

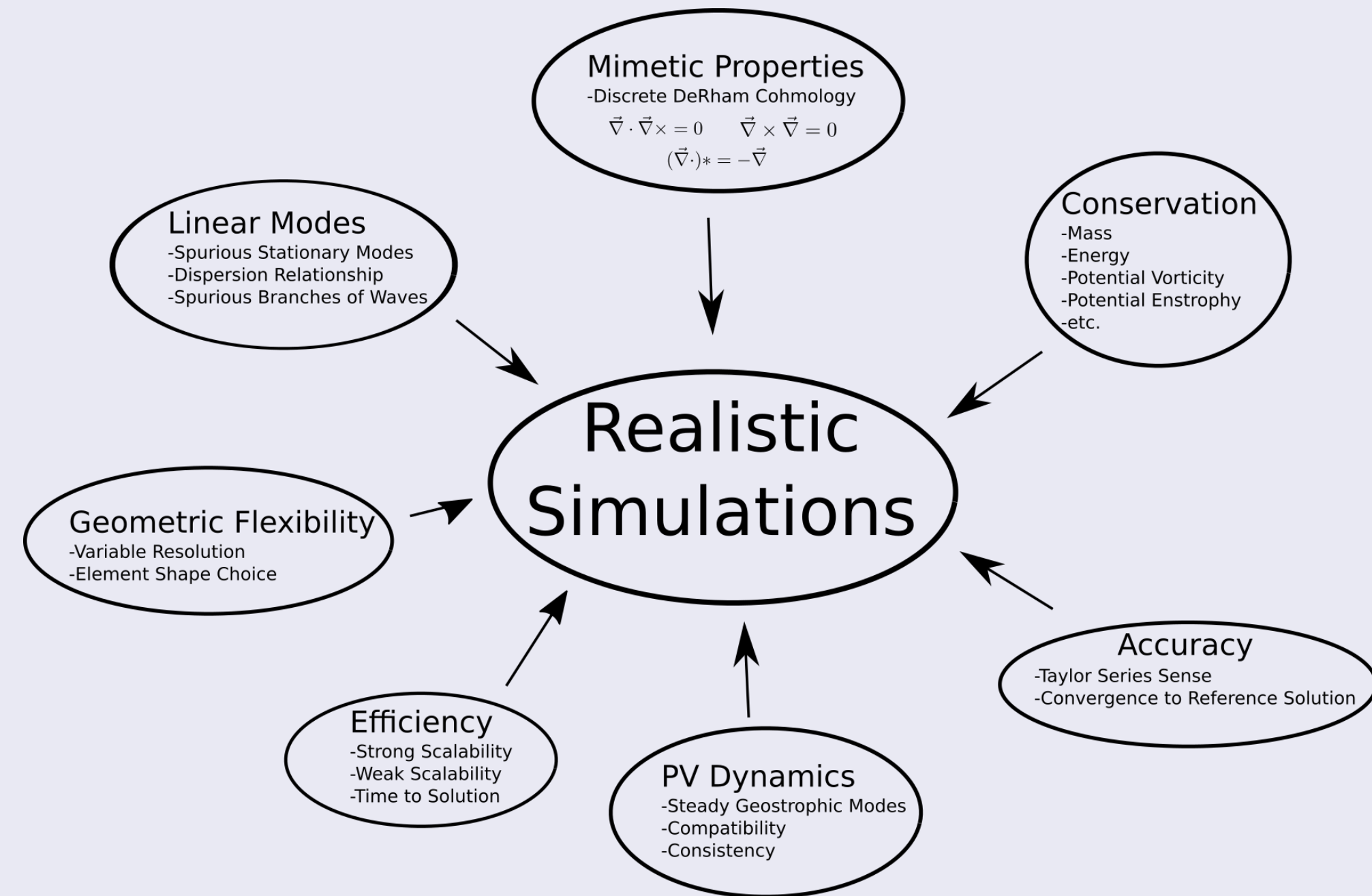
Dynamico-FE: A Hydrostatic Dynamical Core using Higher-Order Structure-Preserving Finite Elements



Chris Eldred¹, Thomas Dubos² and Evaggelos Kritsikis¹
¹LAGA, University of Paris 13 and ²LMD, Ecole Polytechnique
 Contact: chris.eldred@gmail.com
 Paper Number: A31A-0006



(I) Desirable Model Properties



How do we get these properties?

(II) Hamiltonian Formulation in a General Vertical Coordinate

Prognose (μ, S, \vec{v}) using

$$\frac{\partial \mu}{\partial t} + \vec{\nabla} \cdot \frac{\delta \mathcal{H}}{\delta \vec{v}} + \partial_\eta(W) = 0$$

$$\frac{\partial S}{\partial t} + \vec{\nabla} \cdot (s \frac{\delta \mathcal{H}}{\delta \vec{v}}) + \partial_\eta(sW) = 0$$

$$\frac{\partial \vec{v}}{\partial t} + \frac{\vec{\nabla} \times \vec{v}}{\mu} \times \frac{\delta \mathcal{H}}{\delta \vec{v}} + \vec{\nabla} \frac{\delta \mathcal{H}}{\delta \mu} + s \vec{\nabla} \frac{\delta \mathcal{H}}{\delta S} + \frac{W}{\mu} \partial_\eta(\vec{v}) = 0$$

with Hamiltonian \mathcal{H}

$$\mathcal{H} = \mathcal{H}[\mu, \vec{v}, S, z] = \int \mu K(\vec{v}, z) + \mu U(\frac{1}{\mu} \frac{\partial z}{\partial \eta}, \frac{S}{\mu}) + \mu \Phi(z)$$

$$\frac{\delta \mathcal{H}}{\delta \vec{v}} = \mu \vec{u}$$

$$\frac{\delta \mathcal{H}}{\delta \mu} = K + \Phi + U + \mu \frac{\partial U}{\partial \mu}$$

$$\frac{\delta \mathcal{H}}{\delta S} = \mu \frac{\partial U}{\partial S}$$

and diagnose z using

$$\frac{\delta \mathcal{H}}{\delta z} = \mu \frac{\partial K}{\partial z} + \mu \frac{\partial \Phi}{\partial z} - \frac{\partial}{\partial \eta} (\mu \frac{\partial U}{\partial (\frac{\partial z}{\partial \eta})}) = 0 \text{ (Hydrostatic balance)}$$

where $\vec{v} = \vec{u} + \vec{R}(z)$ is the horizontal (covariant) absolute velocity, $\mu = \frac{1}{\alpha} \frac{\partial z}{\partial \eta}$ is the pseudo-density, z the height and $S = \mu s$ the mass-weighted entropy. These equations work for a deep, non-spherical atmosphere with an arbitrary equation of state $U(\alpha, s)$. Lagrangian vertical coordinate defined by $W = 0$. Mass-based vertical coordinate prognoses M_s and specifies $M(\eta) = \int \mu = a(\eta)M_s + b(\eta)M_0 \rightarrow \mu$ is no longer prognostic (it gives the equation for W) \rightarrow redefine $\frac{\delta \mathcal{H}}{\delta \mu}$ to conserve energy (vertical remapping invariance). Details are in [3].

(III) General Mimetic Discretizations: Primal deRham Complex

- Select any 1D Spaces \mathcal{A} and \mathcal{B} such that: $\mathcal{A} \xrightarrow{\frac{d}{dx}} \mathcal{B}$
- Use tensor products to extend to n-dimensions
- Our (novel) choices of \mathcal{A} and \mathcal{B} are guided by linear mode properties and coupling to physics/tracer transport

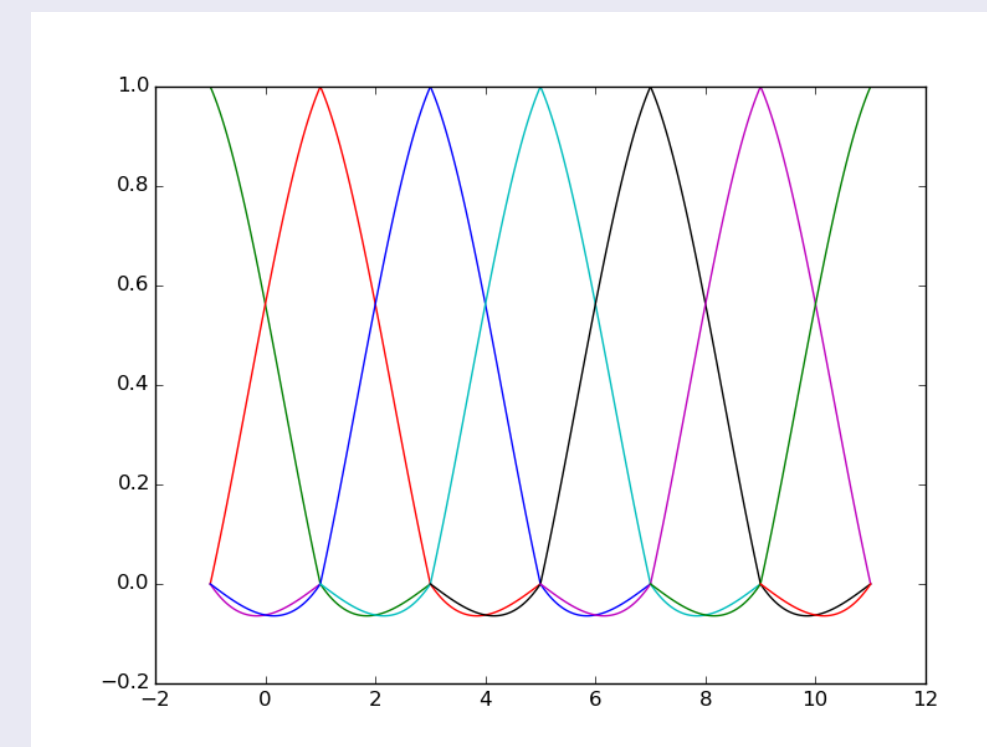
$$\begin{array}{ccccc} \mathbb{W}_0 & \xrightarrow{\frac{d}{dx}} & \mathbb{W}_1 & \xrightarrow{\frac{d}{dx}} & \mathbb{W}_2 & \xrightarrow{\frac{d}{dx}} & \mathbb{W}_3 \\ \downarrow \vec{\nabla} \cdot & & \downarrow \vec{\nabla} \times & & \downarrow \vec{\nabla} \cdot & & \\ \delta & & \delta & & \delta & & \\ (da^k, b^{k+1}) & & (a^k, \delta b^{k+1}) & & & & \end{array}$$

Primal deRham Complex

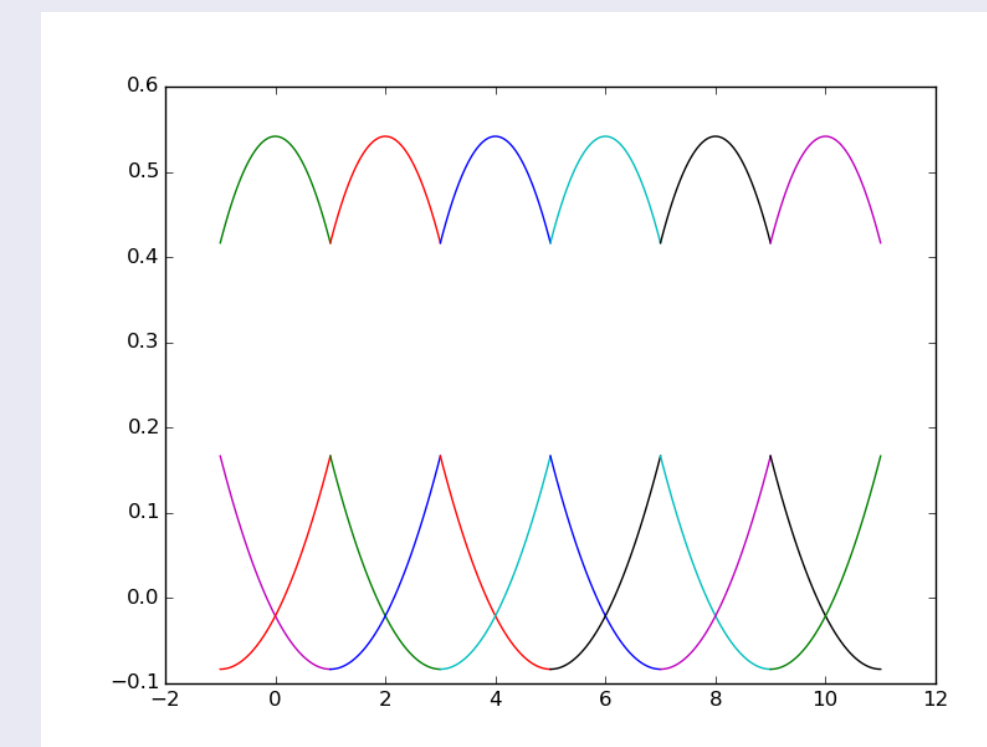
Integration by parts plus $\delta = *d*$ implicitly defines *

(IV) Mimetic Galerkin Differences

H^1 space defined following [1], with L_2 defined to be compatible following [4]. This is an arbitrary order extension of [2]. For 3rd order gives:



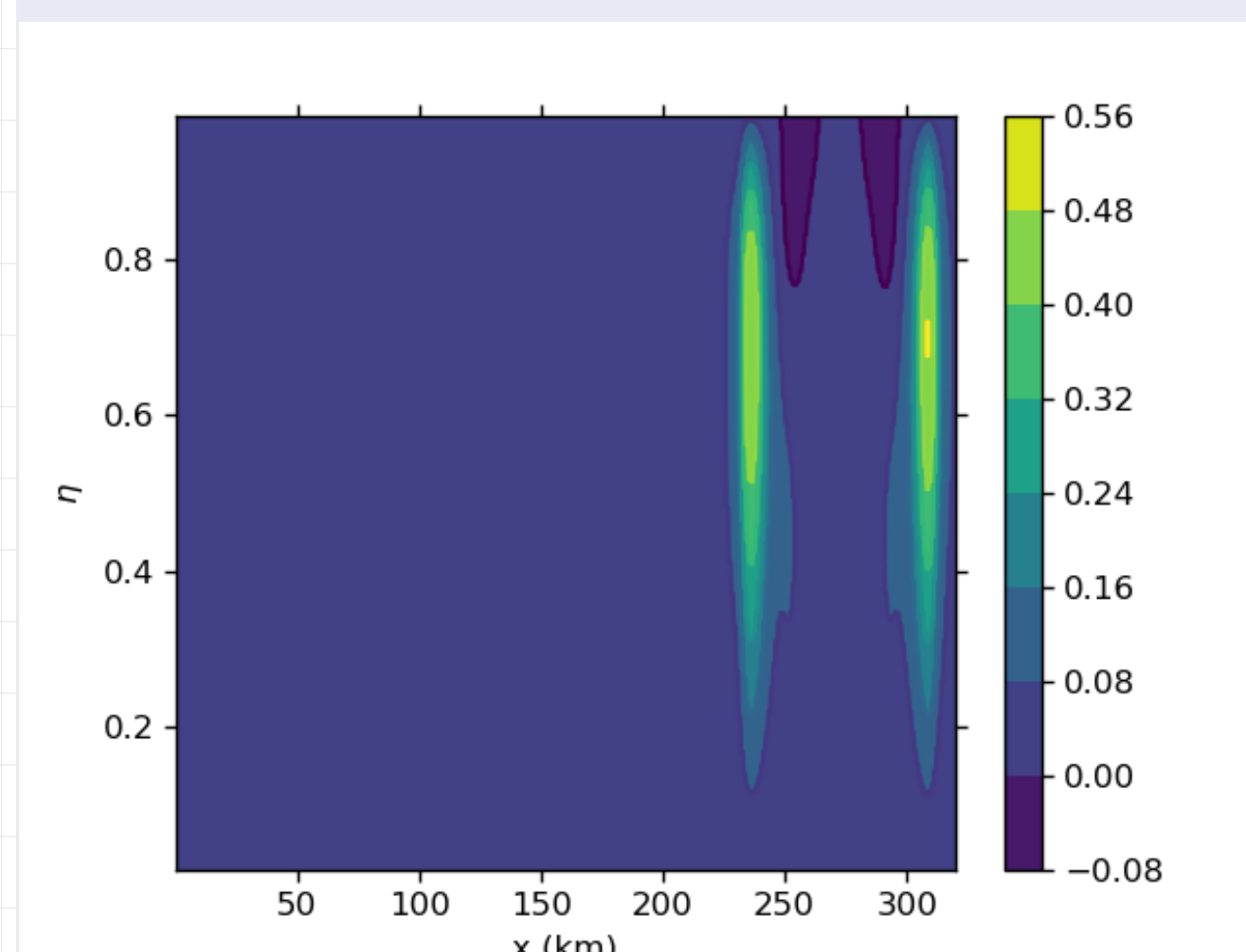
$\mathcal{A} = H_1$ Space (1D)



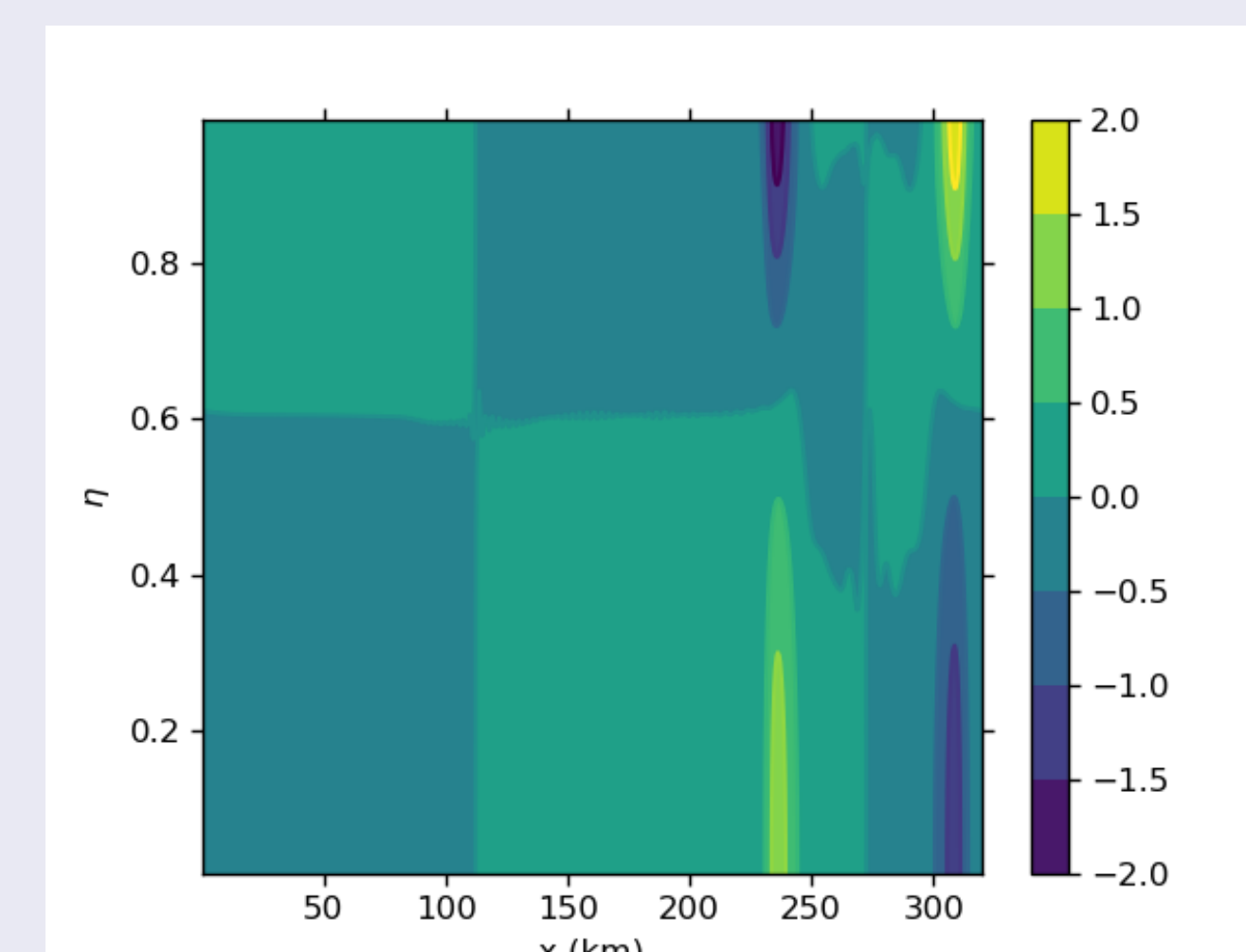
$\mathcal{B} = L_2$ Space (1D)

Single degree of freedom per geometric entity with higher order through larger stencils \rightarrow **no spectral gaps, easy coupling to physics/tracer transport, less local**

(VI) Preliminary Results



θ'

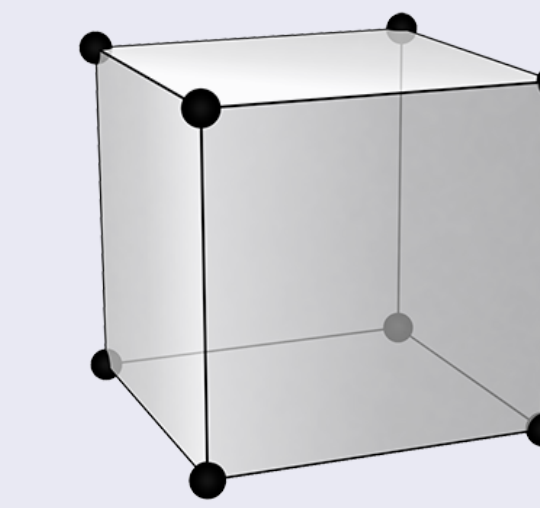


u'

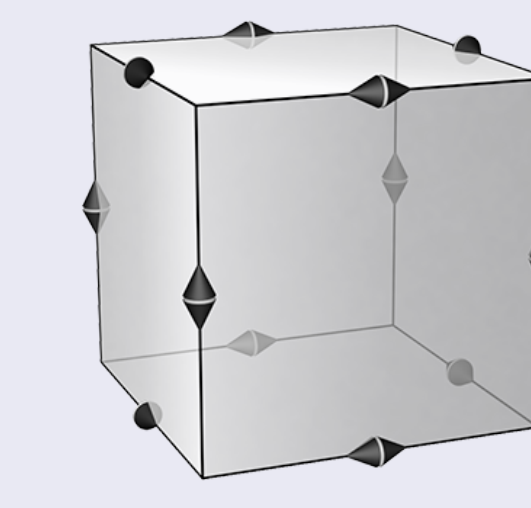
DCMIP 3.1 (Gravity Wave) adapted to slice. $\Delta x = 1km$, 30 vertical levels with a Lagrangian vertical coordinate, $u_0 = 100ms^{-1}$. 4th-order Runge Kutta time stepping. Shown at $t = 1125s$. Prognoses Θ .

(V) Compatible Spaces and Discretization

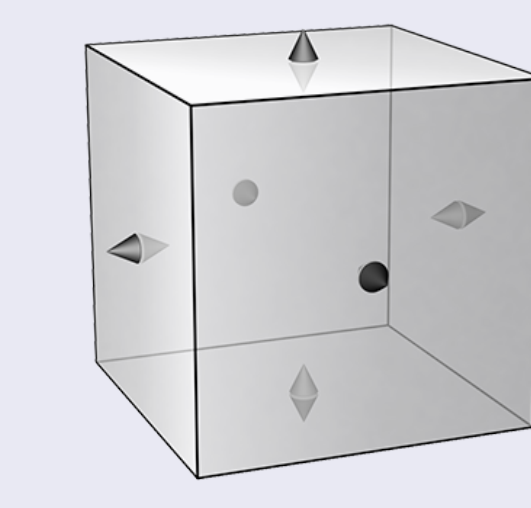
Restrict variables to finite dimensional subspaces, and multiply by test functions to obtain to obtain discrete weak form equations. Choose spaces and staggering as:



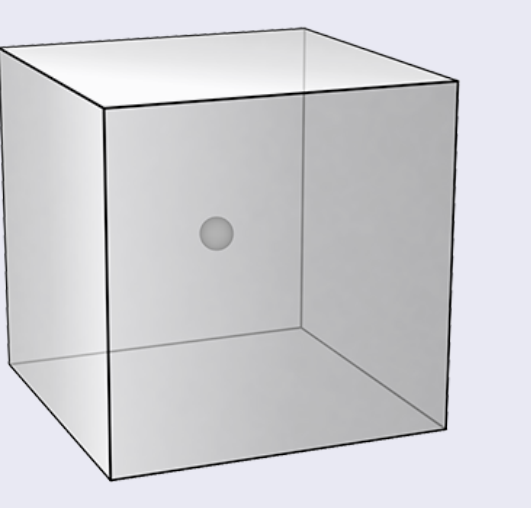
$\psi \in H^1$



$\vec{\zeta} \in H(curl)$



$\vec{u}, \vec{v} \in H(div)$



$\mu, \delta, \chi, S \in L_2$

This approach gives a (quasi-)Hamiltonian semi-discretization that conserves mass, entropy and total energy, independent of the choice of specific spaces!

(VII) Conclusions

- 1 Hydrostatic dynamical core can run with mass-based or Lagrangian vertical coordinates, prognosing Θ or S , on slice or planar grid.
- 2 Possible to obtain most of the desirable properties by combining **mimetic discretization** with a **Hamiltonian formulation**
- 3 Hamiltonian formulation enables treatment of many approximations (deep/shallow, traditional/non-traditional, spherical/non-spherical) with a **unified framework**

(VIII) Future and Ongoing Work

- 1 Extension to the cubed sphere grid
- 2 Corrected $\frac{\delta \mathcal{H}}{\delta \mu}$ for mass-based vertical coordinate
- 3 Extending conservation to time integration (Poisson integrators)
- 4 Introduction of moisture in a consistent way (Multicomponent formulation)
- 5 Incorporation of dissipation and parameterization in a consistent way (Metriplectic formulation, stochastic pdes, higher order moments)
- 6 Computational efficiency
- 7 Extension to non-hydrostatic or semi-hydrostatic (sound-proof) equations

(IX) Themis: Accelerated Computational Science

Themis is a PETSc-based software framework for parallel, high-performance discretization of variational forms through mimetic, tensor-product Galerkin methods.

Available online at https://bitbucket.org/chris_eldred/themis



References

- [1] J.W. Banks, T. Hagstrom. On Galerkin difference methods, Journal of Computational Physics, May 2016
- [2] E. Kritsikis and T. Dubos. Higher-order finite elements for the shallow-water equations on the cubed sphere, PDEs on the Sphere workshop, April 2014
- [3] T. Dubos and M. Tort. Equations of Atmospheric Motion in Non-Eulerian Vertical Coordinates: Vector-Invariant Form and Quasi-Hamiltonian Formulation, Monthly Weather Review, June 2014
- [4] R.R. Hiemstra, D. Toshniwal, R.H.M. Huijsmans, M.I. Gerritsma. High order geometric methods with exact conservation properties, Journal of Computational Physics, January 2014