# Isolatedness of characteristic points for blow-up solutions of a semilinear wave equation (Part 1)

Hatem ZAAG

CNRS & LAGA Université Paris 13

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joint work with Frank Merle, Université de Cergy-Pontoise and CNRS IHES

# The equation

$$\begin{cases} \partial_t^2 u = \partial_x^2 u + |u|^{p-1} u, \\ u(0) = u_0 \text{ and } u_t(0) = u_1, \end{cases}$$

where p>1,  $u(t):x\in\mathbb{R}\to u(x,t)\in\mathbb{R}$ ,  $u_0\in\mathrm{H}^1_{\mathrm{loc},\mathrm{u}}(\mathbb{R})$  and  $u_1\in\mathrm{L}^2_{\mathrm{loc},\mathrm{u}}(\mathbb{R})$  and

$$||v||_{\mathrm{L}^{2}_{\mathrm{loc,u}}(\mathbb{R})} = \sup_{a \in \mathbb{R}} \left( \int_{a-1}^{a+1} |v(x)|^{2} dx \right)^{1/2}.$$

# THE CAUCHY PROBLEM IN $H^1_{loc,u}({\rm I\!R}) \times L^2_{loc,u}({\rm I\!R})$

#### It is a consequence of:

- ▶ the Cauchy problem in  $H^1 \times L^2(\mathbb{R})$ ,
- the finite speed of propagation.

# **Maximal solution in** $H^1_{loc,u}(\mathbb{R}) \times L^2_{loc,u}(\mathbb{R})$

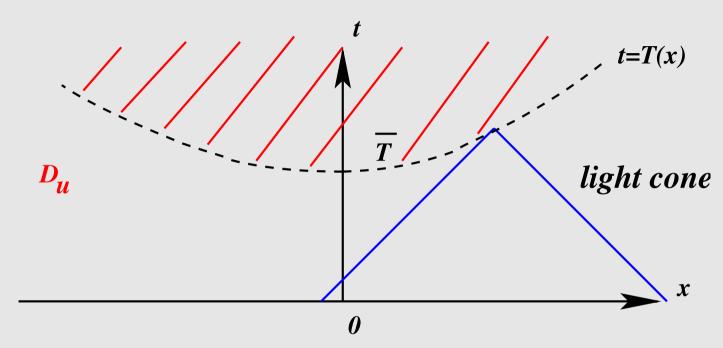
- either it exists for all  $t \in [0, \infty)$  (global solution),
- or it exists for all  $t \in [0, \bar{T})$  (singular solution).

#### **Existence of singular solutions**

It's a consequence of ODE techniques and the finite speed of propagation; see also the energy argument by Levine 1974:

if 
$$(u_0, u_1) \in H^1 \times L^2(\mathbb{R})$$
 and  $\int_{\mathbb{R}} \left( \frac{1}{2} (u_1)^2 + \frac{1}{2} (\partial_x u_0)^2 - \frac{1}{p+1} |u_0|^{p+1} \right) dx < 0$ , then  $u$  is not global.

# Singular solutions: the maximal influence domain



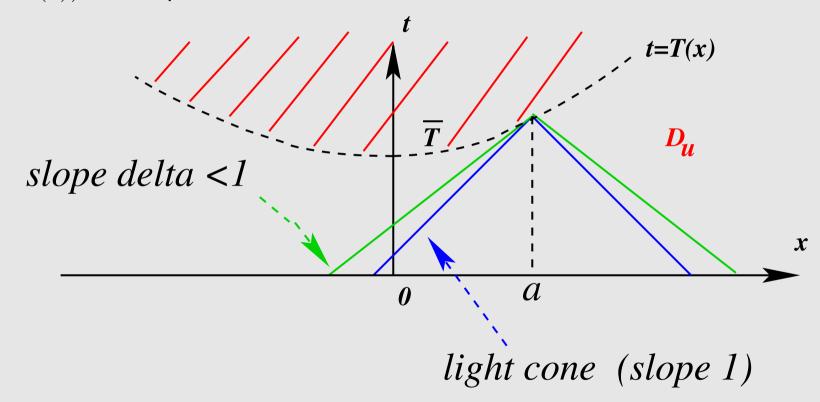
The blow-up set  $t \to T(x)$  is 1-Lipschitz (finite speed of propagation).

**Remark**:  $\bar{T} = \inf T(x)$  is the **blow-up time**. For all  $x \in \mathbb{R}^N$ , there exists a "local" blow-up time T(x).

The aim of this talk: To describe precisely the blow-up set, and the solution near the blow-up set, for an arbitrary blow-up solution.

# Definition: Non characteristic points and characteristic points

A point a is said  $non\ characteristic$  if the domain contains a cone with vertex (a, T(a)) and slope  $\delta < 1$ .



The point is said *characteristic* if not.

- Notation:  $\mathcal{R} \subset \mathbb{R}$  is the set of all *non* characteristic points.
- Notation:  $S \subset \mathbb{R}$  is the set of all characteristic points ( $S \cup \mathcal{R} = \mathbb{R}$ ).

# Known results, for an arbitrary solution

- The blow-up set  $\Gamma = \{(x, T(x))\} \subset \mathbb{R}^2$ .
- By definition,  $\Gamma$  is 1-Lipschitz.
- $\mathcal{R} \neq \emptyset$  (Indeed,  $\bar{x}$  such that  $T(\bar{x}) = \min_{x \in \mathbb{R}} T(x)$  is non characteristic).
- Caffarelli and Friedman (1985 and 1986) had two criteria to have  $\mathcal{R} = \mathbb{R}$  and  $x \mapsto T(x)$  of class  $C^1$  (using the positivity of the fundamental solution):
  - either when  $p \ge 3$ , with  $u_0 \ge 0$ ,  $u_1 \ge 0$  and  $(u_0, u_1) \in C^4 \times C^3(\mathbb{R})$ ,
  - or under conditions on initial data that ensure that

$$u \geq 0$$
 and  $\partial_t u \geq (1 + \delta_0) |\partial_x u|$ 

for some  $\delta_0 > 0$ .

### Questions and new results

- **Existence** 
  - Are there characteristic points? *yes*,  $\mathcal{S} \neq \emptyset$ .
- **Regularity** 
  - Is  $\mathcal{R}$  open? *yes*
  - Is  $\Gamma$  (or  $\Gamma_{\mathcal{R}}$ ) of class  $C^1$  ? *yes*
  - "How is" S? isolated points
  - How does  $\Gamma$  look like near S? corner shaped
- **▶** Asymptotic behavior (profile)
  - How does the solution behave near a non characteristic point? we have the profile
  - and near a characteristic point? we have a precise decomposition into solitons

Rk. Regularity and asymptotic behavior are linked.

# The plan

- Part 1: Existence of characteristic points.
- Part 2: A Liouville theorem and regularity of the blow-up set.
- Part 3: A Lyapunov functional and the blow-up rate.
- Part 4: Asymptotic behavior near *non characteristic* points (the blow-up profile).
- Part 5: Asymptotic behavior near *characteristic* points (decomposition into solitons).

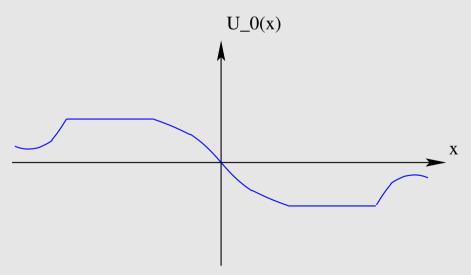
Rk. The order of this presentation goes from the easiest (to state) to the most complicated. The chronological order is actually 3, 4, 1, 2, 5.

# Part 1: Existence of characteristic points

**We recall**: Any solution to the Cauchy problem has (at least) a non characteristic point (the minimum of the blow-up set).

Th. There exist initial data which give solutions with a characteristic point.

**Example**: We take odd initial data, with two large plateaus of different signs. Then, the solution blows up, and the origin is a characteristic point with  $\forall t < T(0), u(0,t) = 0$ .



Th. If we perturb initial data, then then new solution blows up and has a characteristic point.

# Part 2: Regularity of the blow-up set

Near a non characteristic point:

**Th.** The set of non characteristic points  $\mathcal{R}$  is open and T(x) is of class  $C^1$  on this set  $(C^{1,\alpha}$  by N. Nouaili CPDE 2008).

Near a characteristic point:

**Th.** The set of characteristic points S is made of **isolated points**. If  $a \in S$ , then  $T'_1(a) = 1$  and  $T'_r(a) = -1$ .

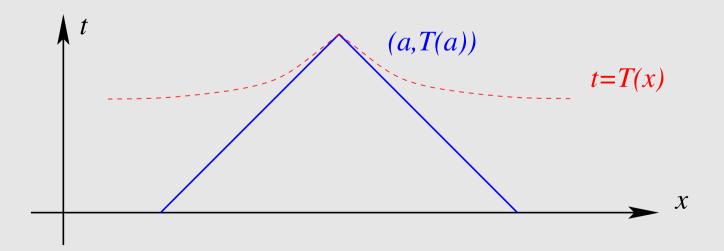
**Cor.** There is no solution with  $a \in S$  and T'(a) = 1.

**Th.** (the corner property) If  $a \in S$ , then for all x near a,

$$\frac{1}{C}|x-a||\log|x-a||^{-\gamma(a)} \le T(x) - T(a) + |x-a| \le C|x-a||\log|x-a||^{-\gamma(a)} \tag{1}$$

where

$$\gamma(a) = \frac{(k(a)-1)(p-1)}{2}$$
 with  $k(a) \in \mathbb{N}, k(a) \ge 2$ .



#### **Comments**

Rk. We recall the result of Caffarelli and Friedman:

If for all  $x \in \mathbb{R}$  and t < T(x), we have  $u(x,t) \ge 0$  and  $\partial_t u \ge (1 + \delta_0) |\partial_x u|$  for some  $\delta_0 > 0$ , then  $\mathcal{R} = \mathbb{R}$ .

Here, We improve their criterion:

If for all  $x \in [a,b]$  and t < T(x), we have  $u(x,t) \ge 0$ , then  $(a,b) \subset \mathcal{R}$ .

Idea of the proof of the regularity in the non characteristic case:

The techniques are based on

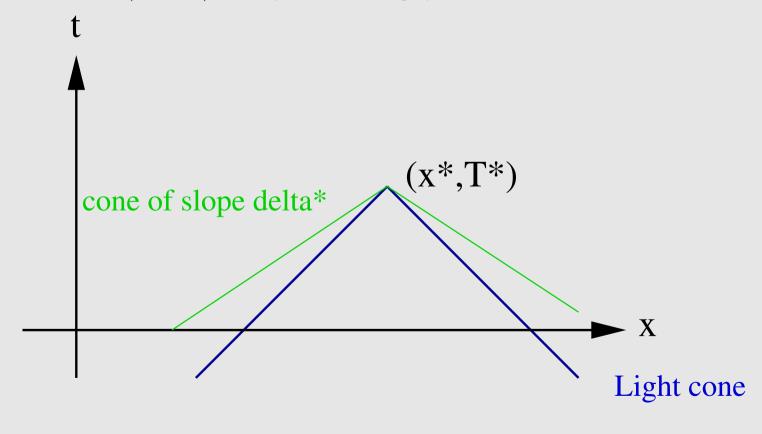
- a very good understanding of the behavior of the solution in selfsimilar variables in the energy space related to the selfsimilar variable (see Part 3 of this talk).
- a Liouville Theorem (see next slide).

Idea of the proof of the regularity in the characteristic case: At the end of the talk.

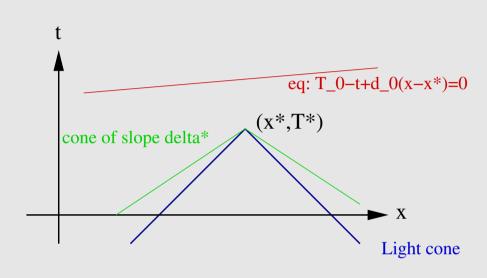
#### A Liouville Theorem

**Th.** Consider u(x,t) a solution of  $u_{tt} = u_{xx} + |u|^{p-1}u$  such that:

- *u* is defined in the *infinite* green cone,
- u is less than  $(T^* t)^{-\frac{2}{p-1}}$  (in  $L^2$  average).



#### A Liouville Theorem



Then,

- either  $u \equiv 0$ ,
- or there exists  $T_0 \ge T^*$ ,  $d_0 \in [-\delta_*, \delta_*]$  and  $\theta_0 = \pm 1$  such that u is actually defined below the red line by

$$u(x,t) = \theta_0 \kappa_0(p) \frac{(1 - d_0^2)^{\frac{1}{p-1}}}{(T_0 - t + d_0(x - x^*))^{\frac{2}{p-1}}}.$$

**Remark**: *u* blows up on the red line.

#### **Comments**

▶ The limiting case  $\delta^* = 1$  is still open.

#### The proof:

- The proof has a completely different structure from the proof for the heat equation.
- The proof is based on various energy arguments and on a dynamical result.

# Part 3: A Lyapunov functional and the blow-up rate

#### **Selfsimilar transformation for all** $x_0 \in \mathbb{R}$

$$w_{x_0}(y,s) = (T(x_0) - t)^{\frac{2}{p-1}} u(x,t), \ y = \frac{x - x_0}{T(x_0) - t}, \ s = -\log(T(x_0) - t).$$

(x,t) in the light cone of vertex  $(x_0,T(x_0)) \iff (y,s) \in (-1,1) \times [-\log T(x_0),\infty)$ .

**Equation on**  $w = w_{x_0}$ : For all  $(y,s) \in (-1,1) \times [-\log T(x_0), \infty)$ :

$$\partial_{ss}^2 w - \frac{1}{\rho} \partial_y (\rho (1 - y^2) \partial_y w) + \frac{2(p+1)}{(p-1)^2} w - |w|^{p-1} w$$

$$=-rac{p+3}{p-1}\partial_s w-2y\partial_{sy}^2 w$$

where 
$$\rho(y) = (1 - |y|^2)^{\frac{2}{p-1}}$$

# A Lyapunov functional (Antonini-Merle)

$$E(w) = \int_{-1}^{1} \left( \frac{1}{2} (\partial_s w)^2 + \frac{1}{2} (\partial_y w)^2 (1 - y^2) + \frac{(p+1)}{(p-1)^2} w^2 - \frac{1}{p+1} |w|^{p+1} \right) \rho dy,$$

Thanks to a Hardy-Sobolev inequality,  $E=E(w,\partial_s w)$  is well defined in the energy space

$$\mathcal{H} = \left\{ q \in H^{1}_{loc} \times L^{2}_{loc}(B) \mid \|q\|^{2}_{\mathcal{H}} \equiv \int_{-1}^{1} \left( q_{1}^{2} + \left( \partial_{y} q_{1} \right)^{2} (1 - y^{2}) + q_{2}^{2} \right) \rho dy < + \infty \right\}.$$

# — LATEX prosper —

# **Properties of the Lyapunov functional** *E*

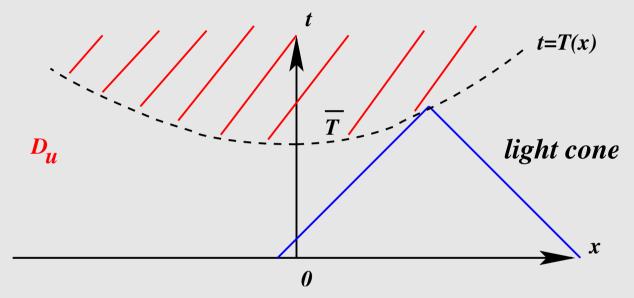
**Lemma 1 (Monotonicity (Antonini-Merle))** For all  $s_1$  and  $s_2$ :

$$E(w(s_2)) - E(w(s_1)) = -\frac{4}{p-1} \int_{s_1}^{s_2} \int_{-1}^{1} (\partial_s w)^2 (1-|y|^2)^{\frac{2}{p-1}-1} dy ds.$$

**Lemma 2 (A blow-up criterion)** Consider a solution W such that  $E(W(s_0)) < 0$  for some  $s_0 \in \mathbb{R}$ , then W blows up in finite time  $S > s_0$ .

# The blow-up rate

We look for a *local* blow-up rate near the singular surface (i.e. near every local blow-up time,  $t \to T(x_0)$ ), in  $H^1 \times L^2$  of the section of the light cone.



**Hint**: Is the rate given by the associated ODE  $v'' = v^p$ ?

# An upper bound on the blow-up rate in selfsimilar variables

Th. For all  $x_0 \in \mathbb{R}$  and  $s \ge -\log T(x_0) + 1$ ,

$$\int_{-1}^{1} \left( \frac{1}{2} (\partial_s w)^2 + \frac{1}{2} (\partial_y w)^2 (1 - |y|^2) + \frac{(p+1)}{(p-1)^2} w^2 + \frac{1}{p+1} |w|^{p+1} \right) \rho dy \le K$$

where the constant K depends only on p and an upper bound on  $T(x_0)$ ,  $1/T(x_0)$  and  $||(u_0, u_1)||$ .

#### Getting rid of the weights

Reducing (-1,1) to  $(-\frac{1}{2},\frac{1}{2})$ , we get:

Cor. For all  $x_0 \in \mathbb{R}$  and  $s \ge -\log T(x_0) + 1$ ,

$$\int_{-\frac{1}{2}}^{\frac{1}{2}} \left( (\partial_s w)^2 + (\partial_y w)^2 + w^2 + |w|^{p+1} \right) dy \le K.$$

# **Upper bound in the original** u(x, t) **variables**

**Th. sup.** *For all*  $x_0 \in \mathbb{R}$  *and*  $t \in [\frac{3}{4}T(x_0), T(x_0))$ :

$$(T(x_0)-t)^{\frac{2}{p-1}}\frac{\|u(t)\|_{L^2(B(x_0,\frac{T(x_0)-t}{2}))}}{(T(x_0)-t)^{1/2}}$$

$$+(T(x_0)-t)^{\frac{2}{p-1}+1}\left(\frac{\|u_t(t)\|_{L^2(B(x_0,\frac{T(x_0)-t}{2}))}}{(T(x_0)-t)^{1/2}}+\frac{\|\partial_{x}u(t)\|_{L^2(B(x_0,\frac{T(x_0)-t}{2}))}}{(T(x_0)-t)^{1/2}}\right)\leq K.$$

**Rk.** We have a lower bound of the same size when  $x_0$  is non characteristic (see Part 4 on profiles near a non characteristic point).

# Idea of the proof of the upper bound

- Selfsimilar transformation and existence of a Lyapunov functional
- Interpolation to gain regularity
- Gagliardo-Nirenberg estimates.

# Part 4: Asymptotic behavior at a non characteristic point

Take  $x_0 \in \mathbb{R}$  non characteristic. Using a covering argument for x near  $x_0$ , we obtain that  $\|(w_{x_0}(s), \partial_s w_{x_0}(s))\|_{H^1 \times L^2(-1,1)}$  is bounded.

**Question**: Does  $w_{x_0}(y,s)$  have a limit or not, as  $s \to \infty$  (that is as  $t \to T(x_0)$ ).

Remark: In the context of Hamiltonian systems, this question is delicate, and there is no natural reason for such a convergence, since the wave equation is time reversible.

See for similar difficulty and approach, results for

- the critical KdV (Martel and Merle),
- NLS (Merle and Raphaël).

# Stationary solutions.

We look for solutions of

$$\frac{1}{\rho} \left( \rho (1 - y^2) w' \right)' - \frac{2(p+1)}{(p-1)^2} w + |w|^{p-1} w = 0.$$

We work in  $\mathcal{H}_0$ , the (stationary energy space) defined by

$$\mathcal{H}_0 = \{ r \in H^1_{loc}(-1,1) \mid ||r||^2_{\mathcal{H}_0} \equiv \int_{-1}^1 \left( r'^2(1-y^2) + r^2 \right) \rho dy < +\infty \}.$$

**Prop.** Consider a stationary solution in  $\mathcal{H}_0$ . Then, either  $w \equiv 0$  or there exist  $d \in (-1,1)$  and  $e = \pm 1$  such that  $w(y) = e\kappa(d,y)$  where

$$\forall (d,y) \in (-1,1)^2, \ \kappa(d,y) = \kappa_0 \frac{(1-d^2)^{\frac{1}{p-1}}}{(1+dy)^{\frac{2}{p-1}}} \text{ and } \kappa_0 = \left(\frac{2(p+1)}{(p-1)^2}\right)^{\frac{1}{p-1}}.$$

**Remark**: We have 3 connected components.  $E(0) = 0 < E(\pm \kappa(d)) = E(\kappa_0)$ .

# Blow-up profile near a non characteristic point

**Th.** There exist  $C_0 > 0$  and  $\mu_0 > 0$  such that if  $x_0$  is **non characteristic**, then there exist  $d(x_0) \in (-1,1)$ ,  $e(x_0) = \pm 1$  and  $s^*(x_0) \ge -\log T(x_0)$  such that: (i) For all  $s \ge s^*(x_0)$ ,

$$\left\| \left( \begin{array}{c} w_{x_0}(s) \\ \partial_s w_{x_0}(s) \end{array} \right) - e(x_0) \left( \begin{array}{c} \kappa(d(x_0), \cdot) \\ 0 \end{array} \right) \right\|_{\mathcal{H}} \le C_0 e^{-\mu_0(s-s^*)}$$

and  $E(w_{x_0}(s) \to E(\kappa_0)$  where the energy space

$$\mathcal{H} = \left\{ q \in H_{loc}^{1} \times L_{loc}^{2}(-1,1) \mid \|q\|_{\mathcal{H}}^{2} \equiv \int_{-1}^{1} \left( q_{1}^{2} + \left( q_{1}^{\prime} \right)^{2} (1 - y^{2}) + q_{2}^{2} \right) \rho dy < +\infty \right\}.$$

(ii) 
$$d(x_0) = T'(x_0)$$
.

**Rk.** We have exp. fast convergence (hence,  $C^{1,\mu_0}$  regularity of  $\mathcal{R}$ , see Nouaili).

**Rk.** 
$$||w_{x_0}(y,s) - e(x_0)\kappa(d(x_0),y)||_{L^{\infty}(-1,1)} \to 0.$$

**Rk.** The parameter of the profile  $d(x_0)$  has a geometrical interpretation  $(T'(x_0))$ .

# Difficulties of the proof of convergence

- ► The set of non zero stationary solutions is made up of non isolated solutions (one parameter family):
  - → we need a modulation technique.
- The linearized operator around a non zero stationary solution is non self-adjoint:
  - we need to use dispersive properties coming from the Lyapunov functional to control the negative part of the spectrum.

# Part 5: Asymptotic behavior at a characteristic point

Th. If  $x_0 \in \mathbb{R}$  is characteristic, then, there exist  $k(x_0) \geq 2$ ,  $e(x_0) = \pm 1$  and continuous  $d_i(s) = -\tanh \zeta_i(s)$  for i = 1, ..., k such that: (i)

$$\left\| \left( \begin{array}{c} w_{x_0}(s) \\ \partial_s w_{x_0}(s) \end{array} \right) - e(x_0) \sum_{i=1}^{k(x_0)} (-1)^i \left( \begin{array}{c} \kappa(d_i(s), \cdot) \\ 0 \end{array} \right) \right\|_{\mathcal{H}} \to 0 \text{ as } s \to \infty,$$

(ii) Introducing

$$\bar{w}_{x_0}(\xi,s) = (1-y^2)^{\frac{1}{p-1}} w_{x_0}(y,s) \text{ with } y = \tanh \xi \text{ and } \zeta_i(x_0) = -\tanh^{-1} d_i(s),$$

we get

$$\|\bar{w}_{x_0}(\xi,s) - e(x_0)\sum_{i=1}^{k(x_0)} (-1)^i \cosh^{-\frac{2}{p-1}}(\xi - \zeta_i(s))\|_{H^1 \cap L^\infty(\mathbb{R})} \to 0 \text{ as } s \to \infty,$$

# Part 5: Asymptotic behavior at a characteristic point (cont.)

(iii) For all  $i = 1, ..., k(x_0)$  and s large enough,

$$\left(i - \frac{(k(x_0) + 1)}{2}\right) \frac{(p - 1)}{2} \log s - C_0 \le \zeta_i(s) \le \left(i - \frac{(k(x_0) + 1)}{2}\right) \frac{(p - 1)}{2} \log s + C_0.$$

(iv) 
$$E(w_{x_0}(s)) \to k(x_0)E(\kappa_0)$$
 as  $s \to \infty$ .

#### Rk.

- As  $s \to \infty$ ,  $w_{x_0}$  becomes like a **decoupled** sum of *equidistant* stationary solutions ("solitons"), with *alternate* signs.
- In the  $\xi$  variable, half of the solitons go to  $-\infty$ , and the other half to  $+\infty$ .
- The main difficulty in the proof is to prove that  $k(x_0) \ge 2$  (the case  $k(x_0) = 0$  is harder to eliminate).
- The  $\zeta_i(s)$  satisfy a Toda system:

$$\frac{1}{c_1}\zeta_i'(s) = e^{-\frac{2}{p-1}(\zeta_i - \zeta_{i-1})} - e^{-\frac{2}{p-1}(\zeta_{i+1} - \zeta_i)} + R_i \text{ with } R_i = o\left(\sum_{j=1}^{k-1} e^{-\frac{2}{p-1}(\zeta_{j+1} - \zeta_j)}\right) \text{ as } s \to \infty.$$

# The energy behavior

**Defining** 

$$k(x_0) = 1$$
 when  $x_0 \in \mathcal{R}$ ,

we get the following:

Cor.

(i) For all  $x_0 \in \mathbb{R}$  and  $s \ge -\log T(x_0)$ , we have

$$E(w_{x_0}(s)) \ge k(x_0)E(\kappa_0).$$

(ii) (An energy criterion for non characteristic points) *If for some*  $x_0 \in \mathbb{R}$  *and*  $s_0 \ge -\log T(x_0)$ , we have

$$E(w_{x_0}(s_0)) < 2E(\kappa_0),$$

then  $x_0 \in \mathcal{R}$ .