Exact integration for products of power of barycentric coordinates over d-simplexes in \mathbb{R}^n

François Cuvelier*

2022/03/04

Abstract

Exact integral computation over a d-simplex in \mathbb{R}^n for products of powers of its barycentric coordinates is done in [9] by using mathematical induction and coordinate mappings. In this note we give a new proof using Laplace transformations without mathematical induction.

Contents

1	Not	ations and definitions	2
2	2.1	ne results on the unit d-simplex unit d-simplex volume	
3	Son	ne results on a d-simplex in \mathbb{R}^n	6
	3.1	Gradients of barycentric coordinates on a <i>d</i> -simplex	6
	3.2	Integration over a <i>d</i> -simplex	8
	3.3	Volume of a d-simplex	11
	3.4		

^{*}Université Paris 13, Sorbonne Paris Cité, LAGA, CNRS UMR 7539, 99 Avenue J-B Clément, F-93430 Villetaneuse, France, cuvelier@math.univ-paris13.fr

This work was supported in part by the ANR project DEDALES under grant ANR-14-CE23-0005.

Local shape functions of a large variety of finite element on a d-simplex $K \subset \mathbb{R}^n$ can be expressed in function of the barycentric coordinates $\{\lambda_0, \ldots, \lambda_d\}$ of K and their derivatives (see [1] for examples).

In [9], the authors give a proof of the magic formula: let $\boldsymbol{\nu} = (\nu_0, \dots, \nu_d) \in \mathbb{N}^{d+1}$, then

$$\int_{K} \prod_{i=0}^{\mathbf{d}} \lambda_{i}^{\nu_{i}}(\mathbf{q}) d\mathbf{q} = \mathbf{d}! |K| \frac{\prod_{i=0}^{\mathbf{d}} \nu_{i}!}{(\mathbf{d} + \sum_{l=0}^{\mathbf{d}} \nu_{i})!}$$
(1)

where |K| is the volume of K. In their proof, mathematical induction and coordinate mappings are mainly used. In this note we give a new proof of this formula using Laplace transformations without mathematical induction.

Firstly we recall definitions of a d-simplex in \mathbb{R}^n and of its barycentric coordinates. Therafter we introduce Laplace transforms to compute the volume of the unit d-simplex $\hat{K} \subset \mathbb{R}^d$ and the *magic formula* (1) over \hat{K} . In the last section, we propose to compute the gradients of the barycentric coordinates by solving linear systems. We also present the mapping of an integral over a d-simplex in \mathbb{R}^n to the reference unit d-simplex, allowing to proove (1).

1 Notations and definitions

Let $n \in \mathbb{N}^*$ be the space dimension and $d \in [0, n]$. We recall the definition of a d-simplex in \mathbb{R}^n as well as its barycentric coordinates.

Definition 1 (d-simplex) A d-simplex $K \subset \mathbb{R}^n$ is the convex hull of (d + 1) points $\mathbf{q}^0, \ldots, \mathbf{q}^d$ of \mathbb{R}^n which form the vertices of K.

$$K = \left\{ \mathbf{q} \in \mathbb{R}^n \mid \mathbf{q} = \sum_{i=0}^{d} \theta_i \mathbf{q}^i, \text{ with } \forall i \in [[0,d]], \ \theta_i \ge 0, \text{ and } \sum_{i=0}^{d} \theta_i = 1 \right\}.$$
(2)

For example, a 2-simplex is a triangle and a 3-simplex is a tetrahedron. It will be always assumed that a d-simplex is **not degenerated**, i.e., the set of vectors $\{\mathbf{q}^i - \mathbf{q}^0\}_{i=1}^d$ is linearly independent.

Definition 2 (Barycentric coordinates) Let $K \subset \mathbb{R}^n$ be a non-degenerate d-simplex and $\{\mathbf{q}^i\}_{i=0}^d$ its vertices. The parametrization of K with a convex combination of the vertices reads as follows

$$K = \left\{ \mathbf{q} \in \mathbb{R}^n \mid \mathbf{q} = \sum_{i=0}^{d} \lambda_i(\mathbf{q}) \mathbf{q}^i, \text{ with } \forall i \in [[0,d]], \lambda_i(\mathbf{q}) \ge 0, \text{ and } \sum_{i=0}^{d} \lambda_i(\mathbf{q}) = 1 \right\}$$
(3)

The coefficients $\lambda_0(\mathbf{q}), \ldots, \lambda_d(\mathbf{q})$ are called the barycentric coordinates on K of \mathbf{q} .

As immediat property, the barycentric coordinates on K satisfy

$$\lambda_i(\mathbf{q}^j) = \delta_{i,j}, \quad \forall (i,j) \in [\![0,d]\!].$$
(4)

2 Some results on the unit d-simplex

The **unit** d-simplex $\hat{K}^{d} \subset \mathbb{R}^{d}$ is defined by the d + 1 vertices

$$\left\{ \hat{\mathbf{q}}^{0}, \hat{\mathbf{q}}^{1}, \cdots, \hat{\mathbf{q}}^{d} \right\} = \left\{ \mathbf{0}, \hat{\mathbf{e}}^{1}, \cdots, \hat{\mathbf{e}}^{d} \right\}$$

where $\{\hat{\boldsymbol{e}}^1, \dots, \hat{\boldsymbol{e}}^d\}$ is the standard basis of \mathbb{R}^d . We have

$$\hat{K}^{d} = \left\{ \hat{\mathbf{q}} \in \mathbb{R}^{d} \mid \hat{\mathbf{q}} = \sum_{i=0}^{d} \hat{\lambda}_{i}(\hat{\mathbf{q}}) \hat{\mathbf{q}}^{i}, \text{ with } \hat{\lambda}_{i}(\hat{\mathbf{q}}) \ge 0, \text{ and } \sum_{i=0}^{d} \hat{\lambda}_{i}(\hat{\mathbf{q}}) = 1 \right\}.$$
 (5)

As immediat property, the barycentric coordinates $(\hat{\lambda}_i)_{i=0}^{\rm d}$ on $\hat{K}^{\rm d}$ satisfy

$$\hat{\lambda}_i(\hat{\mathbf{q}}^j) = \delta_{i,j}, \quad \forall (i,j) \in [\![0,d]\!].$$
(6)

and are explicitly given with $\hat{\mathbf{q}} = (x_1, \cdots, x_d)^{\mathsf{t}} \in \hat{K}^{\mathrm{d}}$ by

$$\hat{\lambda}_0(\hat{\mathbf{q}}) = 1 - \sum_{i=1}^{d} x_i \quad \text{and } \forall i \in [\![1,d]\!], \quad \hat{\lambda}_i(\hat{\mathbf{q}}) = x_i.$$
(7)

Indeed, as $\hat{\mathbf{q}}^0 = \mathbf{0}$, we have

$$\hat{\mathbf{q}} = \sum_{i=0}^{d} \hat{\lambda}_i(\hat{\mathbf{q}}) \hat{\mathbf{q}}^i = \sum_{i=1}^{d} \hat{\mathbf{q}}^i \hat{\lambda}_i(\hat{\mathbf{q}})$$

From $\hat{\mathbf{q}}^i = \hat{\boldsymbol{e}}^i, \ \forall i \in [\![1, d]\!]$, we obtain

$$\sum_{i=1}^{d} \hat{\mathbf{q}}^{i} \hat{\lambda}_{i}(\hat{\mathbf{q}}) = \left(\begin{array}{c|c} \hat{\mathbf{q}}^{1} & \cdots & \hat{\mathbf{q}}^{d} \end{array} \right) \begin{pmatrix} \hat{\lambda}_{1}(\hat{\mathbf{q}}) \\ \vdots \\ \hat{\lambda}_{d}(\hat{\mathbf{q}}) \end{pmatrix} = \mathbb{I}_{d} \begin{pmatrix} \hat{\lambda}_{1}(\hat{\mathbf{q}}) \\ \vdots \\ \hat{\lambda}_{d}(\hat{\mathbf{q}}) \end{pmatrix} = \begin{pmatrix} \hat{\lambda}_{1}(\hat{\mathbf{q}}) \\ \vdots \\ \hat{\lambda}_{d}(\hat{\mathbf{q}}) \end{pmatrix}$$

and thus

$$\hat{\mathbf{q}} = \begin{pmatrix} x_1 \\ \vdots \\ x_d \end{pmatrix} = \begin{pmatrix} \hat{\lambda}_1(\hat{\mathbf{q}}) \\ \vdots \\ \hat{\lambda}_d(\hat{\mathbf{q}}) \end{pmatrix}.$$

From (5), we have

$$\sum_{i=0}^{\mathbf{d}} \hat{\lambda}_i(\hat{\mathbf{q}}) = 1$$

and thus

$$\hat{\lambda}_0(\hat{\mathbf{q}}) = 1 - \sum_{i=1}^{d} \hat{\lambda}_i(\hat{\mathbf{q}}) = 1 - \sum_{i=1}^{d} x_i.$$

2.1 unit d-simplex volume

There are several ways to compute the volume $|\hat{K}|$ of the d-simplex $\hat{K} \subset \mathbb{R}^d$ which is given by the following integral:

$$|\hat{K}| = \int_{\hat{K}} 1d\hat{\mathbf{q}} = \int_{0}^{1} \int_{0}^{1-x_{1}} \int_{0}^{1-x_{1}-x_{2}} \dots \int_{0}^{1-(x_{1}+\dots+x_{d-1})} 1dx_{d} \dots dx_{3}dx_{2}dx_{1}.$$

An elegant way to perform this integration is explained in [6], section 18.10, and uses a Laplace transform. To use this method, we note that

$$\hat{K} = \mathbb{R}^{d}_{+} \cap \{1 - (x_1 + \ldots + x_d) \ge 0\}.$$
 (8)

So we also have

$$|\hat{K}| = \int_{\mathbb{R}^{d}_{+} \cap \{1 - (x_{1} + \dots + x_{d}) \ge 0\}} 1 dx_{d} \dots dx_{1}.$$

By using a dirac function and extending the integration domain to $\mathbb{R}^{d+1}_+,$ we also have

$$|\hat{K}| = \int_{\mathbb{R}^{d+1}_+} \delta(x_1 + \ldots + x_d + x_{d+1} - 1) dx_{d+1} dx_d \ldots dx_1$$

To use the Laplace transform theory, we define the function f by

$$f(t) = \int_{\mathbb{R}^{d+1}_+} \delta(x_1 + \ldots + x_d + x_{d+1} - t) dx_{d+1} dx_d \ldots dx_1$$

so that $|\hat{K}| = f(1)$. The Laplace transform of f is given by

$$\begin{aligned} \mathcal{L}(f)(s) &= \int_{0}^{\infty} f(t)e^{-st}dt \\ &= \int_{\mathbb{R}^{d+1}_{+}} \left(\int_{0}^{\infty} \delta(x_{1} + \ldots + x_{d} + x_{d+1} - t)e^{-st}dt \right) dx_{d+1}dx_{d} \ldots dx_{1} \\ &= \int_{\mathbb{R}^{d+1}_{+}} \exp(-s\sum_{i=1}^{d+1} x_{i})dx_{d+1}dx_{d} \ldots dx_{1} \\ &= \prod_{i=1}^{d+1} \int_{0}^{\infty} \exp(-sx_{i})dx_{i} \\ &= \frac{1}{s^{d+1}}. \end{aligned}$$

By using the inverse Laplace transform table (see [8] for example), we have

$$\mathcal{L}^{-1}(s \mapsto \frac{\mathrm{d}!}{s^{\mathrm{d}+1}})(t) = t^{\mathrm{d}}.$$

As $f = \mathcal{L}^{-1} \circ \mathcal{L}(f)$ and by linearity of the inverse Laplace transform we obtain

$$f(t) = \frac{t^{\mathrm{d}}}{\mathrm{d}!}.$$

So the volume of the unit d-simplex is

$$|\hat{K}| = \frac{1}{\mathrm{d}!} \tag{9}$$

2.2 Magic formula

Let $\boldsymbol{\nu} = (\nu_0, \dots, \nu_d) \in \mathbb{N}^{d+1}$. The magic formula is given by

$$\int_{\hat{K}} \prod_{i=0}^{d} \hat{\lambda}_{i}^{\nu_{i}}(\hat{\mathbf{q}}) d\hat{\mathbf{q}} = \frac{\prod_{i=0}^{d} \nu_{i}!}{(d + \sum_{i=0}^{d} \nu_{i})!}$$
(10)

This formula is often used in \mathbb{P}^1 -Lagrange finite element methods because \mathbb{P}^1 -Lagrange basis functions on a d-simplex are the associated barycentic coordinates. For example, one can refer to [7] (section 8.2.1, page 179, formula (8.3)), [9], [4] section 7.3.3 page 126, [3] for $d \in [1,3]$, [2] as exercise for d = 2 and d = 3. In this section, we propose a proof of this formula using Laplace transform theory. Let $\hat{I}(\nu)$ denote the integral of (10). The barycentic coordinates $\hat{\lambda}_i$ are given in (7) and so with $\hat{\mathbf{q}} = (x_1, \ldots, x_d)$ and using (8) we obtain

$$\hat{I}(\nu) = \int_{\hat{K}} (1 - \sum_{i=1}^{d} x_i)^{\nu_0} \prod_{i=1}^{d} x_i^{\nu_i} dx_d \dots dx_1$$
$$= \int_{\mathbb{R}^d_+ \cap \{1 - (x_1 + \dots + x_d) \ge 0\}} (1 - \sum_{i=1}^{d} x_i)^{\nu_0} \prod_{i=1}^{d} x_i^{\nu_i} dx_d \dots dx_1$$

From section 2.1, by using a dirac function and by extending the integration domain to \mathbb{R}^{d+1}_+ we obtain with $\nu_{d+1} = \nu_0$

$$\hat{I}(\nu) = \int_{\mathbb{R}^{d+1}_+} \delta(x_1 + \ldots + x_d + x_{d+1} - 1) x_{d+1}^{\nu_0} \prod_{i=1}^d x_i^{\nu_i} dx_{d+1} dx_d \ldots dx_1$$
$$= \int_{\mathbb{R}^{d+1}_+} \delta(x_1 + \ldots + x_d + x_{d+1} - 1) \prod_{i=1}^{d+1} x_i^{\nu_i} dx_{d+1} dx_d \ldots dx_1$$

To use the Laplace transform theory, we define the function $f_{\boldsymbol{\nu}}$ by

$$f_{\boldsymbol{\nu}}(t) = \int_{\mathbb{R}^{d+1}_+} \delta(x_1 + \ldots + x_d + x_{d+1} - t) \prod_{i=1}^{d+1} x_i^{\nu_i} dx_{d+1} dx_d \ldots dx_1$$

so that $\hat{I}(\nu) = f_{\nu}(1)$. The Laplace transform of f_{ν} is given by

$$\begin{aligned} \mathcal{L}(f_{\boldsymbol{\nu}})(s) &= \int_{0}^{\infty} f_{\boldsymbol{\nu}}(t) e^{-st} dt \\ &= \int_{\mathbb{R}^{d+1}_{+}} \left(\int_{0}^{\infty} \delta(x_{1} + \ldots + x_{d} + x_{d+1} - t) e^{-st} dt \right) \prod_{i=1}^{d+1} x_{i}^{\nu_{i}} \, dx_{d+1} dx_{d} \ldots dx_{1} \\ &= \int_{\mathbb{R}^{d+1}_{+}} \exp(-s \sum_{i=1}^{d+1} x_{i}) \prod_{i=1}^{d+1} x_{i}^{\nu_{i}} \, dx_{d+1} dx_{d} \ldots dx_{1} \\ &= \prod_{i=1}^{d+1} \int_{0}^{\infty} x_{i}^{\nu_{i}} \exp(-sx_{i}) dx_{i} \\ &= \prod_{i=1}^{d+1} \mathcal{L}(t \mapsto t^{\nu_{i}})(s) \end{aligned}$$

In a classical Laplace transform table (see [8] for example), we have

$$\mathcal{L}(t \mapsto \frac{t^k}{k!})(s) = \frac{1}{s^{k+1}}$$

and by linearity of the Laplace transform

$$\mathcal{L}(t \mapsto t^k)(s) = \frac{k!}{s^{k+1}}$$

So we obtain

$$\mathcal{L}(f_{\boldsymbol{\nu}})(s) = \prod_{i=1}^{d+1} \frac{\nu_i!}{s^{\nu_i+1}} = \frac{\prod_{i=1}^{d+1} \nu_i!}{s^{d+1+\sum_{i=1}^{d+1} \nu_i}}$$

By using the inverse Laplace transform table, we have

$$\mathcal{L}^{-1}(s \mapsto \frac{1}{s^k})(t) = \frac{t^{k-1}}{k-1}$$

With the linearity of the inverse Laplace transform we obtain

$$f_{\boldsymbol{\nu}}(t) = \mathcal{L}^{-1}(\mathcal{L}(f_{\boldsymbol{\nu}})(s))(t)$$

= $\frac{\prod_{i=1}^{d+1} \nu_i!}{(d + \sum_{i=1}^{d+1} \nu_i)!} t^{d + \sum_{i=1}^{d+1} \nu_i}$

As $\hat{I}(\nu) = f_{\nu}(1)$ and $\nu_{d+1} = \nu_0$, the equation (10) is proved.

3 Some results on a d-simplex in \mathbb{R}^n

3.1 Gradients of barycentric coordinates on a *d*-simplex

Lemma 3 Let $K \subset \mathbb{R}^n$ be a non-degenerate d-simplex and and $\{\mathbf{q}^i\}_{i=0}^d$ its vertices. The barycentric coordinates $(\lambda_i(\mathbf{q}))_{i=0}^d$ are solution of the linear system

$$\begin{pmatrix} 1 & 1 & \cdots & 1 \\ 0 & & & \\ \vdots & & & \\ 0 & & & \end{pmatrix} \begin{pmatrix} \lambda_0(\mathbf{q}) \\ \lambda_1(\mathbf{q}) \\ \vdots \\ \lambda_d(\mathbf{q}) \end{pmatrix} = \begin{pmatrix} 1 \\ \mathbb{A}_K(\mathbf{q} - \mathbf{q}^0) \\ \mathbb{A}_K(\mathbf{q} - \mathbf{q}^0) \end{pmatrix}$$
(11)

where $\mathbb{A}_K \in \mathcal{M}_{n,d}(\mathbb{R})$ is defined by

$$\mathbb{A}_{K} = \left(\begin{array}{c|c} \mathbf{q}^{1} - \mathbf{q}^{0} \end{array} \middle| \cdots \middle| \begin{array}{c} \mathbf{q}^{d} - \mathbf{q}^{0} \end{array} \right)$$
(12)

The barycentric coordinates are multivariate polynomials of first degree and their gradients are given by

$$\left(\left. \boldsymbol{\nabla} \lambda_{1}(\mathbf{q}) \right| \cdots \left| \left. \boldsymbol{\nabla} \lambda_{d}(\mathbf{q}) \right. \right) = \mathbb{A}_{K} (\mathbb{A}_{K}^{t} \mathbb{A}_{K})^{-1}$$
(13)

and

$$\boldsymbol{\nabla} \lambda_0(\mathbf{q}) = -\sum_{i=1}^{d} \boldsymbol{\nabla} \lambda_i(\mathbf{q}).$$
(14)

Proof: As $\sum_{i=0}^{d} \lambda_i(\mathbf{q}) = 1$, we have

$$\mathbf{q} = \sum_{i=0}^{d} \lambda_i(\mathbf{q}) \mathbf{q}^i \Longrightarrow \mathbf{q} - \mathbf{q}^0 = \sum_{i=1}^{d} (\mathbf{q}^i - \mathbf{q}^0) \lambda_i(\mathbf{q}) = \mathbb{A}_K \begin{pmatrix} \lambda_1(\mathbf{q}) \\ \vdots \\ \lambda_d(\mathbf{q}) \end{pmatrix}$$

Due to linear independence of $\{ \mathbf{q}^i - \mathbf{q}^0 \}_{i=1}^{\mathrm{d}},$

$$\mathbb{H}_{K} \stackrel{\text{def}}{=} \mathbb{A}_{K}^{t} \mathbb{A}_{K} \in \mathcal{M}_{d,d}(\mathbb{R}) \tag{15}$$

is a regular matrix and the barycentric coordinates are solution of the linear system

$$\mathbb{A}_{K}^{t}\mathbb{A}_{K}\begin{pmatrix}\lambda_{1}(\mathbf{q})\\\vdots\\\lambda_{d}(\mathbf{q})\end{pmatrix} = \mathbb{A}_{K}^{t}(\mathbf{q}-\mathbf{q}^{0}) \text{ and } \sum_{i=0}^{d}\lambda_{i}(\mathbf{q}) = 1.$$

In matrix form these equations can be written as (11) and we deduce that the barycentric coordinates λ_i are multivariate polynomials of first degree. So their gradients are constants on K.

The affine map/transformation \mathcal{F}_K from the unit d-simplex $\hat{K} \subset \mathbb{R}^d$ to $K \subset \mathbb{R}^n$ is given by

$$\mathbf{q} = \mathbb{A}_K \hat{\mathbf{q}} + \mathbf{q}^0 = \mathcal{F}_K(\hat{\mathbf{q}}).$$
(16)

So we have

$$\mathbb{A}_{K}^{t}(\mathbf{q}-\mathbf{q}^{0})=\mathbb{A}_{K}^{t}\mathbb{A}_{K}\hat{\mathbf{q}}=\mathbb{H}_{K}\hat{\mathbf{q}}$$

and thus $\mathcal{F}_K^{-1}: K \subset \mathbb{R}^n \longrightarrow \hat{K} \subset \mathbb{R}^d$ is defined by

$$\hat{\mathbf{q}} = \mathbb{H}_{K}^{-1} \mathbb{A}_{K}^{\mathsf{t}} (\mathbf{q} - \mathbf{q}^{0}) = \mathcal{F}_{K}^{-1} (\mathbf{q}).$$
(17)

So we have

$$\lambda_i(\mathbf{q}) = (\hat{\lambda}_i \circ \mathcal{F}_K^{-1})(\mathbf{q}) \text{ and } \hat{\lambda}_i(\hat{\mathbf{q}}) = (\lambda_i \circ \mathcal{F}_K)(\hat{\mathbf{q}})$$
(18)

One can remark that if d = n then \mathbb{A}_K is a regular square matrix and $\mathbb{H}_K^{-1}\mathbb{A}_K^t = \mathbb{A}_K^{-1}$.

Now, we may compute partial derivative of λ_i and $\forall i \in [0, d], \forall j \in [1, n]$, we obtain with $\hat{\mathbf{q}} = (\hat{x}_1, \dots, \hat{x}_d)$ and $\mathbf{q} = (x_1, \dots, x_n)$

$$\frac{\partial \lambda_i}{\partial x_j}(\mathbf{q}) = \sum_{l=1}^d \frac{\partial \hat{\lambda}_i}{\partial \hat{x}_j} (\mathcal{F}_K^{-1}(\mathbf{q})) \frac{\partial \mathcal{F}_{K,l}^{-1}}{\partial x_j}(\mathbf{q})$$

From (17), denoting $\mathbb{B}_K = \mathbb{H}_K^{-1} \mathbb{A}_K^t \in \mathcal{M}_{d,m}(\mathbb{R})$ gives $\frac{\partial \mathcal{F}_{K,l}^{-1}}{\partial x_j}(\mathbf{q}) = (\mathbb{B}_K)_{l,j}$. The barycentric coordinates are polynomials of first degree, so their gradients are constants and we obtain

$$\boldsymbol{\nabla}\lambda_i = \mathbb{B}_K^{\mathsf{t}} \boldsymbol{\nabla} \lambda_i$$

(in fact \mathbb{B}_K is the Jacobian matrix of \mathcal{F}_K^{-1}). The matrix \mathbb{H}_K is regular and symmetric, so $\mathbb{B}_K^t = \mathbb{A}_K \mathbb{H}_K^{-1}$ and we obtain

$$\boldsymbol{\nabla}\lambda_i = \mathbb{A}_K \mathbb{H}_K^{-1} \hat{\boldsymbol{\nabla}} \hat{\lambda}_i.$$
⁽¹⁹⁾

From (7), we deduced that

$$\left(\begin{array}{c|c} \hat{\boldsymbol{\nabla}} \hat{\lambda}_1 & \cdots & \hat{\boldsymbol{\nabla}} \hat{\lambda}_d \end{array} \right) = \mathbb{I}_d$$

$$(20)$$

and thus

$$\left(\begin{array}{c|c} \boldsymbol{\nabla} \lambda_{1}(\mathbf{q}) & \cdots & \boldsymbol{\nabla} \lambda_{d}(\mathbf{q}) \end{array} \right) = \mathbb{A}_{K} \mathbb{H}_{K}^{-1} \left(\begin{array}{c|c} \hat{\boldsymbol{\nabla}} \hat{\lambda}_{1} & \cdots & \hat{\boldsymbol{\nabla}} \hat{\lambda}_{d} \end{array} \right)$$
$$= \mathbb{A}_{K} \mathbb{H}_{K}^{-1}.$$

As $\sum_{i=0}^{d} \lambda_i(\mathbf{q}) = 1$, we immediately have

$$\nabla \lambda_0(\mathbf{q}) = -\sum_{i=1}^{\mathrm{d}} \nabla \lambda_i(\mathbf{q}).$$

From (13) and (14), we immediatly have:

Remark 4 The gradients of the barycentric coordinates are linear combinations of $\{\mathbf{q}^1 - \mathbf{q}^0, \dots, \mathbf{q}^d - \mathbf{q}^0\}$.

3.2 Integration over a *d*-simplex

If K is a non-degenerated d-simplex in \mathbb{R}^d , from (16) we have $\mathcal{J}_{\mathcal{F}_K}(\hat{\mathbf{q}}) = \mathbb{A}_K$. Then \mathbb{A}_K is a regular square matrix and we have the classical formula:

$$\int_{K} f(\mathbf{q}) d\mathbf{q} = |\det(\mathbb{A}_{K})| \int_{\hat{K}} f \circ \mathcal{F}_{K}(\hat{\mathbf{q}}) d\hat{\mathbf{q}}$$
(21)

The following theorem extend this result to d-simplex in \mathbb{R}^n , with $1 \leq d \leq n$.

Theorem 5 Let $K \subset \mathbb{R}^n$ be a non-degenerated d-simplex and $f: K \longrightarrow \mathbb{R}$.

$$\int_{K} f(\mathbf{q}) d\mathbf{q} = \left| \det(\mathbb{A}_{K}^{t} \mathbb{A}_{K}) \right|^{1/2} \int_{\hat{K}} f \circ \mathcal{F}_{K}(\hat{\mathbf{q}}) d\hat{\mathbf{q}}$$
(22)

where \hat{K} is the unit d-simplex in \mathbb{R}^d , $\mathbb{A}_K \in \mathcal{M}_{n,d}(\mathbb{R})$ is defined by

$$\mathbb{A}_{K} = \left(\begin{array}{c|c} \mathbf{q}^{1} - \mathbf{q}^{0} & \mathbf{q}^{2} - \mathbf{q}^{0} & \cdots & \mathbf{q}^{d} - \mathbf{q}^{0} \end{array} \right)$$
(23)

and $\mathcal{F}_K : \hat{K} \longrightarrow K$ is given by

$$\mathcal{F}_K(\hat{\mathbf{q}}) = \mathbb{A}_K \hat{\mathbf{q}} + \mathbf{q}^0 \tag{24}$$

Proof: The set $\{\boldsymbol{v}^1, \ldots, \boldsymbol{v}^d\}$ is linearly independent so we can extend it to a basis $\{\boldsymbol{v}^1, \ldots, \boldsymbol{v}^n\}$. We denote by $\mathbb{A} \in \mathcal{M}_{n,n}(\mathbb{R})$ the matrix such that the *i*-th column is the vector \boldsymbol{v}^i for all $i \in [1, n]$. So we have

$$\mathbb{A} = \left(\mathbb{A}_K \mid \boldsymbol{v}^{\mathrm{d}+1} \mid \cdots \mid \boldsymbol{v}^n \right)$$
(25)

By the $\mathbb{Q}\mathbb{R}$ -factorization theorem apply to the matrix $\mathbb{A} \in \mathcal{M}_n(\mathbb{R})$, there is an orthogonal matrix $\mathbb{Q} \in \mathcal{M}_n(\mathbb{R})$ and a regular upper triangular matrix $\mathbb{R} \in \mathcal{M}_n(\mathbb{R})$ such that

$$\mathbb{A}=\mathbb{Q}\mathbb{R}$$

So we have

$$\mathbb{Q}^{\mathtt{t}}\mathbb{A}=\mathbb{R}$$

and we define the matrix $\overline{\mathbb{A}} \in \mathcal{M}_{n,d}(\mathbb{R})$ to be the first d columns of \mathbb{R} :

$$\bar{\mathbb{A}} = \mathbb{Q}^{\mathsf{t}} \mathbb{A}_K$$

We can also note that

$$\bar{\mathbb{A}} = \left(\left. \bar{\mathbf{q}}^1 - \bar{\mathbf{q}}^0 \right| \left. \bar{\mathbf{q}}^2 - \bar{\mathbf{q}}^0 \right| \cdots \left| \left. \bar{\mathbf{q}}^d - \bar{\mathbf{q}}^0 \right. \right) = \left(\left. \bar{\mathbf{q}}^1 \right| \left. \bar{\mathbf{q}}^2 \right| \cdots \left| \left. \bar{\mathbf{q}}^d \right. \right).$$

Let $\bar{\mathcal{F}}: \mathbb{R}^n \longrightarrow \mathbb{R}^n$ be the bijective function defined by

$$\bar{\mathcal{F}}(\boldsymbol{x}) = \mathbb{Q}^{t}(\boldsymbol{x} - \boldsymbol{q}^{0}) = \bar{\boldsymbol{x}}$$
(26)

and

$$\mathbf{\bar{q}}^i = \bar{\mathcal{F}}(\mathbf{q}^i) = \mathbb{Q}^t(\mathbf{q}^i - \mathbf{q}^0), \quad \forall i \in [[0, d]].$$

By construction $\bar{\mathbf{q}}^0 = \mathbf{0}$ and, $\forall i \in [\![1,d]\!], \bar{\mathbf{q}}^i$ is the *i*-th column of the upper triangular matrix \mathbb{R} . So we have

$$\forall i \in [[0,d]], \ \bar{\mathbf{q}}^i \in \operatorname{Vect}(\boldsymbol{e}^1, \dots, \boldsymbol{e}^d)$$

where $\{e^1, \ldots, e^n\}$ is the standard basis of \mathbb{R}^n . The set $\{\bar{\mathbf{q}}^0, \ldots, \bar{\mathbf{q}}^d\}$ are the vertices of the d-simplex $\bar{K} = \bar{\mathcal{F}}(K)$ and we deduce

$$\bar{K} \subset \operatorname{Vect}(\boldsymbol{e}^1, \dots, \boldsymbol{e}^{\mathrm{d}}).$$
 (27)

By change of variables, we obtain

$$\int_{K} f(\mathbf{q}) d\mathbf{q} = \int_{\bar{K}} f \circ \bar{\mathcal{F}}^{-1}(\bar{\mathbf{q}}) |\det(\mathcal{J}_{\bar{\mathcal{F}}^{-1}}(\bar{\mathbf{q}}))| d\bar{\mathbf{q}}$$

where $\mathcal{J}_{\bar{\mathcal{F}}^{-1}}$ is the Jacobian matrix of $\bar{\mathcal{F}}^{-1}$. From (26), we have $\mathcal{J}_{\bar{\mathcal{F}}^{-1}}(\bar{\mathbf{q}}) = \mathbb{Q}$ and as \mathbb{Q} is an orthogonal matrix, $\det(\mathcal{J}_{\bar{\mathcal{F}}^{-1}}(\bar{\mathbf{q}})) = 1$. So we obtain

$$\int_{K} f(\mathbf{q}) d\mathbf{q} = \int_{\bar{K}} f \circ \bar{\mathcal{F}}^{-1}(\bar{\mathbf{q}}) d\bar{\mathbf{q}}.$$
(28)

Let $\mathbb{P} \in \mathcal{M}_{d,n}(\mathbb{R})$ defined by

$$\mathbb{P} = \left(\begin{array}{c} \mathbb{I}_d \end{array} \middle| \begin{array}{c} \mathbb{O}_{d,n-d} \end{array} \right)$$

and

$$\forall i \in [\![0,d]\!], \quad \mathbf{\bar{\bar{q}}}^i = \mathbb{P}\mathbf{\bar{q}}^i \in \mathbb{R}^d.$$

From (27), we deduce

$$\forall i \in [\![0, \mathbf{d}]\!], \quad \bar{\mathbf{q}}^i = \mathbb{P}^{\mathsf{t}} \bar{\bar{\mathbf{q}}}^i = \begin{pmatrix} -\bar{\bar{\mathbf{q}}}^i \\ -\bar{\bar{\mathbf{0}}}^i \end{pmatrix}.$$

Let $\bar{g} = f \circ \bar{\mathcal{F}}^{-1}$ and $\bar{\bar{K}}$ be the d-simplex in \mathbb{R}^d with vertices $\bar{\mathbf{q}}^i$, $i \in [0, d]$. We denote by $\mathcal{P} : \bar{\bar{K}} \subset \mathbb{R}^d \longrightarrow \bar{K} \subset \mathbb{R}^n$ the application defined by $\mathcal{P}(\bar{\mathbf{q}}) = \mathbb{P}^t \bar{\mathbf{q}}$. We denote by $\bar{g} : \bar{\bar{K}} \longrightarrow \mathbb{R}$ the application defined by

$$\bar{\bar{g}}(\bar{\bar{\mathbf{q}}}) = \bar{g} \circ \mathcal{P}(\bar{\bar{\mathbf{q}}}) = \bar{g}\left(-\frac{\bar{\bar{\mathbf{q}}}}{\bar{\mathbf{0}}}\right).$$

So we obtain

$$\int_{\bar{K}} \bar{g}(\bar{\mathbf{q}}) d\bar{\mathbf{q}} = \int_{\bar{K}} \bar{\bar{g}}(\bar{\mathbf{q}}) d\bar{\bar{\mathbf{q}}}$$
(29)

Let $\overline{\bar{\mathbb{A}}} \in \mathcal{M}_d(\mathbb{R})$ be the matrix defined by

$$\bar{\bar{\mathbb{A}}} = \left(\left. \bar{\bar{\mathbf{q}}}^1 - \bar{\bar{\mathbf{q}}}^0 \right| \left. \bar{\bar{\mathbf{q}}}^2 - \bar{\bar{\mathbf{q}}}^0 \right| \cdots \left| \bar{\bar{\mathbf{q}}}^d - \bar{\bar{\mathbf{q}}}^0 \right| \right).$$
(30)

We can remark that

 $\overline{\mathbb{A}} = \mathbb{P}\overline{\mathbb{A}} \text{ and } \overline{\mathbb{A}} = \mathbb{P}^{t}\overline{\mathbb{A}}.$

Let $\bar{\bar{\mathcal{F}}}:\hat{K}\longrightarrow\bar{\bar{K}}$ the bijective function defined by

$$\bar{\bar{\mathcal{F}}}(\hat{\mathbf{q}}) = \bar{\bar{\mathbb{A}}}\hat{\mathbf{q}} + \bar{\bar{\mathbf{q}}}^0$$

We can now apply the classical change of variables

$$\begin{split} \int_{\bar{K}} \bar{\bar{g}}(\bar{\mathbf{q}}) d\bar{\mathbf{q}} &= \int_{\hat{K}} \bar{\bar{g}} \circ \bar{\bar{\mathcal{F}}}(\hat{\mathbf{q}}) |\det(\mathcal{J}_{\bar{\mathcal{F}}}(\hat{\mathbf{q}}))| d\hat{\mathbf{q}} \\ &= |\det(\bar{\mathbb{A}})| \int_{\hat{K}} \bar{\bar{g}} \circ \bar{\bar{\mathcal{F}}}(\hat{\mathbf{q}}) d\hat{\mathbf{q}} \end{split}$$

To resume from (22) and (29), we have

$$\int_{K} f(\mathbf{q}) d\mathbf{q} = |\det(\bar{\mathbb{A}})| \int_{\hat{K}} \bar{\bar{g}} \circ \bar{\bar{\mathcal{F}}}(\hat{\mathbf{q}}) d\hat{\mathbf{q}}$$
(31)

We can note that

$$\bar{\bar{g}}\circ\bar{\bar{\mathcal{F}}}=f\circ\bar{\mathcal{F}}^{-1}\circ\mathcal{P}\circ\bar{\bar{\mathcal{F}}}$$

Let $\mathcal{F}_K = \bar{\mathcal{F}}^{-1} \circ \mathcal{P} \circ \bar{\bar{\mathcal{F}}}$, we have as expected $\mathcal{F}_{K'}(\hat{\mathbf{a}}) = \bar{\mathcal{F}}^{-1} \circ \mathcal{P} \circ \bar{\bar{\mathcal{F}}}(\hat{\mathbf{a}})$

$$\mathcal{F}_{K}(\mathbf{q}) = \mathcal{F}^{-1} \circ \mathcal{P} \circ \mathcal{F}(\mathbf{q})$$
$$= \bar{\mathcal{F}}^{-1}(\mathbb{P}^{t}(\bar{\mathbb{A}}\hat{\mathbf{q}}))$$
$$= \bar{\mathcal{F}}^{-1}(\bar{\mathbb{A}}\hat{\mathbf{q}})$$
$$= \mathbb{Q}\bar{\mathbb{A}}\hat{\mathbf{q}} + \mathbf{q}^{0}$$
$$= \mathbb{A}_{K}\hat{\mathbf{q}} + \mathbf{q}^{0}.$$

and we obtain

$$\int_{K} f(\mathbf{q}) d\mathbf{q} = |\det(\bar{\mathbb{A}})| \int_{\hat{K}} f \circ \mathcal{F}_{K}(\hat{\mathbf{q}}) d\hat{\mathbf{q}}$$
(32)

To obtain formula (22), it remains to prove that $|\det(\bar{\bar{\mathbb{A}}})| = |\det(\mathbb{A}_K^t \mathbb{A}_K)|^{1/2}$. We have

$$\mathbb{A}_{K}^{t} \mathbb{A}_{K} = \mathbb{A}_{K}^{t} \mathbb{Q} \mathbb{Q}^{t} \mathbb{A}_{K}$$
 as $\mathbb{A}_{K} = \mathbb{Q} \mathbb{A}$

$$= \overline{\mathbb{A}}^{t} \overline{\mathbb{A}}$$
 as \mathbb{Q} is an orthogonal matrix

$$= \overline{\mathbb{A}}^{t} \mathbb{P} \mathbb{P}^{t} \overline{\mathbb{A}}$$
 as $\overline{\mathbb{A}} = \mathbb{P}^{t} \overline{\mathbb{A}}$

$$= \overline{\mathbb{A}}^{t} \overline{\mathbb{A}}$$
 as $\mathbb{P} \mathbb{P}^{t} = \mathbb{I}_{d}$

As $\overline{\bar{\mathbb{A}}}$ is a square matrix, we have $\det(\overline{\bar{\mathbb{A}}}^t\overline{\bar{\mathbb{A}}}) = \det(\overline{\bar{\mathbb{A}}})^2$ and thus

$$|\det(\bar{\mathbb{A}})| = |\det(\mathbb{A}_K^{t}\mathbb{A}_K)|^{1/2}.$$

3.3Volume of a d-simplex

The volume/measure of the d-simplex $K \subset \mathbb{R}^n$ is given by

$$|K| = \int_{K} 1 d\mathbf{q} \tag{33}$$

Using formula (22) with $f \equiv 1$ gives

$$|K| = \left|\det(\mathbb{A}_{K}^{\mathsf{t}}\mathbb{A}_{K})\right|^{1/2} \int_{\hat{K}} 1d\hat{\mathbf{q}} = \left|\det(\mathbb{A}_{K}^{\mathsf{t}}\mathbb{A}_{K})\right|^{1/2} |\hat{K}|$$

From (9), we finally obtain

$$|K| = \frac{\left|\det\left(\mathbb{A}_{K}^{\mathsf{t}}\mathbb{A}_{K}\right)\right|^{1/2}}{\mathrm{d}!}.$$
(34)

In [5] this formula is proved with geometrical arguments. We can also remark that if $d \equiv n$ then \mathbb{A}_K is a square matrix and we obtain the classical formula

$$K| = \frac{|\det(\mathbb{A}_K)|}{\mathrm{d}!}.$$
(35)

3.4Magic formula

In this section an exact computation of the integral over a d-simplex $K \subset \mathbb{R}^n$ for products of power of its barycentric coordinates given by (1) is proved by using previous results obtained by Laplace transforms. Using formula (22) with $f(\mathbf{q}) = \prod_{i=0}^{d} \lambda_i^{\nu_i}(\mathbf{q})$ gives

$$\int_{K} \prod_{i=0}^{d} \lambda_{i}^{\nu_{i}}(\mathbf{q}) d\mathbf{q} = \left| \det(\mathbb{A}_{K}^{t} \mathbb{A}_{K}) \right|^{1/2} \int_{\hat{K}} \prod_{i=0}^{d} (\lambda_{i} \circ \mathcal{F}_{K}(\hat{\mathbf{q}}))^{\nu_{i}} d\hat{\mathbf{q}}$$

From (18) and (34), we obtain

$$\int_{K} \prod_{i=0}^{\mathbf{d}} \lambda_{i}^{\nu_{i}}(\mathbf{q}) d\mathbf{q} = d! |K| \int_{\hat{K}} \prod_{i=0}^{\mathbf{d}} \hat{\lambda}_{i}^{\nu_{i}}(\hat{\mathbf{q}}) d\hat{\mathbf{q}}$$

Using formula (10) gives (1)

References

- [1] Douglas N. Arnold, Richard S. Falk, and Ragnar Winther. Geometric decompositions and local bases for spaces of finite element differential forms. Computer Methods in Applied Mechanics and Engineering, 198(21):1660 – 1672, 2009. Advances in Simulation-Based Engineering Sciences – Honoring J. Tinsley Oden.
- [2] Susanne C. Brenner and L. Ridgway Scott. The Mathematical Theory of Finite Element Methods. Texts in Applied Mathematics. Springer, New York, Berlin, Paris, 2008.

- [3] M.A. Eisenberg and L. E. Malvern. On finite element integration in natural co-ordinates. Int. J. Numer. Methods Eng., 7:574–575, 1973.
- [4] P. Frey. Chapter 7. the finite element approximation. https://www.ljll.math.upmc.fr/ frey/ma691.html, 2008.
- [5] A.J. Hanson. Geometry for n-Dimensional Graphics, 1994.
- [6] E T Jaynes. Probability Theory The Logic Of Science Edition 1. Cambridge University Press, 2003.
- [7] Alfio Quarteroni. Numerical models for differential problems. Springer, Milan New York, 2009.
- [8] Joel L. Schiff. The Laplace Transform. Springer New York, 1999.
- [9] F.J. Vermolen and A. Segal. On an integration rule for products of barycentric coordinates over simplexes in rn. *Journal of Computational and Applied Mathematics*, 330:289 – 294, 2018.