

Geometric Kramers-Fokker-Planck operators with boundary conditions

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Reims, april. 17th 2018

Outline

Geometric
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operators
with
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conditions

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- Presentation of the problem
- Main results
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- Elements of proofs

Geometric Kramers-Fokker-Planck operators

Geometric
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In the euclidean space, the operator

$$P_{\pm} = \pm(p \cdot \partial_q - \partial_q V(q) \cdot \partial_p) + \frac{-\Delta_p + |p|^2}{2}, \quad x = (q, p) \in \Omega \times \mathbb{R}^d$$

is associated with the Langevin process

$$dq = p dt, \quad dp = -\partial_q V(q) dt - p dt + dW$$

$\overline{Q} = Q \sqcup \partial Q$ riem. mfl'd with bdy, $X = T^*Q$, $\partial X = T^*_{\partial Q} Q$.

Metric $g = g_{ij}(q) dq^i dq^j$, $g^{-1} = (g^{ij})$

$$P_{\pm, Q, g} = \pm \mathcal{Y}_{\mathcal{E}} + \frac{-\Delta_p + |p|_q^2}{2}, \quad \Delta_p = g_{ij}(q) \partial_{p_i} \partial_{p_j}$$

$$\mathcal{E}(q, p) = \frac{|p|_q^2}{2} = \frac{g^{ij}(q) p_i p_j}{2},$$

$$\mathcal{Y}_{\mathcal{E}} = g^{ij}(q) p_i \partial_{q^j} - \frac{1}{2} \partial_{q^k} g^{ij}(q) p_i p_j \partial_{p_k} = g^{ij}(q) p_i e_j, \quad e_j = \partial_{q^j} + \Gamma_{ij}^{\ell} p_{\ell} \partial_{p_j}.$$

acting on $C^{\infty}(\overline{X}; f)$. $P_{\pm, Q, g}$ = scalar part of Bismut's hypoelliptic Laplacian.

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A simple case

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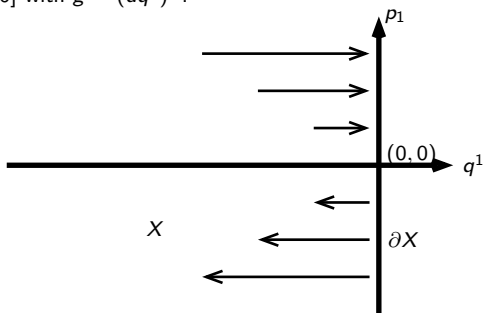
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Take $\bar{Q} = (-\infty, 0]$ with $g = (dq^1)^2$.



Specular reflection: $u(0, -p_1) = u(0, p_1)$ for $p_1 > 0$.

It can be written $\gamma_{\text{odd}} u = 0$ with $\gamma_{\text{odd}} u = \frac{u(0, p_1) - u(0, -p_1)}{2}$.

Absorption: $u(0, p_1) = 0$ for $p_1 < 0$.

It can be written $\gamma_{\text{odd}} u = \text{sign}(p_1) \gamma_{\text{ev}} u$ with $\gamma_{\text{ev}} u = \frac{u(0, p_1) + u(0, -p_1)}{2}$.

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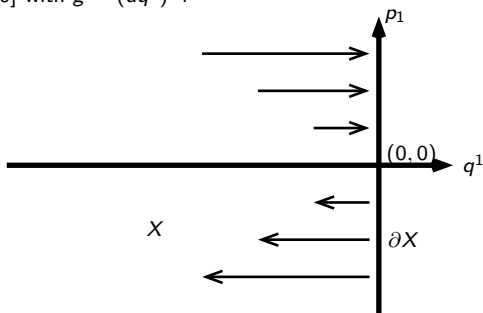
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Metric locally on ∂Q : $(dq^1)^2 \oplus^\perp m(q^1, q')$. Consider \mathfrak{f} -valued functions, \mathfrak{f} Hilbert space.

Let j be a unitary involution in \mathfrak{f} and define along $\partial X = \{q^1 = 0\}$:

$$\gamma_{\text{odd}} = \Pi_{\text{odd}} \gamma = \frac{\gamma(q', p_1, p') - j\gamma(q', -p_1, p')}{2},$$

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$$\gamma_{\text{odd}} u = \pm \text{sign}(p_1) A \gamma_{\text{ev}} u, \quad \Pi_{\text{ev}} A = A \Pi_{\text{ev}}.$$

Formal integration by part

$$\begin{aligned} \text{Re} \langle u, P_{\pm, Q, g} u \rangle &= \frac{\|\nabla_p u\|_{L^2(X, dqdp; \mathfrak{f})}^2 + \| |p|_q u \|_{L^2(X, dqdp; \mathfrak{f})}^2}{2} \pm \frac{1}{2} \int_{\partial X} |\gamma u|(q', p)^2 p_1 dq' dp \\ &= \frac{\|\nabla_p u\|_{L^2(X, dqdp; \mathfrak{f})}^2 + \| |p|_q u \|_{L^2(X, dqdp; \mathfrak{f})}^2}{2} + \underbrace{\text{Re} \langle \gamma_{\text{ev}} u, A \gamma_{\text{ev}} u \rangle_{L^2(\partial X, |p_1| dq' dp; \mathfrak{f})}}. \end{aligned}$$

Assumptions:

- $A = A(q, |p|_q)$ is local in q and $|p|_q$ (local elastic collision at the boundary);
- $A(q, |p|_q) \in \mathcal{L}(L^2(S_{\partial Q}^* Q, |\omega_1| dq' dw; \mathfrak{f}))$ with $\|A(q, r)\| \leq C$ unif.
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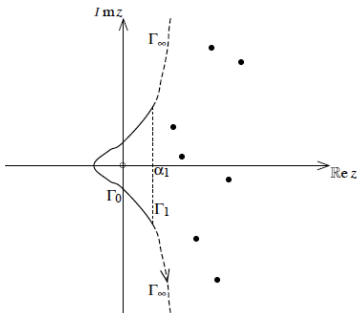
Applications

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Do such boundary conditions with (A, j) define a maximal accretive realization $K_{\pm, A, g}$ of $P_{\pm, Q, g}$?

Can we specify the domain of $K_{\pm, A, g}$ and the regularity (and decay in ρ) estimates for the resolvent ? Global subelliptic estimates ?

$K_{\pm, A, g}$ "cuspidal" ?



Compactness of the resolvent ? Discrete spectrum ? Exponential decay perties of

$$e^{-tK_{\pm, A, g}} = \frac{1}{2i\pi} \int_{\Gamma} e^{-tz} (z - K)^{-1} dz ?$$

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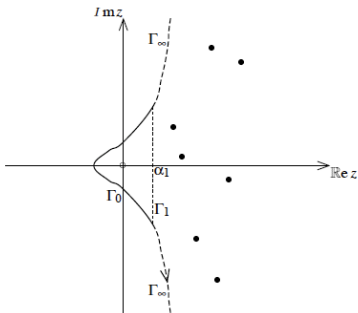
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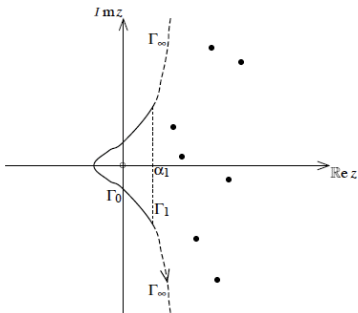
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SDE's: B. Lapeyre (1990) 1D specular reflection, Bossy–Jabir (2011) specular reflection. Bertoin (2007) non-elastic 1D boundary conditions. [Very few results for the PDE interpretation](#)

Quasi Stationary Distribution (\rightarrow molecular dynamics algorithms):
Le Bris–Lelièvre–Luskin–Perez (2012) and Lelièvre–N. (2013) [Elliptic case, Witten Laplacian. But Langevin is a more natural model !](#)

Exponentially small eigenvalues of Witten Laplacians on p -forms in the low temperature limit: Le Peutrec–Viterbo–N. (2013) [Artificial boundary value problems are introduced.](#)

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Notations and first result

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Call $\mathcal{O}_{Q,g} = \frac{-\Delta_p + |p|_q^2}{2}$ and set $\mathcal{H}^{s'}(q) = (d/2 + \mathcal{O}_{Q,g})^{-s'/2} L^2(T_q^*Q, dp; f)$ and globally $\mathcal{H}^{s'} = (d/2 + \mathcal{O}_{Q,g})^{-s'/2} L^2(X, dqdp; f)$. $H^s(Q; \mathcal{H}^{s'})$ is the Sobolev space of H^s -sections of the hermitian fiber bundle $\pi_{\mathcal{H}^{s'}} : \mathcal{H}^{s'} \rightarrow Q$.

Remember the BC's $\gamma_{\text{odd}} u = \pm \text{sign}(p_1) A \gamma_{\text{ev}} u$

- $A \Pi_{\text{ev}} = \Pi_{\text{ev}} A$;
- $A = A(q, |p|_q)$ is local in q and $|p|_q$ (local elastic collision at the boundary);
- $A(q, |p|_q) \in \mathcal{L}(L^2(S_{\partial Q}^* Q, |\omega_1| dq' d\omega; f))$ with $\|A(q, r)\| \leq C$ unif.
- either $\text{Re } A(q, r) \geq c_A > 0$ unif. or $A(q, r) \equiv 0$.

Theorem 1: With the domain $D(K_{\pm, A, g})$ characterized by

$$u \in L^2(Q; \mathcal{H}^1) \quad , \quad P_{\pm, Q, g} u \in L^2(X, dqdp; f), \\ \gamma u \in L^2_{\text{loc}}(\partial X, |p_1| dq' dp; f) \quad , \quad \gamma_{\text{odd}} u = \pm \text{sign}(p_1) A \gamma_{\text{ev}} u,$$

the operator $K_{\pm, A, g} - \frac{d}{2}$ is maximal accretive and

$$\text{Re} \langle u, (K_{\pm, A, g} + \frac{d}{2})u \rangle = \|u\|_{L^2(Q, dq; \mathcal{H}^1)}^2 + \text{Re} \langle \gamma_{\text{ev}} u, A \gamma_{\text{ev}} u \rangle_{L^2(\partial X, |p_1| dq' dp; f)}.$$

The adjoint of $K_{\pm, A, g}$ is $K_{\mp, A^*, g}$.

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The adjoint of $K_{\pm, A, g}$ is $K_{\mp, A^*, g}$.

Subelliptic estimates when $A = 0$

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Theorem 2: When $A=0$ there exists $C > 0$ and for all $\Phi \in C_b^\infty([0, +\infty))$ satisfying $\Phi(0) = 0$ a constant C_Φ such that

$$\begin{aligned} \langle \lambda \rangle^{\frac{1}{4}} \|u\| + \langle \lambda \rangle^{\frac{1}{8}} \|u\|_{L^2(Q; \mathcal{H}^1)} + \|u\|_{H^{1/3}(Q; \mathcal{H}^0)} \\ + \langle \lambda \rangle^{\frac{1}{4}} \|(1 + |p|_q)^{-1} \gamma u\|_{L^2(\partial X, |p_1| dq' dp; f)} \leq C \|(K_{\pm, 0, g} - i\lambda)u\|, \end{aligned}$$

and

$$\|\Phi(d_g(q, \partial Q)) \mathcal{O}_{Q, g} u\| \leq C \|\Phi\|_{L^\infty} \|(K_{\pm, 0, g} - i\lambda)u\| + C_\Phi \|u\|,$$

hold for all $u \in D(K_{\pm, 0, g})$ and all $\lambda \in \mathbb{R}$.

Theorem 3: Assume $\operatorname{Re} A(q, |p|_q) \geq c_A > 0$ uniformly. There exists $C > 0$, for all $t \in [0, \frac{1}{18})$ a constant $C_t > 0$ and for all $\Phi \in C_b^\infty([0, +\infty))$ satisfying $\Phi(0) = 0$ a constant C_Φ such that

$$\begin{aligned} \langle \lambda \rangle^{\frac{1}{4}} \|u\| + \langle \lambda \rangle^{\frac{1}{8}} \|u\|_{L^2(Q; \mathcal{H}^1)} + C_t^{-1} \|u\|_{H^t(Q; \mathcal{H}^0)} \\ + \langle \lambda \rangle^{\frac{1}{8}} \|\gamma u\|_{L^2(\partial X, |p_1| dq' dp; f)} \leq C \|(K_{\pm, A, g} - i\lambda)u\|, \end{aligned}$$

and

$$\|\Phi(d_g(q, \partial Q)) \mathcal{O}_{Q, g} u\| \leq C \|\Phi\|_{L^\infty} \|(K_{\pm, A, g} - i\lambda)u\| + C_\Phi \|u\|,$$

hold for all $u \in D(K_{\pm, A, g})$ and all $\lambda \in \mathbb{R}$.

The operator $K_{\pm, A, g}$ is cuspidal.

When \bar{Q} is compact, $K_{\pm, A, g}^{-1}$ is compact \rightarrow discrete spectrum.

The integration by parts imply $\|u\|_{L^2(Q, \mathcal{H}^1)}^2 \leq \|(K_{\pm, A, g} - i\lambda)u\| \|u\|$ and a potential term $\mp \partial_q V(q) \partial_p$ with V Lipschitz is a nice perturbation \rightarrow All the results are still valid with such a potential term.

PT-symmetry if $UAU^* = A^*$, $UK_{\pm, A, g}U^* = K_{\mp, A^*, g} = K_{\pm, A, g}^*$ when $Uu(q, p) = u(q, -p)$.

The results hold (with additional conditions for the PT-symmetry) when $Q \times f$ is replaced by a hermitian bundle $\pi_F : F \rightarrow Q$ with a metric g^F and a connection ∇^F . The pull-back bundle $F_X = \pi^*F$ with $\pi : \bar{X} = \bar{T}^*\bar{Q} \rightarrow \bar{Q}$ is then endowed with the metric $g^{FX} = \pi^*g^F$ and the connection

$$\nabla_{e_j}^{FX} = \nabla_{\partial_j}^F, \quad \nabla_{\partial_{p_j}}^{FX} = 0.$$

Covariant derivative $\tilde{\nabla}_T^{FX}(s^k(x)f_k) = Ts^k(x)f_k + s^k(x)\nabla_T^{FX}f_k$. $x = (q, p)$.

DEF: General geometric Kramers-Fokker-Planck operator (including hypoelliptic Laplacian)

$$\pm g^{ij}(q)p_i \tilde{\nabla}_{e_j}^{FX} + \mathcal{O}_{Q, g} + M_j^0(q, p) \tilde{\nabla}_{\partial_{p_j}}^{FX} + M^1(q, p),$$

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Specular reflection: $j = 1, A = 0$.

Absorption: $j = 1, A = \text{Id}$.

The two above cases can be interpreted in terms of stochastic processes by completing the Langevin process with a jump process when $X(t)$ hits the boundary:

- For specular reflection the jump changes the velocity (p_1, p') with $p_1 > 0$ into $(-p_1, p')$;
- For the absorption, the particle is sent to an external stationary point ϵ when the particle hits the boundary.

More general jump processes: Set $\partial X_{\pm} = \{(0, q', p_1, p'), \pm p_1 > 0\}$. More general Markov kernel from ∂X_+ to $\partial X_- \sqcup \{\epsilon\}$ can be considered. $\text{Re } A \geq c_A$ means that a positive fraction is sent to ϵ

Doubling the manifold: In the position variable the Neumann and Dirichlet boundary value problems for $-\Delta_q$ can be introduced by considering even and odd solutions after the extension by reflection $(q^1, q') \rightarrow (-q^1, q')$. Here the extension by reflection is $(q^1, q', p_1, p') \rightarrow (-q^1, q', -p_1, p')$.

- Even case=specular reflection: $j = 1$ and $A = 0$.
- Odd case: $j = -1$ and $A = 0 \rightarrow$ does not preserve the positivity. In the elliptic case, considered recently by K.T. Sturm via the stochastic dynamics of signed particles.

Scalar case: $f = \mathbb{C}$

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More general jump processes: Set $\partial X_{\pm} = \{(0, q', p_1, p'), \pm p_1 > 0\}$. More general Markov kernel from ∂X_+ to $\partial X_- \sqcup \{\epsilon\}$ can be considered. $\text{Re } A \geq c_A$ means that a positive fraction is sent to ϵ

Doubling the manifold: In the position variable the Neumann and Dirichlet boundary value problems for $-\Delta_q$ can be introduced by considering even and odd solutions after the extension by reflection $(q^1, q') \rightarrow (-q^1, q')$. Here the extension by reflection is $(q^1, q', p_1, p') \rightarrow (-q^1, q', -p_1, p')$.

- Even case=specular reflection: $j = 1$ and $A = 0$.
- Odd case: $j = -1$ and $A = 0 \rightarrow$ does not preserve the positivity. In the elliptic case, considered recently by K.T. Sturm via the stochastic dynamics of signed particles.

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Scalar case: $f = \mathbb{C}$

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A proposal for “Dirichlet” and “Neumann” realizations of the hypoelliptic Laplacian.

Position space $\bar{Q} = Q \sqcup \partial Q \ni q$, phase-space $X = T^*Q$,
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Hypoelliptic Laplacian: partial differential operator acting on differential forms = sections of $\wedge T^*X$, of which the main part is a scalar geometric KFP operator. (REF Bismut and Lebeau).

With the basis $(e^I \hat{e}_J = e^{i_1} \wedge \dots \wedge e^{i_{|I|}} \wedge \hat{e}_{j_1} \wedge \dots \wedge \hat{e}_{j_{|J|}})$ of $\wedge T_x^*X$, $x \in \partial X$,
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(unitary involution for $k = 0$ or $k = 1$)

“Neumann” realization: Take $k = 0$, $j = \mathbf{j}_0$ and $A = 0$.

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It is a very classical one for boundary value problems (see for example Hörmander-Chap 20 or Boutet de Montvel (1970))

Have a good understanding of the simplest $1D$ -problem.

Use some separation of variables for straight half-spaces.

Look at the general local problem by sending it to the straight half-space problem with a change of variables and try to absorb the corresponding perturbative terms.

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Pb 1 The simplest 1D problem is actually a 2D-problem with p -dependent coefficients. Moreover it looks like a corner problem.

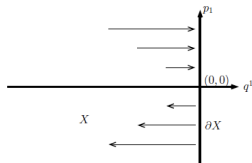


Fig.1: The boundary $\partial X = \{q^1 = 0\}$ and the vector field $p_1 \partial_{q^1}$ are represented. For the absorbing case, the boundary condition says $\gamma u(p_1) = 0$ for $p_1 < 0$ and corresponds to the case ($j = 1$ and $A = 1$).

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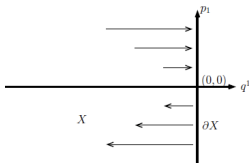


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Pb 2 For a general boundary one has to face the pb of glancing rays.

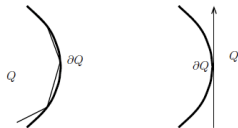


Fig.2: The left picture show a (approximately) gliding ray and the right one a grazing ray.

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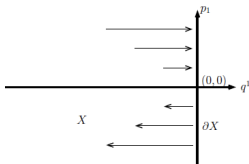


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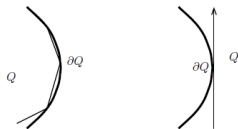


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Pb 1 solved by introducing adapted Fourier series and a quantization of the function $\text{sign}(p_1)$.

Pb 2 solved by introducing a dyadic partition of unity in the p -variable and by using the 2nd resolvent formula for the corresponding semiclassical problems ($\hbar = 2^{-j}$).

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This solves only the basic functional analysis.

There are still a lot of things to be investigated:

Non self-adjoint spectral problems.

Boundary value problems.

Parameter dependent asymptotics (large friction, small temperature=semiclassical).

Multiple wells and tunnel effect...