# Higher-order Structure-Preserving Finite Elements for Atmospheric Dynamical Cores

#### Chris Eldred, Thomas Dubos and Evaggelos Kritsikis University of Paris 13

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# Introduction

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### (Incomplete) List of Desirable Model Properties



# Incomplete History of Structure-Preserving Schemes for Atmospheric Models

- The Beginning: Arakawa and Lamb 1981, Sadourney 1975
- Mimetic Finite Differences: Ringler et. al 2010; Thuburn et. al 2012, 2014, many others
- Mimetic Finite Elements: Cotter et. al 2012,2013; McRae et. al 2014, Thuburn + Cotter 2015, many others
- Hamiltonian Methods: Salmon 2004,2005,2007; Sommer+Nevir 2009; Gassmann 2008,2013; Dubos et. al 2015; Tort et. al 2015, many others

What is structure-preservation?

 Mimetic Discretization: curl-free pressure gradients, discrete product rules, discrete deRham cohomology, etc.

$$\vec{\nabla} \times \vec{\nabla} = 0$$

$$\vec{\nabla}\cdot\vec{\nabla} imes=0$$

$$(ec{
abla}\cdot)^* = -ec{
abla}$$

Quasi-Hamiltonian system: conserves mass, total energy and possibly other invariants

$$\frac{d\mathcal{H}}{dt} = 0$$
$$\frac{d\mathcal{C}}{dt} = 0$$

# Mimetic Discretizations

# General Formulation for Mimetic Discretizations: Primal-Dual Double deRham Complex (Staggered Grids)



$$egin{aligned} \delta &= * d * \ & 
abla^2 &= d \delta + \delta d \ & ec 
abla \cdot ec 
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$$\int_{\Omega} dW = \int_{d\Omega} W$$
$$dd = 0 = \delta\delta$$

# General Formulation for Mimetic Discretizations: Primal deRham Complex



$$\delta = *d*$$
 $abla^2 = d\delta + \delta d$ 
 $end{v} \cdot \vec{
abla} \times = 0 = \vec{
abla} \times \vec{
abla}$ 
 $dd = 0 = \delta\delta$ 



# Hamiltonian

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Non-Canonical Hamiltonian Dynamics

Evolution of an arbitrary functional  $\mathcal{F} = \mathcal{F}[\vec{x}]$  is governed by:

$$\frac{d\mathcal{F}}{dt} = \{\frac{\delta\mathcal{F}}{\delta\vec{x}}, \frac{\delta\mathcal{H}}{\delta\vec{x}}\}$$
(1)

with Poisson bracket  $\{,\}$  antisymmetric (also satisfies Jacobi):

$$\{\frac{\delta\mathcal{F}}{\delta\vec{x}},\frac{\delta\mathcal{G}}{\delta\vec{x}}\} = -\{\frac{\delta\mathcal{G}}{\delta\vec{x}},\frac{\delta\mathcal{F}}{\delta\vec{x}}\}$$
(2)

Also have Casimirs C that satisfy:

$$\{\frac{\delta \mathcal{F}}{\delta \vec{x}}, \frac{\delta \mathcal{C}}{\delta \vec{x}}\} = 0 \quad \forall \mathcal{F}$$
(3)

Neatly encapsulates conservation properties ( $\mathcal{H}$  and  $\mathcal{C}$ ).

## Recap: Mimetic + Hamiltonian = Structure-Preserving

- Fundamentally, structure-preserving schemes can be viewed as a combination of a **mimetic discretization method** plus a **hamiltonian formulation**
- Example: TRiSK Scheme (Primal-Dual ie Staggered Grid):

$$\frac{\partial \mathcal{F}}{\partial t} = \left(\frac{\delta \mathcal{A}}{m_i}, D_2 \frac{\delta \mathcal{B}}{u_e}\right)_{\mathbf{I}} + \left(\frac{\delta \mathcal{A}}{u_e}, \bar{D_1} \frac{\delta \mathcal{B}}{m_i}\right)_{\mathbf{H}} + \left(\frac{\delta \mathcal{A}}{u_e}, \mathbf{Q} \frac{\delta \mathcal{B}}{u_e}\right)_{\mathbf{H}}$$

- D<sub>2</sub>, D
  <sub>1</sub> and Q are mimetic operators (Q is a little complicated)
- $\bullet~(,)_{I},(,)_{J},(,)_{H}$  are inner products (induced by Hodge stars)
- Scheme conserves mass, energy, potential vorticity; has curl-free pressure gradients, steady geostrophic modes, etc.

### Dynamico



- Primal-Dual: Mimetic finite differences (based on TRiSK scheme): C grid horizontal, Lorenz vertical
- Icosahedral grid
- Hydrostatic primitive equations: Lagrangian and mass-based vertical coordinates
- Conserves mass, energy and entropy
- See Dubos et. al 2015 for more information

Reconstruction Operator  $(\mathbf{W})$  in TRiSK

$$\mathbf{W} = \sum_{e' \in ECP(e)} W_{e,e'}$$
$$\mathbf{W} = -\mathbf{W}^{T}$$
$$-\mathbf{R}D_2 = \bar{D}_2\mathbf{W}$$

Given normal fluxes, reconstruct tangential fluxes Satisfying: Steady geostrophic modes AND energy conservation AND accuracy



### Issues with TRiSK

#### **Operator Accuracy**



#### Spurious Branches of Dispersion Relationship

Hexagonal grid means 3:1 ratio of wind to mass dofs (should be 2:1)  $\rightarrow$  spurious branch of Rossby waves with unphysical behaviour

#### How do we fix them?



- Fix spurious branches: Quadrilateral (cubed-sphere) grid, correct 2:1 ratio of dofs
- Fix accuracy: Use Primal approach (mimetic Galerkin methods)
- Keep the same equations and Hamiltonian structure
- Keep the same mimetic and conservation properties

# Mimetic Galerkin Discretization

#### Mimetic Galerkin Discretization



$$\delta = *d*$$

$$\nabla^2 = d\delta + \delta d$$

$$\vec{\nabla} \cdot \vec{\nabla} \times = 0 = \vec{\nabla} \times \vec{\nabla}$$

$$dd = 0 = \delta\delta$$



# General Approach to Mimetic Galerkin Spaces

#### Mimetic Spaces

Select 1D Spaces  $\mathcal{A}$  and  $\mathcal{B}$  such that  $: \mathcal{A} \xrightarrow{\frac{d}{dx}} \mathcal{B}$ 

(4)

- Use tensor products to extend to n-dimensions
- Works for ANY set of spaces A and B that satisfy this property (mimetic finite elements use P<sub>n</sub> and P<sub>DG,n-1</sub>)
- Mimetic spectral element, Mimetic isogeometric methods (B-splines) all fall under this framework
- Our (different) choices of  $\mathcal{A}$  and  $\mathcal{B}$  are guided by linear mode properties and coupling to physics/tracer transport
- See Hiemstra et. al 2014 (and references therein)

 $P_2 - P_{1,DG}$  Dispersion Relationship



#### Mimetic Galerkin Differences: Basis

![](_page_19_Figure_2.jpeg)

Single degree of freedom per geometric entity (physics coupling) Higher order by larger stencils (less local) 3rd Order Elements

### Mimetic Galerkin Differences- Dispersion

![](_page_20_Figure_2.jpeg)

Inertia-Gravity Wave Dispersion Relationship (1D) for 3rd Order Elements

#### Spectral gap is gone

Can show that dispersion relation is O(2n) where *n* is the order More details in a forthcoming paper

#### Overview of 3D Spaces

![](_page_21_Figure_2.jpeg)

$$\mathbb{W}_0 \xrightarrow{\vec{\nabla}} \mathbb{W}_1 \xrightarrow{\vec{\nabla} \times} \mathbb{W}_2 \xrightarrow{\vec{\nabla} \cdot} \mathbb{W}_3$$

$$\begin{split} \mathbb{W}_0 &= \mathcal{A} \otimes \mathcal{A} \otimes \mathcal{A} = H_1 = \text{Continuous Galerkin} \\ \mathbb{W}_1 &= (\mathcal{B} \otimes \mathcal{A} \otimes \mathcal{A})\hat{i} + \ldots = H(\textit{curl}) = \text{Nedelec} \\ \mathbb{W}_2 &= (\mathcal{A} \otimes \mathcal{B} \otimes \mathcal{B})\hat{i} + \ldots = H(\textit{div}) = \text{Raviart-Thomas} \\ \mathbb{W}_3 &= \mathcal{B} \otimes \mathcal{B} \otimes \mathcal{B} = L_2 = \text{Discontinuous Galerkin} \end{split}$$

## Grid Staggering for HPE

![](_page_22_Figure_2.jpeg)

Follows from differential geometry and Tonti diagram Galerkin Version of a C Grid Question: Where should ⊖ live?

# Hamiltonian Formulation

# Poisson Brackets (Lagrangian Vertical Coordinate)

#### Poisson Brackets

From Dubos and Tort 2014, evolution of  $\mathcal{F}[\vec{x}] = \mathcal{F}[\mu, \vec{v}, \Theta, z]$  is

$$\frac{d\mathcal{F}}{dt} = \{\frac{\delta\mathcal{F}}{\delta\vec{x}}, \frac{\delta\mathcal{H}}{\delta\vec{x}}\}_{SW} + \{\frac{\delta\mathcal{F}}{\delta\vec{x}}, \frac{\delta\mathcal{H}}{\delta\vec{x}}\}_{\Theta} + \langle\frac{\delta\mathcal{F}}{\delta z}\frac{\partial z}{\partial t}\rangle$$
(5)  
$$\{\frac{\delta\mathcal{F}}{\delta\vec{x}}, \frac{\delta\mathcal{H}}{\delta\vec{x}}\}_{SW} = \langle\frac{\delta\mathcal{H}}{\delta\vec{v}} \cdot \vec{\nabla}\frac{\delta\mathcal{F}}{\delta\mu} - \frac{\delta\mathcal{H}}{\delta\vec{v}} \cdot \vec{\nabla}\frac{\delta\mathcal{F}}{\delta\mu}\rangle + \langle\frac{\vec{\nabla} \times \vec{v}}{\mu} \cdot (\frac{\delta\mathcal{F}}{\delta\vec{v}} \times \frac{\delta\mathcal{H}}{\delta\vec{v}})\rangle$$
(6)  
$$\{\frac{\delta\mathcal{F}}{\delta\vec{x}}, \frac{\delta\mathcal{H}}{\delta\vec{x}}\}_{\Theta} = \langle\theta(\frac{\delta\mathcal{H}}{\delta\vec{v}} \cdot \vec{\nabla}\frac{\delta\mathcal{F}}{\delta\Theta} - \frac{\delta\mathcal{H}}{\delta\vec{v}} \cdot \vec{\nabla}\frac{\delta\mathcal{F}}{\delta\Theta})\rangle$$
(7)  
where  $\mu$  is the pseudo-density,  $\vec{v} = \vec{u} - \vec{R}$  is the absolute  
(covariant) velocity,  $\Theta = u\theta$  is the mass-weighted potential

temperature and z is the height.

### Equations of Motion: Lagrangian Vertical Coordinate

#### Equations of Motion

Choose  $\mathcal{F} = \int \hat{\mu}$  ( or  $\int \hat{v} / \int \hat{\Theta} / \int \hat{z}$ ) to get:  $\int \hat{\mu} \left( \frac{\partial \mu}{\partial t} + \vec{\nabla} \cdot \left( \frac{\delta \mathcal{H}}{\delta \vec{v}} \right) \right) = 0$ (8) $\int \hat{\Theta} \left( \frac{\partial \Theta}{\partial t} + \vec{\nabla} \cdot \left( \theta \frac{\delta \mathcal{H}}{\delta \vec{v}} \right) \right) = 0$ (9)  $\int \hat{v} \left( \frac{\partial \vec{v}}{\partial t} + \frac{\zeta_{v}}{\mu} \times \frac{\delta \mathcal{H}}{\delta \vec{v}} + \theta \vec{\nabla} (\frac{\delta \mathcal{H}}{\delta \Theta}) + \vec{\nabla} (\frac{\delta \mathcal{H}}{\delta \mu}) \right) = 0$ (10) $\int \hat{z} \frac{\delta \mathcal{H}}{\delta z} = \int \hat{z} \left( g \mu + \frac{\partial p}{\partial n} \right) = 0$ (11)

Note that these are ALL 2D except for hydrostatic balance (11)

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### Hamiltonian (Lagrangian Vertical Coordinate)

#### Hamiltonian and Functional Derivatives

$$\mathcal{H} = \mathcal{H}[\mu, \vec{v}, \Theta, z] = \int \mu(\frac{\vec{u} \cdot \vec{u}}{2} + U(\frac{1}{\mu} \frac{\partial z}{\partial \eta}, \frac{\Theta}{\mu}) + gz) \qquad (12)$$
$$\int \hat{v} \frac{\delta \mathcal{H}}{\delta \vec{v}} = \int \hat{v} (\mu \vec{u}) \qquad (13)$$

$$\int \hat{\mu} \frac{\delta \mathcal{H}}{\delta \mu} = \int \hat{\mu} \left( \frac{\vec{u} \cdot \vec{u}}{2} + gz \right)$$
(14)

$$\int \hat{\Theta} \frac{\delta \mathcal{H}}{\delta \Theta} = \int \hat{\Theta} \frac{\partial U}{\partial \theta} = \int \hat{\Theta} \pi$$
(15)

$$\int \hat{z} \frac{\delta \mathcal{H}}{\delta z} = \int \hat{z} \left( g \mu + \frac{\partial \rho}{\partial \eta} \right)$$
(16)

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## Conservation

#### Energy

- Arises purely from anti-symmetry of the brackets PLUS  $\frac{\delta \mathcal{H}}{\delta z}=0$
- Mimetic Galerkin methods automatically ensure an anti-symmetric bracket
- $\bullet$  Works for ANY choice of  ${\cal H}$
- Something similar can be done with a mass-based vertical coordinate, although it is slightly more complicated

#### Mass, Potential Vorticity and Entropy

- These are Casimirs
- Can show that this discretization also conserves them

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# Remaining Issues and Questions

#### Hydrostatic Balance

- Oynamico: Column-wise direct solution
- ② Can this also be done with Galerkin approach?

#### Grid Staggering: Placement of $\Theta$

- Dynamico: Lorenz staggering ( $\Theta$  and  $\mu$  are collocated)
- **③** Galerkin Equivalent:  $\mu, \Theta \in W_3$  (Admits a spurious computational mode in the vertical)
- Solution
   Charney-Phillips: Θ ∈ W<sub>2,vert</sub> (Avoids computational mode, complicates formulation)
- Differential Geometry: Θ is a 0-form → Θ ∈ W<sub>0</sub> (Excessive horizontal averaging → computational mode/poor dispersion properties?)

# Summary and Conclusions

# Summary and Conclusions

#### Summary

- There is a general, effective procedure for devising mimetic, conservative numerical schemes:
- **2** Mimetic discretizations + Hamiltonian formulation

#### Future Work

- Computational efficiency: preconditioning/solvers, matrix assembly
- 2 Mass-based vertical coordinate
- Onhydrostatic equations
- Past Inivscid, Adiabatic Dry Dynamics: Subgrid Turbulence, Moisture, Tracers, Physics Coupling