

Low temperature
asymptotics for
Quasi Stationary
Distributions in a
bounded domain

Francis Nier,
IRMAR, Univ.
Rennes 1
Joint work with
T. Lelièvre

Quasi-Stationary
Distributions

Parallel Replica
and
Hyperdynamics

Witten Laplacians

Results

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Sept. 9, 2013

Outline

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Definition of QSD

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Results

Consider the SDE in \mathbb{R}^d

$$dX_t = -\partial_x f(X_t)dt + \sqrt{2\beta^{-1}}dW_t \quad , \quad \beta = \frac{1}{kT} .$$

with $f \in \mathcal{C}^\infty(\mathbb{R}^d)$.

Ω domain of \mathbb{R}^d . A trajectory starting with $X_0 = x \in \Omega$ will be denoted $X_{x,t}$ and the exit time τ_x equals

$$\tau_x = \min \{ t \in [0, +\infty), X_{x,t} \in \partial\Omega \} .$$

Definition:

A probability measure μ is a QSD in Ω if for

$$\forall \varphi \in \mathcal{C}^\infty(\bar{\Omega}), \forall t > 0, \mathbb{E}(\varphi(X_{x,t}) | t < \tau_x) = \int_{\Omega} \varphi(x) d\mu(x)$$

when the law of $X_{x,0}$ is μ .

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Use of QSD

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QSD can be defined with more general stochastic processes. They model the stochastic evolution of a population in a domain with possible individual escape or end. Applications (after P.K. Pollett): Biology and ecology, Chemical kinetics, Epidemics, Genetics, Reliability, Telecommunication, Medicine, Neurosciences.

They provide the theoretical framework for analysing stochastic algorithms **Parallel Replica** and Fleming-Viot algorithms.

On his website

<http://www.maths.uq.edu.au/~pkp/papers/qsd/qsd.pdf>

P.K. Pollett lists 429 references, including 5 textbooks, from 1947 (Yaglom) with premises in 1931 (S. Wright), until 2012.

The name Quasi-Stationary Distribution was given in Barlett's book (1960).

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In our case with a regular bounded domain Ω , there exists a unique QSD, $\mu = Z_{v1\Omega}^{-1} v(x) 1_{\Omega}(x) e^{-\beta f(x)} dx$, where v is the non-negative eigenvector of the Dirichlet realization L^D of $L = -\partial_x f(x) \cdot \partial_x + \beta^{-1} \Delta_x$.

For any regular φ and any $x \in \Omega$, $\mathbb{E}(\varphi(X_{x,t}) | t < \tau_x)$ converges to $\int_{\Omega_+} \varphi(x) d\mu(x)$ as $t \rightarrow +\infty$.

When $X_{x,0}$ is distributed according the μ , the exit times τ and exit point X_{τ} are independent random variables with

- τ has the exponential law with coefficient λ_1 , the first eigenvalue of L^D
- X_{τ} has the law given by

$$\mathbb{E}(\varphi(X_{\tau})) = -\frac{\int_{\partial\Omega} \varphi \partial_n (v e^{\beta f}) d\sigma}{\beta \lambda_1 \int_{\Omega} v e^{-\beta f} dx} = \frac{\int_{\partial\Omega} \varphi \partial_n (v e^{\beta f}) d\sigma}{\int_{\partial\Omega} \partial_n (v e^{\beta f}) d\sigma}.$$

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Aim

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In a complex system (i.e. with a large number of degrees of freedom), one may identify a finite set of metastable states S . One wants to find the effective Markov process on S from suitable simulations of the whole system.

For a metastable domain $\Omega \in S$, the quasistationary distribution and the laws of τ and X_τ give respectively the exit process to the exterior of Ω and the reaction paths and proportions to other states of S .

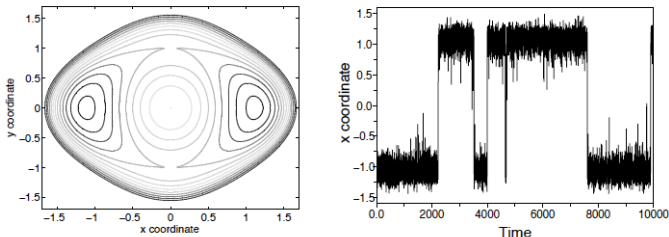


Figure: An example with $\#S = 2$.

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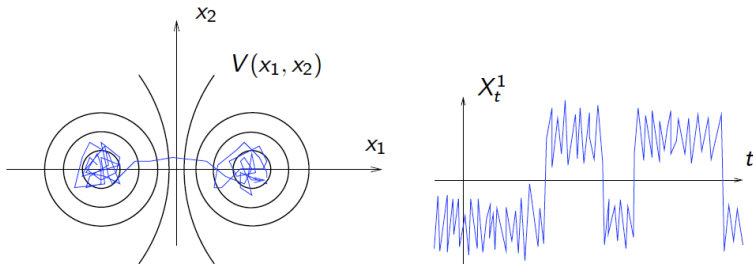


Figure: Here the metastable states are identified with $x_1 > 0$ and $x_1 < 0$

Parallel Replica

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This method was introduced by A. Voter (1998).

One step consists in sampling as follows the exit time and exit position of a QSD in Ω : On parallel processors a large number K of independent particles $X_{x_k, t}$, $k = 1, \dots, K$ are simulated and the simulation is stopped as soon as one particle hits the boundary (time T). It can be proved that the distribution of $(X_{x_k, T})_{1 \leq k \leq K}$ approximates the QSD.

REF Le Bris–Lelièvre–Luskin–Perez (2012).

Hyperdynamics

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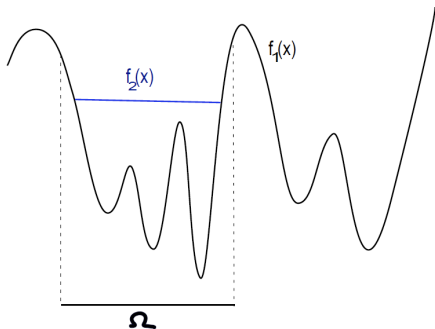


Figure: The potential is increased from f_1 to f_2 inside Ω .

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Guess : The exit time law coefficient is changed by multiplying with the “boost” factor

$$\frac{\lambda_1(f_2)}{\lambda_1(f_1)} \simeq \frac{\int_{\Omega} e^{-\beta f_1(x)} dx}{\int_{\Omega} e^{-\beta f_2(x)} dx} = \frac{\int_{\Omega} e^{-\beta(f_1(x)-f_2(x))} e^{-\beta f_2(x)} dx}{\int_{\Omega} e^{-\beta f_2(x)} dx}$$

The exit point law is not changed

$$\frac{\partial_n(v_2 e^{\beta f_2}) d\sigma}{\int_{\partial\Omega} \partial_n(v_2 e^{\beta f_2}) d\sigma} = \frac{\partial_n(v_1 e^{\beta f_1}) d\sigma}{\int_{\partial\Omega} \partial_n(v_1 e^{\beta f_1}) d\sigma}$$

Question : Can we prove those statements in the small temperature regime $\beta \rightarrow +\infty$?

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Fokker-Planck operators and Witten Laplacians

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Remember

$$L = (\partial_x f(x) - \beta^{-1} \partial_x) \cdot \partial_x .$$

Conjugating with $e^{-\frac{\beta f}{2}}$ (or $e^{\frac{\beta f}{2}}$) leads to

$$\frac{4}{\beta} e^{-\frac{\beta f}{2}} L e^{\frac{\beta f}{2}} = \Delta_{f,h}^{(0)}$$

$$\Delta_{f,h}^{(0)} = -h^2 \Delta_x + |\partial_x f(x)|^2 - h \Delta f(x), \quad h = \frac{2}{\beta},$$

to be studied in $L^2(\Omega, dx)$. This is the semiclassical Witten Laplacian acting on functions.

More generally

$$\Delta_{f,h} = (d_{f,h} + d_{f,h}^*)^2 = d_{f,h}^* d_{f,h} + d_{f,h} d_{f,h}^*$$

$$\text{with } d_{f,h} = e^{-\frac{f}{h}} (hd) e^{\frac{f}{h}} \quad , \quad d_{f,h}^* = e^{\frac{f}{h}} (hd^*) e^{-\frac{f}{h}} .$$

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Witten Laplacians on compact manifolds

No boundary. f is a Morse function on Ω .

In his famous article of 81, E. Witten showed that the number of $\mathcal{O}(h^{3/2})$ eigenvalues of $\Delta_{f,h}^{(p)}$ is the number $m_p(\Omega)$ of critical points with index $p \rightarrow$ analytic approach to Morse inequalities. Continued in particular by Bismut and other collaborators for more general invariants (cf book by W. Zhang 01).

Result proved by Helffer–Sjöstrand in 85. The $\mathcal{O}(h^{3/2})$ eigenvalues are actually exponentially small.

After the works of Bovier–Eckhoff–Gayraud–Klein (04,04) using Freidlin–Wentzell and potential theory (capacities), accurate values of the exponentially small eigenvalues of $\Delta_{f,h}^{(0)}$ were given in Helffer–Klein–N. (04). Idea: Think of the eigenvalues of $\Delta_{f,h}^{(0)}|_{F_h} = d_{f,h}^* d_{f,h}|_{F_h}$ as the squared singular values of $d_{f,h}|_{F_h}$ (see also Bismut–Zhang 94), then Gaussian elimination (Le Peutrec 09) summarizes former tricky inductions.

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$\bar{\Omega}$ compact, f Morse, $f|_{\partial\Omega}$ Morse, $\|\nabla f\|$ does not vanish on $\partial\Omega$

In 95 Chang and Liu, proved that the number of $\mathcal{O}(h^{6/5})$ -eigenvalues of Dirichlet or Neumann Witten Laplacians was equal to the number of generalized critical points.

$$D(\Delta_{f,h}^N) = \left\{ \omega \in \bigwedge W^{2,2}(\Omega), \quad \mathbf{n}\omega = 0, \quad \mathbf{n}d_{f,h}\omega = 0 \right\}.$$

$$D(\Delta_{f,h}^D) = \left\{ \omega \in \bigwedge W^{2,2}(\Omega), \quad \mathbf{t}\omega = 0, \quad \mathbf{t}d_{f,h}^*\omega = 0 \right\}.$$

Idea: generalized saddle points Neumann (resp. Dirichlet) means $f = +\infty$ (resp $f = -\infty$) outside $\bar{\Omega}$.

Helffer-N. (06) computes accurately the exp. small eigenvalues of $\Delta_{f,h}^{D,(0)}$. Le Peutrec (10) did it for $\Delta_{f,h}^{N,(0)}$.

Inspired by Chang–Liu, Laudenbach (11) gave a topological proof of Morse inequalities for boundary manifolds.

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In 95 Chang and Liu, proved that the number of $\mathcal{O}(h^{6/5})$ -eigenvalues of Dirichlet or Neumann Witten Laplacians was equal to the number of generalized critical points.

$$D(\Delta_{f,h}^N) = \left\{ \omega \in \bigwedge W^{2,2}(\Omega), \quad \mathbf{n}\omega = 0, \quad \mathbf{n}d_{f,h}\omega = 0 \right\}.$$

$$D(\Delta_{f,h}^D) = \left\{ \omega \in \bigwedge W^{2,2}(\Omega), \quad \mathbf{t}\omega = 0, \quad \mathbf{t}d_{f,h}^*\omega = 0 \right\}.$$

Idea: generalized saddle points Neumann (resp. Dirichlet) means $f = +\infty$ (resp $f = -\infty$) outside $\bar{\Omega}$.

Helffer-N. (06) computes accurately the exp. small eigenvalues of $\Delta_{f,h}^{D,(0)}$. Le Peutrec (10) did it for $\Delta_{f,h}^{N,(0)}$.

Inspired by Chang–Liu, Laudenbach (11) gave a topological proof of Morse inequalities for boundary manifolds.

Boundary Witten Laplacians

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Morse function ?

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$$\frac{\lambda_1(f_2)}{\lambda_1(f_1)} \simeq \frac{\int_{\Omega} e^{-\beta f_1(x)} dx}{\int_{\Omega} e^{-\beta f_2(x)} dx}$$
$$\frac{\int_{\partial\Omega} \partial_n(v_2 e^{\beta f_2}) d\sigma}{\int_{\partial\Omega} \partial_n(v_2 e^{\beta f_2}) d\sigma} = \frac{\int_{\partial\Omega} \partial_n(v_1 e^{\beta f_1}) d\sigma}{\int_{\partial\Omega} \partial_n(v_1 e^{\beta f_1}) d\sigma}$$

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suggest that this should hold for non Morse functions.

Idea: Formulate all the assumptions in terms of restricted Witten Laplacians.

Framework

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$$\Omega_- \subset\subset \Omega_+ (= \Omega). \quad \nabla f \neq 0 \text{ in } \overline{\Omega_+} \setminus \Omega_- \quad , \quad \partial_n f|_{\Omega_\pm} > 0.$$

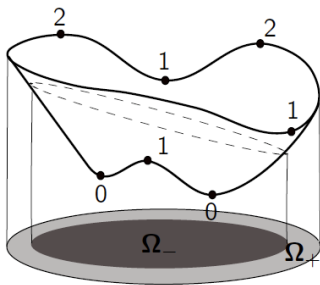


Figure: A typical 2D Morse example. The indices of the generalized critical points are labelled. .

Assumptions

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Results

There exist a constant $c_0 > 0$ and a function $\nu : (0, h_0) \rightarrow (0, +\infty)$ with

$$\forall \varepsilon > 0, \exists C_\varepsilon > 1, \quad \frac{1}{C_\varepsilon} e^{-\frac{\varepsilon}{h}} \leq \nu(h) \leq h,$$

and such that the following hypotheses are fulfilled.

Hyp 1 For $p = 0, 1$,

$$\begin{aligned} \# \left[\sigma(\Delta_{f,h}^{N,(p)}(\Omega_-)) \cap [0, \nu(h)] \right] &= m_p^N(\Omega_-) \\ \left[\sigma(\Delta_{f,h}^{N,(p)}(\Omega_-)) \cap [0, \nu(h)] \right] &\subset \left[0, e^{-\frac{c_0}{h}} \right] \end{aligned}$$

with $m_p^N(\Omega_-)$ independent of h . In a neighborhood \mathcal{V}_- of $\partial\Omega_-$ and for $\lambda(h) \in \sigma(\Delta_{f,h}^{N,(p)}(\Omega_-)) \cap (0, \nu(h)]$, the corresponding eigenvector satisfies

$$\|\psi(h)\|_{L^2(\mathcal{V}_-)} = \tilde{O}(\sqrt{\lambda(h)}).$$

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Hyp 2

$$\begin{aligned} \# \left[\sigma(\Delta_{f,h}^{D,(1)}(\Omega_+ \setminus \overline{\Omega_-})) \cap [0, \nu(h)] \right] &= m_1^D(\Omega_+ \setminus \overline{\Omega_-}), \\ \left[\sigma(\Delta_{f,h}^{D,(1)}(\Omega_+ \setminus \overline{\Omega_-})) \cap [0, \nu(h)] \right] &\subset \left[0, e^{-\frac{c_0}{h}} \right], \end{aligned}$$

with $m_1^D(\Omega_+ \setminus \overline{\Omega_-})$ independent of $h \in (0, h_0)$.

Hyp 3

$$\{x \in \Omega_-, \partial f(x) = 0\} \subset \left\{ x \in \Omega_-, f(x) < \min_{\partial\Omega_+} f - c_0 \right\}.$$

Notation

$$\kappa_f = \min_{\partial\Omega_+} f - \min_{\Omega_+} f.$$

Assumptions cont'

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Theorem 1

For $p = 0, 1$

$$m_0^D(\Omega_+) := \# \left[\sigma(\Delta_{f,h}^{D,(0)}(\Omega_+)) \cap [0, \nu(h)] \right] = m_0^N(\Omega_-),$$

$$m_1^D(\Omega_+) := \# \left[\sigma(\Delta_{f,h}^{D,(1)}(\Omega_+)) \cap [0, \nu(h)] \right] = m_1^N(\Omega_-) + m_1^D(\Omega_+ \setminus \overline{\Omega_-}),$$

$$\sigma(\Delta_{f,h}^{D,(p)}(\Omega_+)) \cap [0, \nu(h)] \subset [0, e^{-\frac{\epsilon}{h}}].$$

If $(u_k^{(1)})_{1 \leq k \leq m_1^D(\Omega_+ \setminus \overline{\Omega_-})}$ is an orthonormal system of the spectral subspace $\text{Ran} 1_{[0, \nu(h)]}(\Delta_{f,h}^{D,(1)}(\Omega_+ \setminus \overline{\Omega_-}))$, then

$$\begin{aligned} \lambda_1^{(0)}(\Omega_+) &= \frac{h^2 \sum_{k=1}^{m_1^D(\Omega_+ \setminus \overline{\Omega_-})} \left| \int_{\partial\Omega_+} e^{-\frac{f}{h}} \mathbf{i}_{n_\sigma} u_k^{(1)} d\sigma \right|^2}{\int_{\Omega_+} e^{-\frac{2f(x)}{h}} dx} (1 + \mathcal{O}(e^{-\frac{\epsilon}{h}})) \\ &= \tilde{\mathcal{O}}(e^{-\frac{2\kappa_f}{h}}). \end{aligned}$$

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Theorem 1 cont'

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The non negative normalized eigenfunction $u_1^{(0)}$ satisfies

$$\left\| u_1^{(0)} - \frac{e^{-\frac{f}{h}}}{\left(\int_{\Omega_+} e^{-\frac{2f(x)}{h}} dx \right)^{1/2}} \right\|_{W^{1,2}(\Omega_+)} = \mathcal{O}(e^{-\frac{c}{h}}),$$
$$\left\| d_{f,h} u_1^{(0)} + \sum_{k=1}^{m_1^D(\Omega_+ \setminus \overline{\Omega_-})} \frac{h \int_{\partial\Omega_+} e^{-\frac{f(\sigma)}{h}} \mathbf{i}_{n_\sigma} \overline{u_k^{(1)}} d\sigma}{\left(\int_{\Omega_+} e^{-\frac{2f(x)}{h}} dx \right)^{1/2}} u_k^1 \right\|_{\Lambda^1 W^{n,2}(\mathcal{V})} = \mathcal{O}(e^{-\frac{\kappa_f + c_V}{h}}),$$

where \mathcal{V} is any neighborhood of $\partial\Omega_+$ lying in $\Omega_+ \setminus \overline{\Omega_-}$ and $c_V > 0$ does not depend on (n, h) .

Changing f

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Corollary

Let f_1 and f_2 satisfy Hyp 1,2,3 and $f_1 = f_2$ in $\Omega_+ \setminus \overline{\Omega_-}$, then

$$\frac{\lambda_1^{(0)}(f_2)}{\lambda_1^{(0)}(f_1)} = \frac{\int_{\Omega_+} e^{-2\frac{f_1(x)}{h}} dx}{\int_{\Omega_+} e^{-2\frac{f_2(x)}{h}} dx} (1 + \mathcal{O}(e^{-\frac{\epsilon}{h}})),$$

$$\frac{\partial_n \left[e^{-\frac{f_2}{h}} u_1^{(0)}(f_2) \right] \Big|_{\partial\Omega_+}}{\| \partial_n \left[e^{-\frac{f_2}{h}} u_1^{(0)}(f_2) \right] \|_{L^1(\partial\Omega_+)}}} = \frac{\partial_n \left[e^{-\frac{f_1}{h}} u_1^{(0)}(f_1) \right] \Big|_{\partial\Omega_+}}{\| \partial_n \left[e^{-\frac{f_1}{h}} u_1^{(0)}(f_1) \right] \|_{L^1(\partial\Omega_+)}}} + \mathcal{O}(e^{-\frac{\epsilon}{h}}) \quad \text{in } L^1(\partial\Omega_+).$$

Theorem 2

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Assume Hyp 1,2,3 and that $f|_{\partial\Omega_+}$ is a Morse function. Then

$$\lambda_1^{(0)}(\Omega_+) = \frac{\int_{\partial\Omega_+} 2\partial_n f(\sigma) e^{-2\frac{f(\sigma)}{h}} d\sigma}{\int_{\Omega_+} e^{-2\frac{f(x)}{h}} dx} \times (1 + \mathcal{O}(h))$$
$$-\frac{\partial_n \left[e^{-\frac{f}{h}} u_1^{(0)} \right] \Big|_{\partial\Omega_+}}{\| \partial_n \left[e^{-\frac{f}{h}} u_1^{(0)} \right] \|_{L^1(\partial\Omega_+)}} = \frac{(2\partial_n f) e^{-\frac{2f}{h}} \Big|_{\partial\Omega_+}}{\| (2\partial_n f) e^{-\frac{2f}{h}} \|_{L^1(\partial\Omega_+)}} + \mathcal{O}(h) \quad \text{in } L^1(\partial\Omega_+, d\sigma).$$

A non-Morse potential

An example with one degenerate minimum and no saddle points.

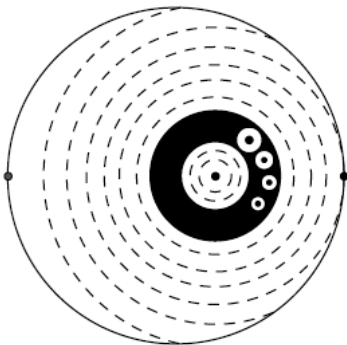


Figure: An example with $\#S = 2$.

In this example when Ω_- is a disk, concentric with Ω_+ ,
 $m_0^N(\Omega_-) = 1$, $m_2^N(\Omega_-) = 5$ and .

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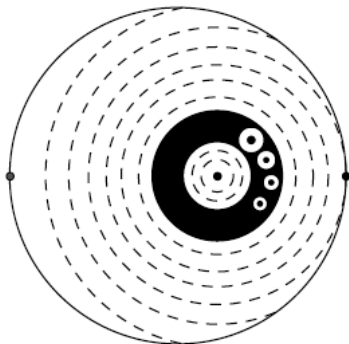


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