# ON THE DEPENDENCE OF THE BLOW-UP TIME WITH RESPECT TO THE INITIAL DATA IN A SEMILINEAR PARABOLIC PROBLEM

PABLO GROISMAN, JULIO D. ROSSI, AND HATEM ZAAG

ABSTRACT. We find a bound for the modulus of continuity of the blow-up time for the semilinear parabolic problem  $u_t = \Delta u + |u|^{p-1}u$ , with respect to the initial data.

#### Introduction.

In this paper we study the dependence with respect to the initial data of the maximal time of existence for solutions of the following problem

(1.1) 
$$\begin{cases} u_t = \Delta u + |u|^{p-1}u, & (x,t) \in \Omega \times (0,T), \\ u(x,0) = u_0(x), & x \in \Omega, \end{cases}$$

where p is superlinear and subcritical, that is, p>1 and p(N-2)< N+2. Here  $\Omega$  is a bounded convex smooth domain in  $\mathbb{R}^N$  or the whole  $\mathbb{R}^N$ . In case we are dealing with a bounded domain we assume Dirichlet boundary conditions, that is u(x,t)=0, for  $(x,t)\in\partial\Omega\times(0,T)$ . Also we assume that the initial data,  $u_0(x)$ , is regular in order to guarantee existence, uniqueness and regularity of the solution.

A remarkable fact is that the solution of parabolic problems may develop singularities in finite time, no matter how smooth the initial data are. It is well known that for many differential equations or systems the solutions can become unbounded in finite time (a phenomena that is known as blow-up). Typical examples where this happens are problems involving nonlinear reaction terms in the equation like (1.1), see [18], [19], [21] and the references therein.

Key words and phrases. Parabolic equations, blow-up time, continuity. 2000 Mathematics Subject Classification. 35K55, 35B40.

Supported by ANPCyT PICT No. 03-00000-00137, CONICET and Fundación Antorchas. (Argentina).

For our problem, if the maximal solution is defined on a finite time interval, [0,T) with  $T<+\infty$ , then

$$\lim_{t \nearrow T} \|u(\cdot, t)\|_{L^{\infty}} = +\infty.$$

We say that u blows up at time T.

The existence of blowing up solutions for (1.1) has been proved by several authors, see [1], [5], [7], [14], etc. The study of the blow-up behaviour for blow-up solutions of (1.1) has attracted a considerable attention in recent years, see for example, [2], [3], [6], [11], [8], [12], [13], [15], [18], [22]. In particular, Giga and Kohn [10] and recently, Giga et al. [9] have proved the following upper bound for any solution u of (1.1) that blows up at time T:

$$(1.2) \forall t \in [0, T), ||u(t)||_{L^{\infty}} \le Cv(t - T) = C\kappa(T - t)^{-\frac{1}{p-1}}$$

for some C > 0, where  $\kappa = (p-1)^{-\frac{1}{p-1}}$  and v is the solution of the associated ODE  $v' = v^p$ ,  $v(0) = +\infty$ . Let us mention that an easy use of the maximum principle shows that, unlike the constant C in (1.2), the power of (T-t) is optimal, in the sense that

$$(1.3) \forall t \in [0, T), ||u(t)||_{L^{\infty}} \ge v(t - T) = \kappa (T - t)^{-\frac{1}{p-1}}.$$

Later, the following Liouville Theorem has been proved by Merle and Zaag in [16] and [18]:

Consider U a solution of (1.1) defined for all  $(x,t) \in \mathbb{R}^N \times (-\infty,T)$ such that for all  $(x,t) \in \mathbb{R}^N \times (-\infty,T)$ ,  $|U(x,t)| \leq C(T-t)^{-\frac{1}{p-1}}$ . Then, either  $U \equiv 0$  or  $U(x,t) = [(p-1)(T^*-t)]^{-\frac{1}{p-1}}$  for some  $T^* \geq T$ .

This result is the key to find the best constant C in (1.2), as will be presented later.

From now on we assume that we are dealing with an initial datum  $u_0$  which produces a solution u that blows up at time T with

$$||u_0||_{L^{\infty}} \leq M_0$$
 and  $T \leq T_0$ .

For every h(x) with  $L^{\infty}$ -norm small enough, the solution  $u_h$  of problem (1.1) that has initial datum  $u_0(x) + h(x)$  also blows up in finite time, that we call  $T_h$ . Indeed, it is known that the blow-up time is continuous with respect to the initial data  $u_0$ ; see [20], [2], [15] if  $\Omega$  is bounded; if  $\Omega = \mathbb{R}^N$ , this was done in [4] under (1.2). Hence,  $T_h \to T$  as  $||h||_{L^{\infty}} \to 0$ .

Our interest here is to provide a bound for  $|T-T_h|$  in terms of  $||h||_{L^{\infty}}$ . To this end, we obviously need estimates on  $u_h$ , uniform with respect to h. It is to be noticed that before the proof of the Liouville Theorem

of [18], no such estimates were available in the literature. We should emphasize that the only exception to this statement is the fact that the argument of [10] and [9] in the proof of (1.2) does hold uniformly with respect to initial data. Indeed, a simple use of the continuity of  $u_h(t_0)$  for a fixed  $t_0$ , with respect to initial data, shows that Giga and Kohn actually proved that

There exists  $C = C(M_0, T_0) > 0$  such that for all h small enough,  $\forall (x, t) \in \Omega \times [0, T_h)$ ,

$$(1.4) |u_h(x,t)| \le Cv(t-T_h) = C\kappa(T_h-t)^{-\frac{1}{p-1}}.$$

See Theorem 2 in [18] and its proof for a simple sketch of this argument. Using the Liouville Theorem, the estimate (1.4) has been refined in [17] and [18] uniformly with respect to h:

Given  $\varepsilon > 0$ , there exists  $\tau = \tau(M_0, T_0, \varepsilon)$  such that for every h with  $||h||_{L^{\infty}}$  small, the solution of problem (1.1) verifies

$$(1.5) |u_h(x,t)| \le \kappa (T_h - t)^{-\frac{1}{p-1}} + \left(\frac{N\kappa}{2p} + \varepsilon\right) \frac{(T_h - t)^{-\frac{1}{p-1}}}{|\ln(T_h - t)|},$$

for every  $(x,t) \in \Omega \times [T_h - \tau, T_h)$ .

We will call this fact hypothesis (H). Note that all the results of [18] hold under (1.2) which is true for all subcritical p with no sign condition. Hence (H) is valid for any subcritical p.

Now we state our result,

**Theorem 1.1.** Assume that hypothesis (H) holds. Let T and  $T_h$  be the blow-up times for u and  $u_h$ , two solutions of problem (1.1) with initial data  $u_0$  and  $u_0 + h$  respectively, then for every  $\varepsilon > 0$  there exist a positive constant  $C = C(M_0, T_0, \varepsilon)$  such that for every h smaller than some  $\eta(M_0, T_0, \varepsilon) > 0$ ,

$$(1.6) |T - T_h| \le C ||h||_{L^{\infty}} |\ln(||h||_{L^{\infty}})|^{\frac{N+2}{2} + \varepsilon}.$$

**Remark**: We suspect that an estimate of the form  $|T - T_h| \le C||h||_{L^{\infty}}$  is valid, as has been proved by Herrero and Velázquez [13] in one dimension. However, we cannot get rid of the logarithmic term in (1.6).

**Remark**: We want to remark that the restriction on p, namely  $1 if <math>N \ge 3$ , is not technical. Indeed, if it does not hold, then the blow-up time is not even continuous as a function of the initial data, see [8].

## Proof of Theorem 1.1

First, let us give an idea of the proof. Let u and  $u_h$  be the solutions of (1.1) with initial data  $u_0$  and  $u_0 + h$  respectively. When  $T \neq T_h$ , the function  $\|e(\cdot,t)\|_{L^{\infty}} = \|(u_h - u)(\cdot,t)\|_{L^{\infty}}$  will grow from its initial value  $\|h\|_{L^{\infty}}$  to  $+\infty$ , as  $t \to \min(T, T_h)$ . If  $t_0$  is the first time where  $\|e(t)\|_{L^{\infty}}$  reaches a given size, we get a bound on  $\|e(t)\|_{L^{\infty}}$  that allows us to control  $T - t_0$  and  $T_h - t_0$  in terms of  $\|h\|_{L^{\infty}}$ . We then use a triangular inequality to conclude.

The hypothesis (H) is crucial in getting a sharp estimate on  $||e(t)||_{L^{\infty}}$ . Given  $\varepsilon > 0$ , it is easy to see from (1.5) that (H) holds uniformly for u (resp. for  $u_h$ ), for all  $t \in [T - \tau, T)$  (resp.  $t \in [T - \tau, T_h)$ ), where  $\tau = \tau(M_0, T_0, \epsilon)$ , whenever  $||h||_{L^{\infty}}$  is small enough. Since  $T - \frac{\tau}{2}$  is independent of h, we prove in the following lemma that the size of e at  $t = T - \frac{\tau}{2}$  is comparable to its size at t = 0. Therefore, we can make a shift in time to simplify the notation and assume that (H) holds for u (resp. for  $u_h$ ) for all  $t \in [0, T)$  (resp.  $t \in [0, T_h)$ ). We claim the following lemma:

**Lemma 1.1.** Given  $\varepsilon > 0$  there exists a constant  $C = C(M_0, T_0, \varepsilon) > 0$  such that for every h with  $||h||_{L^{\infty}}$  small enough, it holds that  $u_h$ , the solution with initial datum  $u_0 + h$ , is defined on  $[0, T - \frac{\tau}{2}]$  where  $\tau = \tau(M_0, T_0, \epsilon)$ , and moreover

(1.7) 
$$||u - u_h||_{L^{\infty}} (T - \frac{\tau}{2}) \le C||h||_{L^{\infty}}.$$

**Proof:** According to (1.4), there exists a constant  $M(M_0, T_0, \epsilon)$  such that

$$M \ge ||u||_{L^{\infty}(\Omega \times [0, T - \frac{\tau}{2}])}.$$

Consider a uniformly Lipschitz function f(u) that coincides with the power nonlinearity  $|u|^{p-1}u$  for values of u with  $|u| \leq M+1$ . Now let v and  $v_h$  the solutions of a problem like (1.1) with the source term  $|u|^{p-1}u$  replaced by f(u) and initial datum  $u_0$  and  $u_0 + h$  respectively. As f(u) is uniformly Lipschitz we have that both v and  $v_h$  are global and moreover there exists a constant C such that

$$||v - v_h||_{L^{\infty}} (T - \frac{\tau}{2}) \le C ||h||_{L^{\infty}}.$$

Now we only have to observe that if  $||h||_{L^{\infty}}$  is small enough we have that both v and  $v_h$  verify that

$$||v||_{L^{\infty}(\Omega \times [0, T - \frac{\tau}{2}])}, ||v_h||_{L^{\infty}(\Omega \times [0, T - \frac{\tau}{2}])} \le M + 1.$$

Hence, by uniqueness,  $v(T - \frac{\tau}{2})$  and  $v_h((T - \frac{\tau}{2})$  coincides with  $u(T - \frac{\tau}{2})$  and  $u_h((T - \frac{\tau}{2}))$  respectively and the result follows.

Now we can prove Theorem 1.1,

**Proof of Theorem 1.1:** First of all we observe that for  $u_0$  fixed, the set of functions h such that  $u_h$  blows up in finite time is open. So if u blows up, the same occurs with  $u_h$  for h small enough. From Lemma 1.1, taking as initial datum  $u(T - \frac{\tau}{2})$  and  $u_h((T - \frac{\tau}{2}))$  if necessary, we can assume that hypothesis (H) is valid for all t in [0, T) and  $[0, T_h)$  respectively.

Let  $e(x,t) = u_h(x,t) - u(x,t)$ . We have that for every  $(x,t) \in \Omega \times [0, \min(T, T_h))$ ,

(1.8) 
$$\begin{cases} e_t = \Delta e + |u_h|^{p-1} u_h - |u|^{p-1} u = \Delta e + p|\xi(x,t)|^{p-1} e, \\ \|e(x,0)\|_{L^{\infty}} \le C \|h\|_{L^{\infty}}, \end{cases}$$

where  $\xi(x,t)$  lies between u(x,t) and  $u_h(x,t)$ . Given  $\varepsilon > 0$ , we can define for  $||h||_{L^{\infty}}$  small enough, a time  $t_0 \in (0,T]$  as the maximal time such that

(1.9) 
$$||e(\cdot,t)||_{L^{\infty}} \le \varepsilon \frac{(T-t)^{-\frac{1}{p-1}}}{|\ln(T-t)|} \quad \text{for all } t \in [0,t_0).$$

For times  $t \in [0, t_0)$  we have that

$$|\xi(x,t)| \le |u(x,t)| + \varepsilon \frac{(T-t)^{-\frac{1}{p-1}}}{|\ln(T-t)|}$$

and as we are assuming (H) we get that

$$|\xi(x,t)| \le \kappa (T-t)^{-\frac{1}{p-1}} + \left(\frac{N\kappa}{2p} + 2\varepsilon\right) \frac{(T-t)^{-\frac{1}{p-1}}}{|\ln(T-t)|},$$

hence,

$$p|\xi(x,t)|^{p-1} \le \frac{p}{p-1}(T-t)^{-1} + \left(\frac{N}{2} + C\varepsilon\right)\frac{(T-t)^{-1}}{|\ln(T-t)|} \equiv \alpha(t).$$

Now, let  $\overline{e}(t)$  be the solution of the ODE

$$\overline{e}'(t) = \alpha(t)\overline{e}(t)$$

with initial datum  $\overline{e}(0) = C||h||_{L^{\infty}}$ . Integrating we obtain

$$\overline{e}(t) = C \|h\|_{L^{\infty}} (T-t)^{-\frac{p}{p-1}} |\ln(T-t)|^{\frac{N}{2} + C\varepsilon}$$

and, since e is a solution of (1.8), by a comparison argument we have for every  $t \in [0, t_0)$ ,

$$(1.10) ||e(\cdot,t)||_{L^{\infty}} \le \overline{e}(t) = C||h||_{L^{\infty}}(T-t)^{-\frac{p}{p-1}}|\ln(T-t)|^{\frac{N}{2}+C\varepsilon}.$$

We want to use this to obtain a bound on  $|T - t_0|$  in terms of  $||h||_{L^{\infty}}$ . To this end let us introduce for  $||h||_{L^{\infty}}$  small enough, the time  $t_1 = t_1(h) \leq T$  as the maximal time such that

$$C\|h\|_{L^{\infty}}(T-t)^{-\frac{p}{p-1}}|\ln(T-t)|^{\frac{N}{2}+C\varepsilon} \le \varepsilon \frac{(T-t)^{-\frac{1}{p-1}}}{|\ln(T-t)|}$$

for all  $t \in [0, t_1)$ .

It is clear from (1.10) that  $t_1 \leq t_0$ , hence  $0 \leq T - t_0 \leq T - t_1$ . Note that

$$T - t_1(h) \sim \frac{C}{\varepsilon} ||h||_{L^{\infty}} |\ln ||h||_{L^{\infty}} |\frac{N+2}{2} + C\epsilon$$
 as  $||h||_{L^{\infty}} \to 0$ .

Therefore,

$$T - t_0 \le T - t_1 \le \frac{C}{\varepsilon} \|h\|_{L^{\infty}} |\ln \|h\|_{L^{\infty}}|^{\frac{N+2}{2} + C\varepsilon}.$$

Now we observe that (1.4) and the definition of e yield

$$|T_h - t_0| \le C \left( \|u_h(t_0, \cdot)\|_{L^{\infty}(\Omega)} \right)^{1-p} \le C \left( \|u(t_0, \cdot)\|_{L^{\infty}} + \|e(t_0, \cdot)\|_{L^{\infty}} \right)^{1-p}.$$

Using (1.3) and (1.9), we get

$$|T_h - t_0| \le C \left( \|u(t_0, \cdot)\|_{L^{\infty}(\Omega)} - \varepsilon \frac{(T - t_0)^{-\frac{1}{p-1}}}{\ln(T - t_0)} \right)^{1-p}$$

$$\leq C \left( \kappa (T - t_0)^{-\frac{1}{p-1}} - \varepsilon \frac{(T - t_0)^{-\frac{1}{p-1}}}{|\ln(T - t_0)|} \right)^{1-p} \leq C(T - t_0),$$

and so we obtain

$$|T - T_h| \le |T - t_0| + |T_h - t_0| \le$$

$$|T-t_0|+C|T-t_0| \leq \frac{C}{\varepsilon} ||h||_{L^{\infty}} \left|\ln(||h||_{L^{\infty}})\right|^{\frac{N+2}{2}+C\varepsilon}.$$

This completes the proof.

**Remark**: The above proof gives a constant  $C(M_0, T_0, \varepsilon)$  for (1.6) with  $C(M_0, T_0, \varepsilon) \to +\infty$  as  $\varepsilon \to 0$ .

## Acknowledgments

This work was started during a visit of JDR to Ecole Normale Supérieure and Univ. Versailles, he wants to thanks for the hospitality found in those places.

#### References

- [1] J. Ball. Remarks on blow-up and nonexistence theorems for nonlinear evolution equations. Quart. J. Math. Oxford, Vol. 28, (1977), 473-486.
- [2] P. Baras and L. Cohen. Complete blow-up after  $T_{max}$  for the solution of a semilinear heat equation. J. Funct. Anal. 71, (1987), 142-174.
- [3] X. Y. Chen and H. Matano. Convergence, asymptotic periodicity and finite point blow up in one-dimensional semilinear heat equations. J. Differential Equations, Vol. 78, (1989), 160-190.
- [4] C. Fermanian Kammerer, F. Merle, and H. Zaag. Stability of the blow-up profile of non-linear heat equations from the dynamical system point of view. Math. Ann., 317, (2000), 195-237.
- [5] A. Friedman. Remarks on nonlinear diffusion equations. In Proc. Sympos. Appl. Math., Vol. XVII, Amer. Math. Soc., Providence, R.I., (1965), 2-23.
- [6] A. Friedman and J. B. Mc Leod. Blow up of positive solutions of semilinear heat equations. Indiana Univ. Math. J., Vol. 34, (1985), 425-447.
- [7] H. Fujita. On the blowing up of solutions of the Cauchy problem for  $u_t = \Delta u + u^{1+\alpha}$ . J. Fac. Sci. Univ. Tokyo Sect. I, (1966), 109-124.
- [8] V. Galaktionov and J. L. Vazquez. Continuation of blow-up solutions of non-linear heat equations in several space dimensions. Comm. Pure Appl. Math. 50, (1997), 1-67.
- [9] Y. Giga, Matsai and Sasayama. Personal communication, (2001).
- [10] Y. Giga and R. V. Kohn. Characterizing blowup using similarity variables. Indiana Univ. Math. J. Vol. 36 (1987), 1-40.
- [11] Y. Giga and R. V. Kohn. Nondegeneracy of blow up for semilinear heat equations. Comm. Pure Appl. Math. Vol. 42, (1989), 845-884.
- [12] M. A. Herrero and J. J. L. Velazquez. Flat blow up in one-dimensional, semi-linear parabolic problems, Differential Integral Equations. Vol. 5(5), (1992), 973-997.
- [13] M. A. Herrero and J. J. L. Velazquez. Generic behaviour of one-dimensional blow up patterns. Ann. Scuola Norm. Sup. di Pisa, Vol. XIX (3), (1992), 381-950
- [14] H. A. Levine. Some nonexistence and instability theorems for solutions of formally parabolic equations of the form  $Pu_t = -Au + F(u)$ . Arch. Rational Mech. Anal., Vol. 51, (1973), 371-386.
- [15] F. Merle. Solution of a nonlinear heat equation with arbitrarily given blow-up points. Comm. Pure Appl. Math. Vol. XLV, (1992), 263-300.
- [16] F. Merle and H. Zaag. Optimal estimates for blowup rate and behavior for nonlinear heat equations. Comm. Pure Appl. Math., Vol. 51, (1998), 139-196.
- [17] F. Merle and H. Zaag. Refined uniform estimates at blow-up and applications for nonlinear heat equations. Geom. Funct. Anal., Vol. 8, (1998), 1043-1085.
- [18] F. Merle and H. Zaag. A Liouville theorem for vector-valued nonlinear heat equations and applications. Math. Ann. Vol. 316, (2000), 103-137.
- [19] C. V. Pao. Nonlinear parabolic and elliptic equations. Plenum Press 1992.
- [20] P. Quittner. Continuity of the blow-up time and a priori bounds for solutions in superlinear parabolic problems. Houston J. Math, to appear.
- [21] A. Samarskii, V. A. Galaktionov, S. P. Kurdyumov and A. P. Mikhailov. Blow-up in quasilinear parabolic equations. Walter de Gruyter, Berlin, (1995).

[22] H. Zaag. One dimensional behavior of singular N dimensional solutions of semilinear heat equations. Comm. Math. Phys. Vol. 225, (2002), 523-549.

Universidad de San Andrés, Vito Dumas 284 (1644) Victoria, Pcia. de Buenos Aires, Argentina

E-mail address: pgroisman@udesa.edu.ar

DEPARTAMENTO DE MATEMÁTICA, FCEYN., UBA, (1428) BUENOS AIRES, ARGENTINA.

 $E ext{-}mail\ address: jrossi@dm.uba.ar}$ 

CNRS UMR 8553, DÉPARTMENT DE MATHÉMATIQUES ET APPLICATIONS, ECOLE NORMALE SUPÉRIEURE, 45 RUE D'ULM, F-75230, PARIS CEDEX 05, FRANCE.

 $E ext{-}mail\ address: {\tt Hatem.Zaag@ens.fr}$