Extension of the 1981 Arakawa and Lamb Scheme to Arbitrary Grids

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Introduction

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Key Papers

Arakawa and Lamb 1981

A Potential Enstrophy and Energy Conserving Scheme for the Shallow Water Equations

AKIO ARAKAWA AND VIVIAN R. LAMB¹

Salmon 2004

Poisson-Bracket Approach to the Construction of Energy- and Potential-Enstrophy-Conserving Algorithms for the Shallow-Water Equations

RICK SALMON

Thuburn, Cotter and Dubos 2012

A FRAMEWORK FOR MIMETIC DISCRETIZATION OF THE ROTATING SHALLOW-WATER EQUATIONS ON ARBITRARY POLYGONAL GRIDS*

J. THUBURN[†] AND C. J. COTTER[‡]

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Broad Research Overview: Dynamical Cores

- Develop discrete models of the atmosphere
- ② Dynamical core: deals with "resolved processes"
- O Parameterizations: deal with "unresolved processes"
- Model: dynamical core + parameterizations



Key Principles of Numerical Modeling

- In the solving arbitrary PDEs- physical system
- O No analytic solutions
- $\textcircled{O} Differential equations \rightarrow algebraic equations$
- O algebraic solutions have the same properties as the differential solutions?



Philosophy of Dynamical Core Design

Models should respect underlying physics

- Energetics
- PV dynamics
- 3 Wave dynamics
- Onserved quantities- energy, potential enstrophy
- Mimetic properties

$$\vec{\nabla}\times\vec{\nabla}\phi=\mathbf{0}$$



Shallow Water Equations

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Vector Invariant Shallow Water Equations

$$\frac{\partial h}{\partial t} + \vec{\nabla} \cdot (h\vec{u}) = 0$$

$$\frac{\partial \vec{u}}{\partial t} + q\hat{k} \times (h\vec{u}) + \vec{\nabla}\Phi = 0$$

$$\mathbb{J} = \begin{pmatrix} 0 & -\vec{\nabla} \cdot \\ -\vec{\nabla} & -q\hat{k} \times \end{pmatrix}$$

$$\mathcal{H} = \frac{1}{2}g(h,h) + \frac{1}{2}(\vec{F},\vec{u})$$

$$\frac{\partial \vec{x}}{\partial t} = \mathbb{J}\frac{\delta \mathcal{H}}{\delta \vec{x}}$$

$$\frac{\delta \mathcal{H}}{\delta \vec{x}} = \begin{pmatrix} \Phi \\ h\vec{u} \end{pmatrix}$$

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Shallow Water Equations: (Subset of) Properties

Mimetic

No Spurious Vorticity Production:

$$\vec{\nabla} \times \vec{\nabla} \phi = \mathbf{0}$$

Pressure Gradient is Energy Conserving:

$$(\vec{\nabla})^* = -\vec{\nabla} \cdot$$

Conserved Quantities

Total Energy
$$\int_{\Omega} \frac{gh^2}{2} + \frac{h|\vec{u}|^2}{2}$$

Potential Enstrophy $\int_{\Omega} h \frac{q^2}{2}$

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Arakawa and Lamb 1981 Scheme: Desirable Properties

(A Subset of) Desirable Properties

- No spurious vorticity production (curl-free gradient)
- Energy-conserving pressure gradient force (divergence and gradient are adjoints)
- Total energy and potential enstrophy conservation



Arakawa and Lamb 1981 Scheme: Limitations

(A Subset of) AL81 Shortcomings

Restricted to logically square, orthogonal grids

TRiSK: Ringler, Skamarok, Klemp, Thuburn, Cotter, Dubos

- General, non-orthogonal polygonal grids
- Choose between total energy and potential enstrophy conservation



Logically square, orthogonal grid

General, non-orthogonal grid

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Recap: What am I trying to accomplish?

- How can Arakawa and Lamb 1981 be extended to arbitrary, non-orthogonal polygonal grids?
- In a way that preserves all of its desirable properties, and does not add new limitations?



AL81 on arbitrary grids

Extension of AL81 to Arbitrary Grids

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Mimetic Methods (Discrete Exterior Calculus)

Mimetic Operators

$$ec{
abla} imesec{
abla}\phi=0 \Longleftrightarrow ar{D_2}ar{D_1}=0 \ (ec{
abla})^*=-ec{
abla}\cdot \Longleftrightarrow D_2=-ar{D_1}^T$$



(Discrete) Exterior Derivative



Primal-Dual Grid

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Conservation Laws (Hamiltonian Mechanics)

Hamiltonian (Energy)

$$\mathbb{J}^{\mathcal{T}}=-\mathbb{J}$$

$\ensuremath{\mathcal{H}}$ is positive definite

Casimirs (Potential Enstrophy)

$$\mathbb{J}\frac{\delta \mathcal{Z}}{\delta \vec{x}} = 0$$

Discrete Conservation

Enforce these conditions in discrete case \rightarrow conservation!

Recap: Conservative, Mimetic Methods



Conservative, Mimetic Methods

 Use mimetic operators to build a discrete (quasi-)Hamiltonian system

Onifies two important lines of research

Generalized C Grid Discretization: Hamiltonian-DEC

- Discrete variables are $m_i = \int h dA$ and $u_e = \int \vec{u} \cdot \vec{dl}$
- C grid staggering (*m_i* at cell centers, *u_e* at edges)



Q operator is the remaining hurdle

General Form of **Q**



Diagram of Q operator stencil

Discrete Conservation

Energy

$$\mathbb{J}^{\mathsf{T}} = -\mathbb{J} \longrightarrow \mathbf{Q} = -\mathbf{Q}^{\mathsf{T}} \longrightarrow \alpha_{e,e',v} = -\alpha_{e',e,v}$$

Potential Enstrophy

$$\mathbb{J}\frac{\delta \mathcal{Z}}{\delta \vec{x}} = 0 \longrightarrow \text{ linear system of equations} \longrightarrow \mathbf{A}\vec{\alpha} = \vec{b}$$

Also want $\mathbf{Q} \rightarrow \mathbf{W}$ when $q_v = \text{const}$ (steady geostrophic modes)

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Solving $\mathbf{A}\vec{\alpha}=\vec{b}$

Issue: System is too large

- Geodesic grid: 90 coefficients per cell, all coefficients are interdependent \rightarrow not feasible for realistic grids
- Cubed sphere grid is similar (24 coefficients per cell)

Solution: Subystem Splitting

$$\mathbf{A}\vec{lpha} = \vec{b} \longrightarrow \sum_{i} \mathbf{A}_{i}\vec{lpha}_{i} = \vec{b}_{i}$$

Split into independent subsystems for each cell!

System has been solved for various planar and spherical grids

Recap: What have I accomplished?

What has been done?

- Arakawa and Lamb 1981 extended to arbitrary grids via new **Q**
- Coefficients can be precomputed (efficiently)





Test Case Results

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Variants of **Q** Operator

Total Energy

$$\mathbf{Q}F_e = rac{1}{2}q_e\mathbf{W}F_e + rac{1}{2}\mathbf{W}q_eF_e$$

Works for ANY choice of q_e (APVM, CLUST, etc.)

Potential Enstrophy

$$\mathbf{Q}F_{\mathrm{e}}=q_{\mathrm{e}}\mathbf{W}F_{\mathrm{e}}$$

Requires that q_e is arithmetic mean

Total Energy and Potential Enstrophy

$$\mathbf{Q}F_e = Q(q_v,F_e)$$

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Grids

Grids

- Cubed Sphere: 6x384x384, 880K cells (25km resolution)
- Geodesic: G8, 640K cells (30km resolution)



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Operators

Operators

$$\mathbf{R} = \sum_{v \in VC(i)} \frac{A_{iv}}{A_i}$$

$$\mathbf{\Phi} = \frac{A_{ie}}{A_e} + \frac{A_{je}}{A_e}$$

$$\mathbf{I} = \frac{1}{A_i}$$
$$\mathbf{J} = \frac{1}{A_v}$$

$$\mathbf{H} = \frac{le}{de}$$

 ${\bf H}$ different for cubed-sphere

Time Stepping

- Adams-Bashford 3rd Order (15s cubed-sphere, 22.5s geodesic)
- 10 simulated days, output every 6 hours

Dissipation

- $\vec{\nabla}^2 \vec{u}$ specified in Galewsky et. al added to that test case
- Schemes stable without it

Galewsky et. al (Unstable Jet)- Geodesic (Tweaked)







Galewsky et. al (Unstable Jet)- Cubed Sphere (Thuburn)







Order of Accuracy- Laplacian on Primal



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Order of Accuracy- Laplacian on Dual



Computed for $\psi = \cos(\theta) \sin(\lambda)$

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Order of Accuracy- \mathbf{Q} (also \mathbf{W})

 L_2 Error

L_{∞} Error



Order of Accuracy- Zonal Jet (Williamson Test Case 2)



Height errors

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Conclusions

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Summary and Conclusions

Conclusions

- Preserves all desirable properties of AL81

Future Work

- Consistent W (and Q)
- Anticipated Potential Vorticity + other dissipation options
- **③** Cubed sphere accuracy- Hodge stars, grid optimization
- In Effects of extra Rossby modes on geodesic grid?
- Somparison to Z-grid scheme (Salmon 2007)