

Higher-order Structure-Preserving Finite Elements for Atmospheric Dynamical Cores

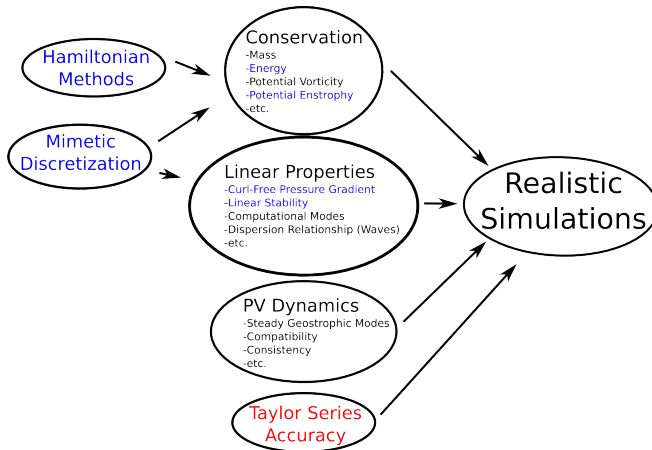
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Introduction

Elements of Dynamical Core Design



History of Structure-Preserving Schemes for Atmospheric Models

- The Beginning: Arakawa and Lamb 1981, Sadourney 1975
- Mimetic Finite Differences: Ringler et. al 2010; Thuburn et. al 2012, 2014, many others
- Mimetic Finite Elements: Cotter et. al 2012,201; McRae et. al 2014, many others
- Hamiltonian: Salmon 2004,2005,2007; Sommer+Nevir 2009; Gassmann 2008,2013; Dubos et. al 2015; Tort et. al 2015, many others
- Fundamentally, structure-preserving schemes can be viewed as a combination of a **mimetic discretization method** plus a **hamiltonian formulation**

What is structure-preservation?

- 1 Discrete deRham cohomology (mimetic discretization)

$$\vec{\nabla} \times \vec{\nabla} = 0$$

$$\vec{\nabla} \cdot \vec{\nabla} \times = 0$$

- 2 Quasi-Hamiltonian system (conserves total energy, and possibly other invariants)

$$\frac{d\mathcal{H}}{dt} = 0$$

$$\frac{d\mathcal{C}}{dt} = 0$$

General Formulation for Mimetic Discretizations: Primal-Dual Double deRham Complex

$$\begin{array}{ccccccc}
 & & d & & d & & d \\
 & & \vec{\nabla} & & \vec{\nabla} \times & & \vec{\nabla} \cdot \\
 \mathbb{W}_0 & \xrightarrow{\quad} & \mathbb{W}_1 & \xrightarrow{\quad} & \mathbb{W}_2 & \xrightarrow{\quad} & \mathbb{W}_3 \\
 \downarrow * & & \downarrow * & & \downarrow * & & \downarrow * \\
 \mathbb{W}^3 & \xleftarrow{\quad} & \mathbb{W}^2 & \xleftarrow{\quad} & \mathbb{W}^1 & \xleftarrow{\quad} & \mathbb{W}^0 \\
 & & d & & d & & d \\
 & & \vec{\nabla} \cdot & & \vec{\nabla} \times & & \vec{\nabla}
 \end{array}$$

$$\delta = *d*$$

$$\nabla^2 = d\delta + \delta d$$

$$\vec{\nabla} \cdot \vec{\nabla} \times = 0 = \vec{\nabla} \times \vec{\nabla}$$

$$\int_{\Omega} dW = \int_{d\Omega} W$$

$$dd = 0 = \delta\delta$$

General Formulation for Mimetic Discretizations: Primal deRham Complex

$$\begin{array}{ccccccc}
 & d & & d & & d & \\
 & \vec{\nabla} \rightarrow & & \vec{\nabla} \times & & \vec{\nabla} \cdot & \\
 W_0 & \xrightarrow{\quad} & W_1 & \xrightarrow{\quad} & W_2 & \xrightarrow{\quad} & W_3 \\
 & \nwarrow \text{---} & \nwarrow \text{---} & \nwarrow \text{---} & & & \\
 & \vec{\nabla} \cdot & & \vec{\nabla} \times & & \vec{\nabla} & \\
 & \delta & & \delta & & \delta & \\
 & (da^k, b^{k+1}) = (a^k, \delta b^{k+1}) & & & & &
 \end{array}$$

$$\delta = *d*$$

$$\nabla^2 = d\delta + \delta d$$

$$\vec{\nabla} \cdot \vec{\nabla} \times = 0 = \vec{\nabla} \times \vec{\nabla}$$

$$\int_{\Omega} dW = \int_{d\Omega} W$$

$$dd = 0 = \delta\delta$$

Non-Canonical Hamiltonian Dynamics

Poisson Brackets

Evolution of an arbitrary functional $\mathcal{F} = \mathcal{F}[\vec{x}]$ is governed by:

$$\frac{d\mathcal{F}}{dt} = \left\{ \frac{\delta\mathcal{F}}{\delta\vec{x}}, \frac{\delta\mathcal{H}}{\delta\vec{x}} \right\} \quad (1)$$

with Poisson bracket $\{, \}$ antisymmetric (also satisfies Jacobi):

$$\left\{ \frac{\delta\mathcal{F}}{\delta\vec{x}}, \frac{\delta\mathcal{G}}{\delta\vec{x}} \right\} = - \left\{ \frac{\delta\mathcal{G}}{\delta\vec{x}}, \frac{\delta\mathcal{F}}{\delta\vec{x}} \right\} \quad (2)$$

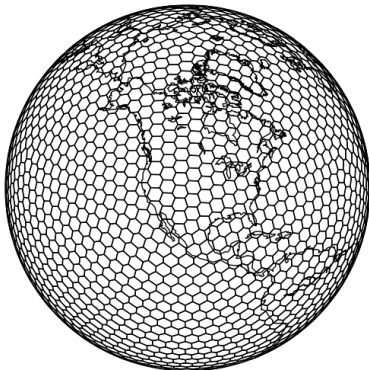
Also have Casimirs \mathcal{C} that satisfy:

$$\left\{ \frac{\delta\mathcal{F}}{\delta\vec{x}}, \frac{\delta\mathcal{C}}{\delta\vec{x}} \right\} = 0 \quad \forall \mathcal{F} \quad (3)$$

Neatly encapsulates conservation properties (\mathcal{H} and \mathcal{C}).

Dynamico

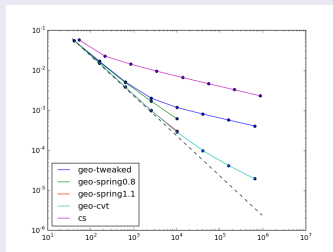
G4 - tweaked Primal



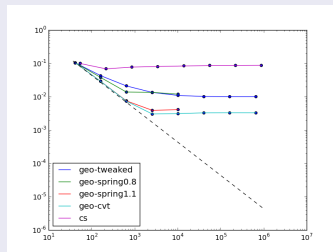
- 1 Primal-Dual: Mimetic finite differences (based on TRiSK scheme): C grid horizontal, Lorenz vertical
- 2 Icosahedral grid
- 3 Hydrostatic primitive equations: Lagrangian and mass-based vertical coordinates
- 4 Conserves mass, energy and entropy
- 5 See Dubos et. al 2015 for more information

Issues in Dynamico

Operator Accuracy



L_2 for W



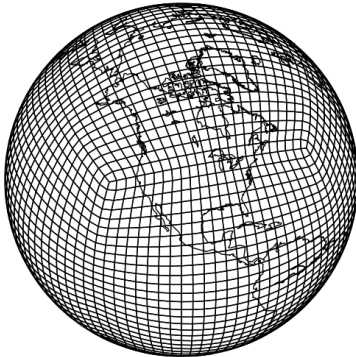
L_∞ for W

Spurious Branches of Dispersion Relationship

Hexagonal grid means 3:1 ratio of wind to mass dofs (should be 2:1) \rightarrow spurious branch of Rossby waves with unphysical behaviour

How do we fix them?

C4 - 6x24x24 - thuburn Primal



- 1 Fix spurious branches:
Quadrilateral (cubed-sphere)
grid, correct 2:1 ratio of dofs
- 2 Fix accuracy: Use Primal
approach (mimetic Galerkin
methods)
- 3 Keep the same equations
and Hamiltonian structure
- 4 Keep the same mimetic and
conservation properties

Mimetic Galerkin Discretization

Mimetic Galerkin Discretization

$$\begin{array}{ccccccc}
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 & \text{---} & & \text{---} & & \text{---} & \\
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$$\int_{\Omega} dW = \int_{d\Omega} W$$

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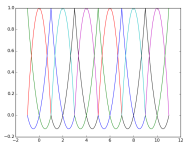
General Approach to Mimetic Galerkin Spaces

Mimetic Spaces

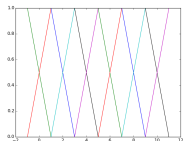
Select 1D Spaces \mathcal{A} and \mathcal{B} such that : $\mathcal{A} \xrightarrow{\frac{d}{dx}} \mathcal{B}$ (4)

- Use tensor products to extend to n-dimensions
- Works for ANY set of spaces \mathcal{A} and \mathcal{B} that satisfy this property (mimetic finite elements use P_n and $P_{DG,n-1}$)
- Mimetic spectral element, Mimetic isogeometric methods (B-splines) all fall under this framework
- Our (different) choices of \mathcal{A} and \mathcal{B} are guided by linear mode properties and coupling to physics/tracer transport
- See Hiemstra et. al 2014 (and references therein)

$P_2 - P_{1,DG}$ Dispersion Relationship



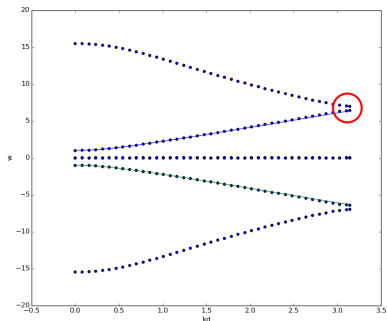
$\mathcal{A} = H_1$ Space



$\mathcal{B} = L_2$ Space

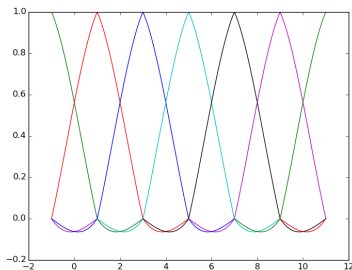
Multiple dofs per element \rightarrow breaks translational invariance \rightarrow spectral gaps

We have developed an alternative: mimetic Galerkin differences

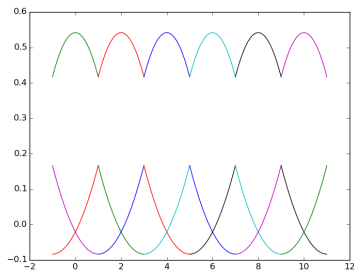


Inertia-Gravity Wave Dispersion Relationship (1D)

Mimetic Galerkin Differences: Basis



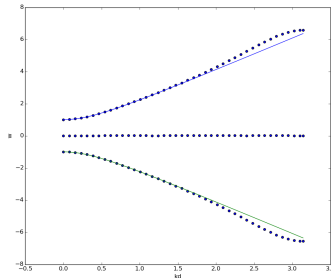
$\mathcal{A} = H_1$ Space



$\mathcal{B} = L_2$ Space

Single degree of freedom per geometric entity (**physics coupling**)
Higher order by larger stencils (**less local**)

Mimetic Galerkin Differences- Dispersion



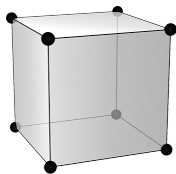
Inertia-Gravity Wave Dispersion Relationship (1D)

Spectral gap is gone

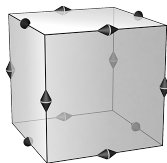
Can show that dispersion relation is $O(2n)$ where n is the order

More details in a forthcoming paper

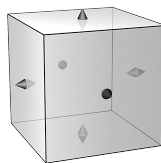
Overview of 3D Spaces



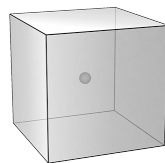
W_0



W_1



W_2



W_3

$$W_0 \xrightarrow{\vec{\nabla}} W_1 \xrightarrow{\vec{\nabla} \times} W_2 \xrightarrow{\vec{\nabla} \cdot} W_3$$

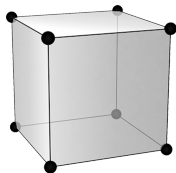
$W_0 = \mathcal{A} \otimes \mathcal{A} \otimes \mathcal{A} = H_1 = \text{Continuous Galerkin}$

$W_1 = (\mathcal{B} \otimes \mathcal{A} \otimes \mathcal{A})\hat{i} + \dots = H(\text{curl}) = \text{Nedelec}$

$W_2 = (\mathcal{A} \otimes \mathcal{B} \otimes \mathcal{B})\hat{i} + \dots = H(\text{div}) = \text{Raviart-Thomas}$

$W_3 = \mathcal{B} \otimes \mathcal{B} \otimes \mathcal{B} = L_2 = \text{Discontinuous Galerkin}$

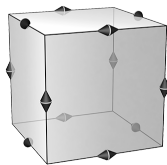
Grid Staggering for HPE



W_0

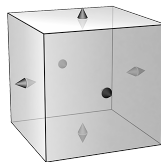
ψ, \ominus

Differential
Geometry



W_1

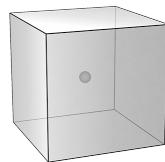
$\vec{\zeta}$



W_2

$\vec{v}, \vec{u}, W, \ominus$

CP Grid



W_3

$\mu, \ominus, \delta, \chi$

L Grid

Follows from differential geometry and Tonti diagram
Galerkin Version of a C Grid

Question: Where should \ominus live?

Hamiltonian Formulation

Poisson Brackets (Lagrangian Vertical Coordinate)

Poisson Brackets

From Dubos and Tort 2014, evolution of $\mathcal{F}[\vec{x}] = \mathcal{F}[\mu, \vec{v}, \Theta, z]$ is

$$\frac{d\mathcal{F}}{dt} = \left\{ \frac{\delta \mathcal{F}}{\delta \vec{x}}, \frac{\delta \mathcal{H}}{\delta \vec{x}} \right\}_{sw} + \left\{ \frac{\delta \mathcal{F}}{\delta \vec{x}}, \frac{\delta \mathcal{H}}{\delta \vec{x}} \right\}_{\Theta} + \left\langle \frac{\delta \mathcal{F}}{\delta z} \frac{\partial z}{\partial t} \right\rangle \quad (5)$$

$$\left\{ \frac{\delta \mathcal{F}}{\delta \vec{x}}, \frac{\delta \mathcal{H}}{\delta \vec{x}} \right\}_{sw} = \left\langle \frac{\delta \mathcal{H}}{\delta \vec{v}} \cdot \vec{\nabla} \frac{\delta \mathcal{F}}{\delta \mu} - \frac{\delta \mathcal{H}}{\delta \vec{v}} \cdot \vec{\nabla} \frac{\delta \mathcal{F}}{\delta \mu} \right\rangle + \left\langle \frac{\vec{\nabla} \times \vec{v}}{\mu} \cdot \left(\frac{\delta \mathcal{F}}{\delta \vec{v}} \times \frac{\delta \mathcal{H}}{\delta \vec{v}} \right) \right\rangle \quad (6)$$

$$\left\{ \frac{\delta \mathcal{F}}{\delta \vec{x}}, \frac{\delta \mathcal{H}}{\delta \vec{x}} \right\}_{\Theta} = \left\langle \theta \left(\frac{\delta \mathcal{H}}{\delta \vec{v}} \cdot \vec{\nabla} \frac{\delta \mathcal{F}}{\delta \Theta} - \frac{\delta \mathcal{H}}{\delta \vec{v}} \cdot \vec{\nabla} \frac{\delta \mathcal{F}}{\delta \Theta} \right) \right\rangle \quad (7)$$

where μ is the pseudo-density, $\vec{v} = \vec{u} - \vec{R}$ is the absolute (covariant) velocity, $\Theta = \theta\mu$ is the mass-weighted potential temperature and z is the height.

Equations of Motion: Lagrangian Vertical Coordinate

Equations of Motion

Choose $\mathcal{F} = \int \hat{\mu}$ (or $\int \hat{v} / \int \hat{\Theta} / \int \hat{z}$) to get:

$$\int \hat{\mu} \left(\frac{\partial \mu}{\partial t} + \vec{\nabla} \cdot \left(\frac{\delta \mathcal{H}}{\delta \vec{v}} \right) \right) = 0 \quad (8)$$

$$\int \hat{\Theta} \left(\frac{\partial \Theta}{\partial t} + \vec{\nabla} \cdot \left(\theta \frac{\delta \mathcal{H}}{\delta \vec{v}} \right) \right) = 0 \quad (9)$$

$$\int \hat{v} \left(\frac{\partial \vec{v}}{\partial t} + \frac{\zeta_v}{\mu} \times \frac{\delta \mathcal{H}}{\delta \vec{v}} + \theta \vec{\nabla} \left(\frac{\delta \mathcal{H}}{\delta \Theta} \right) + \vec{\nabla} \left(\frac{\delta \mathcal{H}}{\delta \mu} \right) \right) = 0 \quad (10)$$

$$\int \hat{z} \frac{\delta \mathcal{H}}{\delta z} = \int \hat{z} \left(g\mu + \frac{\partial p}{\partial \eta} \right) = 0 \quad (11)$$

Note that these are ALL 2D except for hydrostatic balance (11)

Hamiltonian (Lagrangian Vertical Coordinate)

Hamiltonian and Functional Derivatives

$$\mathcal{H} = \mathcal{H}[\mu, \vec{v}, \Theta, z] = \int \mu \left(\frac{\vec{u} \cdot \vec{u}}{2} + U \left(\frac{1}{\mu} \frac{\partial z}{\partial \eta}, \frac{\Theta}{\mu} \right) + gz \right) \quad (12)$$

$$\int \hat{v} \frac{\delta \mathcal{H}}{\delta \vec{v}} = \int \hat{v} (\mu \vec{u}) \quad (13)$$

$$\int \hat{\mu} \frac{\delta \mathcal{H}}{\delta \mu} = \int \hat{\mu} \left(\frac{\vec{u} \cdot \vec{u}}{2} + gz \right) \quad (14)$$

$$\int \hat{\Theta} \frac{\delta \mathcal{H}}{\delta \Theta} = \int \hat{\Theta} \frac{\partial U}{\partial \theta} = \int \hat{\Theta} \pi \quad (15)$$

$$\int \hat{z} \frac{\delta \mathcal{H}}{\delta z} = \int \hat{z} \left(g\mu + \frac{\partial p}{\partial \eta} \right) \quad (16)$$

Conservation

Energy

- Arises purely from anti-symmetry of the brackets PLUS $\frac{\delta \mathcal{H}}{\delta z} = 0$
- Mimetic Galerkin methods automatically ensure an anti-symmetric bracket
- Works for ANY choice of \mathcal{H}
- Something similar can be done with a mass-based vertical coordinate, although it is slightly more complicated

Mass and Entropy

- These are Casimirs
- Can show that this discretization also conserves them

Remaining Issues and Questions

Hydrostatic Balance

- 1 Dynamico: Column-wise direct solution
- 2 Can this also be done with Galerkin approach?

Grid Staggering: Placement of Θ

- 1 Dynamico: Lorenz staggering (Θ and μ are collocated)
- 2 Galerkin Equivalent: $\mu, \Theta \in \mathbb{W}_3$ (Admits a spurious computational mode in the vertical)
- 3 Charney-Phillips: $\Theta \in \mathbb{W}_{2,vert}$ (Avoids computational mode, more difficult to obtain conservation, complicates formulation)
- 4 Differential Geometry: Θ is a 0-form $\rightarrow \Theta \in \mathbb{W}_0$ (Excessive horizontal averaging \rightarrow computational mode/poor dispersion properties?)

Summary and Conclusions

Summary and Conclusions

Summary

- 1 There is a general, effective procedure for devising numerical schemes that mimic important characteristics of the real atmosphere
- 2 **Combine mimetic discretizations with a Hamiltonian formulation**

Future Work

- 1 Computational efficiency: preconditioning, assembly, etc.
- 2 Mass-based vertical coordinate
- 3 Non-hydrostatic equations
- 4 Subgrid Turbulence, Moisture, Tracers, Physics Coupling