{s_1}^{s_2} \int_{-1}^1 (\partial_s w)^2 (1 - |y|^2)^{\frac{2}{p-1}-1} dy ds.$$

**Lemma 2 (A blow-up criterion)** *Consider a solution  $W$  such that  $E(W(s_0)) < 0$  for some  $s_0 \in \mathbb{R}$ , then  $W$  blows up in finite time  $S > s_0$ .*

## An upper bound on the blow-up rate in selfsimilar variables

**Th.** For all  $x_0 \in \mathbb{R}$  and  $s \geq -\log T(x_0) + 1$ ,

$$\int_{-1}^1 \left( \frac{1}{2} (\partial_s w)^2 + \frac{1}{2} (\partial_y w)^2 (1 - |y|^2) + \frac{(p+1)}{(p-1)^2} w^2 + \frac{1}{p+1} |w|^{p+1} \right) \rho dy \leq K$$

where the constant  $K$  depends only on  $p$  and an upper bound on  $T(x_0)$ ,  $1/T(x_0)$  and  $\|(u_0, u_1)\|$ .

### Getting rid of the weights

Reducing  $(-1, 1)$  to  $(-\frac{1}{2}, \frac{1}{2})$ , we get:

**Cor.** For all  $x_0 \in \mathbb{R}$  and  $s \geq -\log T(x_0) + 1$ ,

$$\int_{-\frac{1}{2}}^{\frac{1}{2}} \left( (\partial_s w)^2 + (\partial_y w)^2 + w^2 + |w|^{p+1} \right) dy \leq K.$$

## Upper bound in the original $u(x, t)$ variables

**Th. sup.** For all  $x_0 \in \mathbb{R}$  and  $t \in [\frac{3}{4}T(x_0), T(x_0))$ :

$$(T(x_0) - t)^{\frac{2}{p-1}} \frac{\|u(t)\|_{L^2(B(x_0, \frac{T(x_0)-t}{2}))}}{(T(x_0) - t)^{1/2}} \\ + (T(x_0) - t)^{\frac{2}{p-1} + 1} \left( \frac{\|u_t(t)\|_{L^2(B(x_0, \frac{T(x_0)-t}{2}))}}{(T(x_0) - t)^{1/2}} + \frac{\|\partial_x u(t)\|_{L^2(B(x_0, \frac{T(x_0)-t}{2}))}}{(T(x_0) - t)^{1/2}} \right) \leq K.$$

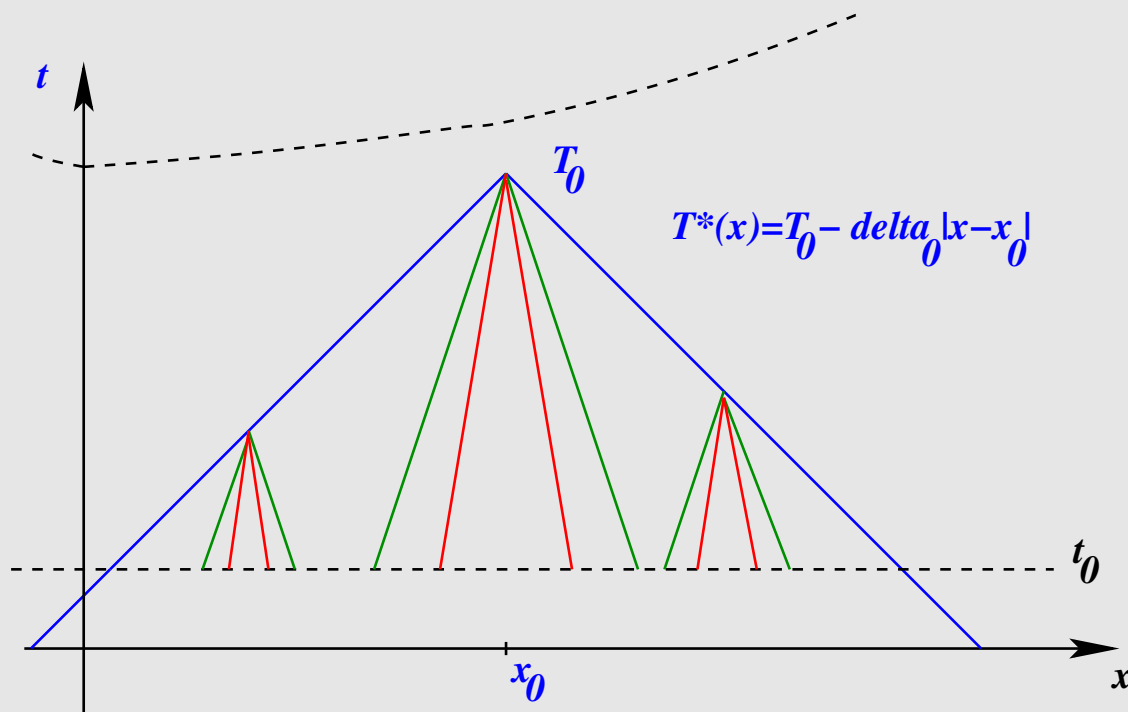
**Rk.** We have a lower bound of the same size when  $x_0$  is non characteristic (see Part 4 on profiles near a non characteristic point).



## Covering technique at a *non characteristic point*

If  $x_0$  is non characteristic point, then we can recover the estimate in whole section of the light-cone (or in the  $y$  variable, on the whole interval  $(-1, 1)$ ):

**Prop.** If  $x_0 \in \mathcal{R}$ , then for all  $s \geq -\log T(x_0)$ , we have  
 $\|(\omega_{x_0}(s), \partial_s \omega_{x_0}(s))\|_{H^1 \times L^2(-1,1)} \leq C_0.$



Blue: slope  $\delta_0 < 1$ , Green: slope 1 (light-cone), Red: slope 2.

## Asymptotic behavior at a *non characteristic* point

Take  $x_0 \in \mathbb{R}$  **non characteristic**. Using the energy structure, we obtain that  $\|(w_{x_0}(s), \partial_s w_{x_0}(s))\|_{H^1 \times L^2(-1,1)}$  is bounded.

**Question:** Does  $w_{x_0}(y, s)$  have a limit or not, as  $s \rightarrow \infty$  (that is as  $t \rightarrow T(x_0)$ ).

**Remark:** In the context of Hamiltonian systems, **this question is delicate**, and there is no natural reason for such a convergence, since **the wave equation is time reversible**.

See for similar difficulty and approach, results for

- ▷ the **critical KdV** (Martel and Merle),
- ▷ **NLS** (Merle and Raphaël).

## Stationary solutions.

We look for solutions of

$$\frac{1}{\rho} \left( \rho(1-y^2)w' \right)' - \frac{2(p+1)}{(p-1)^2} w + |w|^{p-1} w = 0.$$

We work in  $\mathcal{H}_0$ , the (stationary energy space) defined by

$$\mathcal{H}_0 = \left\{ r \in H_{loc}^1(-1,1) \mid \|r\|_{\mathcal{H}_0}^2 \equiv \int_{-1}^1 \left( r'^2(1-y^2) + r^2 \right) \rho dy < +\infty \right\}.$$

**Prop.** Consider a stationary solution in  $\mathcal{H}_0$ . Then, either  $w \equiv 0$  or there exist  $d \in (-1,1)$  and  $e = \pm 1$  such that  $w(y) = e\kappa(d,y)$  where

$$\forall (d,y) \in (-1,1)^2, \quad \kappa(d,y) = \kappa_0 \frac{(1-d^2)^{\frac{1}{p-1}}}{(1+dy)^{\frac{2}{p-1}}} \text{ and } \kappa_0 = \left( \frac{2(p+1)}{(p-1)^2} \right)^{\frac{1}{p-1}}.$$

**Remark:** We have 3 connected components.  $E(0) = 0 < E(\pm\kappa(d)) = E(\kappa_0)$ .

## Blow-up profile near a non characteristic point

**Th.** *There exist  $C_0 > 0$  and  $\mu_0 > 0$  such that if  $x_0$  is **non characteristic**, then there exist  $d(x_0) \in (-1, 1)$ ,  $e(x_0) = \pm 1$  and  $s^*(x_0) \geq -\log T(x_0)$  such that :*

(i) *For all  $s \geq s^*(x_0)$ ,*

$$\left\| \begin{pmatrix} w_{x_0}(s) \\ \partial_s w_{x_0}(s) \end{pmatrix} - e(x_0) \begin{pmatrix} \kappa(d(x_0), \cdot) \\ 0 \end{pmatrix} \right\|_{\mathcal{H}} \leq C_0 e^{-\mu_0(s-s^*)}$$

*and  $E(w_{x_0}(s)) \rightarrow E(\kappa_0)$  where the energy space*

$$\mathcal{H} = \left\{ q \in H_{loc}^1 \times L_{loc}^2(-1, 1) \mid \|q\|_{\mathcal{H}}^2 \equiv \int_{-1}^1 \left( q_1^2 + (q_1')^2 (1-y^2) + q_2^2 \right) \rho dy < +\infty \right\}.$$

(ii)  $d(x_0) = T'(x_0)$ .

**Rk.** We have exp. fast convergence (hence,  $C^{1,\mu_0}$  regularity of  $\mathcal{R}$ , see Nouaili).

**Rk.**  $\|w_{x_0}(y, s) - e(x_0)\kappa(d(x_0), y)\|_{L^\infty(-1,1)} \rightarrow 0$ .

**Rk.** The parameter of the profile  $d(x_0)$  has a geometrical interpretation  $(T'(x_0))$ .

## Difficulties of the proof of convergence

- ▷ The set of non zero stationary solutions is made up of non isolated solutions (one parameter family):  
→ we need a **modulation technique**.
- ▷ The linearized operator around a non zero stationary solution is **non self-adjoint**:  
→ we need to use dispersive properties coming from the Lyapunov functional to control the negative part of the spectrum.

## The aim of the talk: the proof the convergence

Consider  $x_0 \in \mathcal{R}$  and write  $w$  instead of  $w_{x_0}$ . We proceed in 3 parts:

- Part 1: Approaching the set of (non zero) stationary solutions.
- Part 2: Study of the linearized operator around a stationary solution and decomposition of the solution.
- Part 3: Convergence to a stationary solution.

## Part 1: Approaching the set of stationary solutions

We claim the following

**Prop.** For some  $d_0 \in (-1, 1)$  and  $e_0 = \pm 1$ , we have

$$\inf_{|d| < d_0} \left\| \begin{pmatrix} w(s) \\ \partial_s w(s) \end{pmatrix} - e_0 \begin{pmatrix} \kappa(d) \\ 0 \end{pmatrix} \right\|_{H^1 \times L^2(-1,1)} \rightarrow \text{as } s \rightarrow \infty.$$

Consider the set of stationary solutions  $Stat = \{0, \pm\kappa(d) \mid |d| < 1\}$ . Since

- ▷ for  $s$  large,  $0 < \epsilon_0(p) \leq \|(w(s), \partial_s w(s))\|_{H^1 \times L^2(-1,1)} \leq C_0$ ,
- ▷  $\|\kappa(d)\|_{H^1(-1,1)} \rightarrow +\infty$  as  $|d| \rightarrow 1$ ,
- ▷  $Stat$  is made of 3 connected components  $\{0\}$ ,  $\{\kappa(d) \mid |d| < 1\}$  and  $\{-\kappa(d) \mid |d| < 1\}$ ,

it is enough to prove that

$$\inf_{w^* \in Stat} \left\| \begin{pmatrix} w(s) \\ \partial_s w(s) \end{pmatrix} - \begin{pmatrix} w^* \\ 0 \end{pmatrix} \right\|_{H^1 \times L^2(-1,1)} \rightarrow \text{as } s \rightarrow \infty.$$

## The proof of Part 1

We proceed in 2 steps:

- In Step 1, we use compactness to prove the convergence in  $L^\infty(-1, 1)$ .
- In Step 2, we use the energy localization in the  $u$  variable to gain control in  $H^1(-1, 1)$ .



## Step 1: Compactness and convergence in $L^\infty(-1,1)$

Consider an arbitrary sequence  $s_n \rightarrow \infty$ . We will show that for some  $w^* \in S$  and up to a subsequence, we have

$$\left\| \begin{pmatrix} w(s_n) \\ \partial_s w(s_n) \end{pmatrix} - \begin{pmatrix} w^* \\ 0 \end{pmatrix} \right\|_{H^1 \times L^2(-1,1)} \rightarrow 0 \text{ as } n \rightarrow \infty.$$

From the bound

$$\| (w(s), \partial_s w(s)) \|_{H^1 \times L^2(-1,1)} \leq C_0$$

for  $s$  large enough and compactness, we see that for some  $(w^*, v^*) \in H^1 \times (-1,1)$  and up to a subsequence,

$$\begin{pmatrix} w(s_n) \\ \partial_s w(s_n) \end{pmatrix} \rightharpoonup \begin{pmatrix} w^* \\ v^* \end{pmatrix} \text{ weakly in } H^1 \times L^2(-1,1) \text{ as } n \rightarrow \infty.$$

and

$$\| w(s_n) - w^* \|_{L^\infty(-1,1)} \rightarrow 0 \text{ as } n \rightarrow \infty.$$

## Step 1: Compactness and convergence in $L^\infty(-1, 1)$ (cont.)

From the dissipation of the Lyapunov functional:

$$\int_{-\log T(x_0)}^{\infty} \int_{-1}^1 \frac{\partial_s w(y, s)^2}{1 - y^2} \rho(y) dy \leq E(w(-\log T(x_0))) \leq C_0,$$

we prove that

$$v^* \equiv 0, \quad w^* \in \text{Stat} \text{ and } w(y, s_n + s) \rightarrow w^*(y) \text{ as } n \rightarrow \infty$$

uniformly for  $|y| < 1$  and  $|s| \leq M$ , for any  $M > 0$ .

## Step 2: Convergence in $H^1$ through energy localization in the $u$ variable

Going to the  $u(x, t)$  variable, writing a Duhamel formulation in the light cone and coming back to the  $w(y, s)$  variable, we get a Duhamel formulation in  $w$ , for  $s \in [s_n - M, s_n]$ , yielding

$$\left\| \begin{pmatrix} w(s_n) \\ \partial_s w(s_n) \end{pmatrix} - \begin{pmatrix} w^* \\ 0 \end{pmatrix} \right\|_{H^1 \times L^2(-1,1)} \leq C(M) \|w(s_n - M) - w^*\|_{L^\infty(-1,1)} + C_0 e^{-\frac{2M}{p-1}}.$$

Fixing  $M$  then  $n$  large enough, we get to the conclusion which we recall:

**Prop.** For some  $d_0 \in (-1, 1)$  and  $e_0 = \pm 1$ , we have

$$\inf_{|d| < d_0} \left\| \begin{pmatrix} w(s) \\ \partial_s w(s) \end{pmatrix} - e_0 \begin{pmatrix} \kappa(d) \\ 0 \end{pmatrix} \right\|_{H^1 \times L^2(-1,1)} \rightarrow \text{as } s \rightarrow \infty.$$

## Part 2: Linearization around a possible limit

Let us assume that  $w(y, s) \rightarrow e_0 \kappa(d, y)$  in the energy space  $\mathcal{H}$ . Let

$$q(y, s) = w(y, s) - e_0 \kappa(d, y).$$

To simplify the notation, we assume that

$$e_0 = 1, \quad p = 2 \text{ and } w \geq 0.$$

## Part 2: Linearization around a possible limit (cont.)

For all  $s \geq -\log T(x_0)$ ,

$$\frac{\partial}{\partial s} \begin{pmatrix} q_1 \\ q_2 \end{pmatrix} = L_d \begin{pmatrix} q_1 \\ q_2 \end{pmatrix} + \begin{pmatrix} 0 \\ q_1^2 \end{pmatrix}$$

where

$$L_d \begin{pmatrix} q_1 \\ q_2 \end{pmatrix} = \begin{pmatrix} q_2 \\ \mathcal{L}q_1 + \psi(d, y)q_1 - 5q_2 - 2yq_2' \end{pmatrix},$$

$$\mathcal{L}q_1 = \frac{1}{\rho} \partial_y (\rho(1 - y^2) \partial_y q_1),$$

$$\psi(d, y) = 6 \left( 2 \frac{(1-d^2)}{(1+dy)^2} - 1 \right)$$

## Properties of the linear operator

- The operator  $L_d$  is not self-adjoint in the energy space.
- Its spectrum is given by

$$\lambda_n = 1 - n \text{ and } \mu_n = -6 - n, \quad n \in \mathbb{N}.$$

In particular, it has  $\lambda = 1$  and  $\lambda = 0$  as eigenvalues. The others are negative.

Two problems :

- ▷ How to control the zero eigenvalue? *By modulation.*
- ▷ How to control the negative part? *By a linear version of the Lyapunov functional.*

## Decomposition of the solution

For  $\lambda = 1$  or  $0$ , we introduce the eigenfunction  $F_\lambda^d(\mathbf{y})$  such that

$$L_d F_\lambda^d = \lambda F_\lambda^d$$

and the projector  $\pi_\lambda^d$  on  $F_\lambda^d$ .

**Rk.** We have

$$F_0^d(\mathbf{y}) = C(d) \begin{pmatrix} \partial_d \kappa(d, \mathbf{y}) \\ 0 \end{pmatrix}$$

and  $F_1^d$  is coming from the choice of the scaling time in the definition of  $w = w_{x_0}$ .

We decompose  $q$  as

$$q(\mathbf{y}, s) = \pi_1^d(q(s)) F_1^d(\mathbf{y}) + \pi_0^d(q(s)) F_0^d(\mathbf{y}) + q_-(\mathbf{y}, s).$$

## Control of the negative part

Let us introduce the symmetric bilinear form

$$\varphi_d(q, r) = \int_{-1}^1 (-q_1 (\mathcal{L}r_1 + \psi(d, y)r_1) + q_2 r_2) \rho dy$$

where  $\mathcal{L}r_1 + \psi(d, y)r_1$  already apperas in the definition of  $L_d$ :

$$L_d \begin{pmatrix} r_1 \\ r_2 \end{pmatrix} = \begin{pmatrix} r_2 \\ \mathcal{L}r_1 + \psi(d, y)r_1 - 5r_2 - 2yr'_2 \end{pmatrix}.$$



## Control of the negative part (cont.)

We recall the decomposition:

$$r(y) = \pi_1^d(r)F_1^d(y) + \pi_0^d(r)F_0^d(y) + r_-(y).$$

We claim the following:

**Prop.**

(i) If  $r \in \mathcal{H}$  and  $\pi_1^d(r) = \pi_0^d(r) = 0$ , then

$$\frac{1}{C_0} \|r\|_{\mathcal{H}}^2 \leq \varphi_d(r, r) \leq C_0 \|r\|_{\mathcal{H}}^2.$$

(ii) If  $r \in \mathcal{H}$ , then,

$$\frac{1}{C_0} \|r\|_{\mathcal{H}}^2 \leq \varphi_d(r_-, r_-) + \sum_{\lambda=0}^1 |\pi_\lambda^d(r)|^2 \leq C_0 \|r\|_{\mathcal{H}}^2.$$

## Modulation technique

I recall that we know that

$$\inf_{|d| < d_0} \left\| \begin{pmatrix} w(s) \\ \partial_s w(s) \end{pmatrix} - \begin{pmatrix} \kappa(d) \\ 0 \end{pmatrix} \right\|_{H^1 \times L^2(-1,1)} \rightarrow \text{as } s \rightarrow \infty$$

for some  $d_0 \in (-1, 1)$ . We want to prove the convergence to some  $\kappa(d^*, y)$ .

We introduce

$$q(y, s) = w(y, s) - \kappa(d(s), y)$$

where  $d(s) \in (-1, 1)$  is chosen so that

$$\pi_0^{d(s)}(q(s)) = 0 \text{ in } q(y, s) = \pi_1^{d(s)}(q(s))F_1^{d(s)}(y) + \pi_0^{d(s)}(q(s))F_0^{d(s)}(y) + q_-(y, s).$$

This is possible because  $F_0^d(y) = C(d) \begin{pmatrix} \partial_d \kappa(d, y) \\ 0 \end{pmatrix}$ .

## Modulation technique (cont.)

The decomposition becomes

$$q(y, s) = \pi_1^{d(s)}(q(s)) F_1^{d(s)}(y) + 0 + q_-(y, s)$$

and if we define

$$\alpha_1(s) = \pi_1^{d(s)}(q(s)) \text{ and } \alpha_-(s) = \sqrt{\varphi_{d(s)}(q_-, q_-)},$$

then

$$\frac{1}{C_0} \|q(s)\|_{\mathcal{H}} \leq |\alpha_1(s)| + |\alpha_-(s)| \leq C_0 \|q(s)\|_{\mathcal{H}}.$$

## Projection of the equation on the components of $q$

We recall the equation:

$$\frac{\partial}{\partial s} \begin{pmatrix} q_1 \\ q_2 \end{pmatrix} = L_d \begin{pmatrix} q_1 \\ q_2 \end{pmatrix} + \begin{pmatrix} 0 \\ q_1^2 \end{pmatrix} - d'(s) \begin{pmatrix} \partial_d \kappa(d) \\ 0 \end{pmatrix}.$$

If  $d(s) = \tanh \zeta(s)$ , then

$$\begin{aligned} |\zeta'(s)| &\leq C \|q(s)\|_{\mathcal{H}}^2, \\ |\alpha_1'(s) - \alpha_1(s)| &\leq C \|q(s)\|_{\mathcal{H}}^2, \\ \left(\frac{1}{2}\alpha_-(s)^2 + R_-\right)' &\leq -4 \int_{-1}^1 q_{-,2}^2 \frac{\rho}{1-y^2} dy + C \|q(s)\|_{\mathcal{H}}^3 \\ \text{with } |R_-(s)| &\leq C \|q(s)\|_{\mathcal{H}}^3 \\ \frac{d}{ds} \int_{-1}^1 q_1 q_2 \rho &\leq -\frac{4}{5} \alpha_-(s)^2 + C \int_{-1}^1 q_{-,2}^2 \frac{\rho}{1-y^2} dy + C \alpha_1^2. \end{aligned}$$

## Decreasing of the function

If

$$f(s) = \alpha_-(s)^2 + 2R_-(s) + \eta \int_{-1}^1 q_1 q_2 \rho$$

and  $\eta > 0$  is small, then

$$\begin{aligned} f'(s) &\leq -2\mu f(s), \\ \frac{1}{C_0} \|q(s)\|_{\mathcal{H}}^2 &\leq f(s) \leq C_0 \|q(s)\|_{\mathcal{H}}^2. \end{aligned}$$

Thus,

$$\|q(s)\|_{\mathcal{H}}^2 \leq C e^{-\mu s} \text{ with } q(y, s) = w(y, s) - \kappa(d(s), y).$$

But does  $w(y, s)$  converge?

## Convergence of the modulation parameter

Recalling that  $|\zeta'(s)| \leq C\|q(s)\|_{\mathcal{H}}^2$ , we see that  $\zeta(s)$  converges and so does  $d(s) = \tanh \zeta(s)$ .

Finally, we see that

$$\|w(y, s) - \kappa(d^*, y)\|_{\mathcal{H}} \leq Ce^{-\mu s}$$

for some  $d^* \in (-1, 1)$ .